



دانشگاه سمنان

Semnan University  
Faculty of Mechanical Engineering

دانشکده مهندسی مکانیک



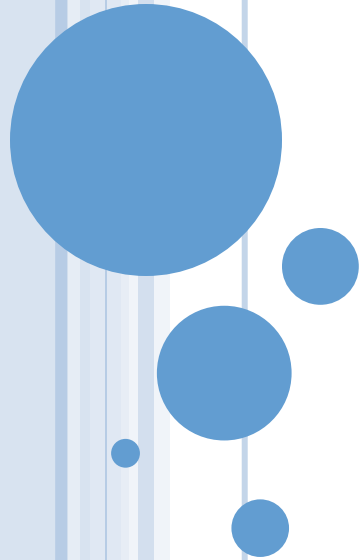
دانشکده مهندسی مکانیک

درس مکاترونیک ۱

# MECHATRONICS 1

Section 3:

**Dynamics Review**



**Reference:**

**Engineering Mechanics  
Dynamics**

9th Edition

Meriam, Kraige & Bolton

Chapters 1 - 6



□ CONTENTS:

- ➔ ❖ Chapter 1: **Introduction to Dynamics**
- ❖ Chapter 2: Kinematics of Particles
- ❖ Chapter 3: Kinetics of Particles
- ❖ Chapter 4: Kinetics of Systems of Particles
- ❖ Chapter 5: Plane Kinetics of Rigid Bodies
- ❖ Chapter 6: Plane Kinematics of Rigid Bodies

1/1

# History and Modern Applications

## History of Dynamics



© Fine Art Images/SuperStock

Galileo Galilei  
Portrait of Galileo Galilei  
(1564–1642) (oil on canvas),  
Sustermans, Justus  
(1597–1681) (school of)/  
Galleria Palatina, Florence,  
Italy/Bridgeman Art Library.

## Applications of Dynamics

WENN Ltd/Alamy Stock Photo



Artificial hand



1/2

## Basic Concepts

- ❑ **Space**
- ❑ ***Time***
- ❑ ***Mass***
- ❑ ***Force***
- ❑ ***A Particle***
- ❑ ***A Rigid Body***
- ❑ ***Vector* and *Scalar* quantities**



**1/3** Newton's Laws

- **Law I.** A particle remains at rest or continues to move with uniform velocity (in a straight line with a constant speed) if there is no unbalanced force acting on it.
- **Law II.** The acceleration of a particle is proportional to the resultant force acting on it and is in the direction of this force.

$$\mathbf{F} = m\mathbf{a}$$

- **Law III.** The forces of action and reaction between interacting bodies are equal in magnitude, opposite in direction, and collinear.



## 1/4 Units

- International System of metric units (SI)
- U.S. Customary system of units

Quantity	Dimensional Symbol	SI Units		U.S. Customary Units			
		Unit	Symbol	Unit	Symbol		
Mass	M	Base units	kilogram	kg	slug	—	
Length	L		meter*	m	Base units	foot	ft
Time	T		second	s		second	sec
Force	F		newton	N	pound	lb	

\*Also spelled *metre*.



The U.S. standard kilogram at the National Bureau of Standards.

Omikron/Photo Researchers, Inc.

SI Units	U.S. Customary Units
$(1 \text{ N}) = (1 \text{ kg})(1 \text{ m/s}^2)$	$(1 \text{ lb}) = (1 \text{ slug})(1 \text{ ft/sec}^2)$
$\text{N} = \text{kg} \cdot \text{m/s}^2$	$\text{slug} = \text{lb} \cdot \text{sec}^2/\text{ft}$

## 1/5 Gravitation

- Newton's law of gravitation, which governs the mutual attraction between bodies

$$F = G \frac{m_1 m_2}{r^2}$$

$F$  = the mutual force of attraction between two particles

$G$  = a universal constant called the *constant of gravitation*

$m_1, m_2$  = the masses of the two particles

$r$  = the distance between the centers of the particles

$$G = 6.673(10^{-11}) \text{ m}^3/(\text{kg} \cdot \text{s}^2)$$

$$\mathbf{W} = m\mathbf{g}$$

9.81 m/s<sup>2</sup> in SI units

32.2 ft/sec<sup>2</sup> in U.S. customary units



1/6

## Dimensions

- A given dimension such as length can be expressed in a number of different units such as meters, millimeters, or kilometers.
- Thus, a *dimension* is different from a *unit*.
- The *principle of dimensional homogeneity* states that all physical relations must be dimensionally homogeneous

$$F = ML/T^2$$



❑ CONTENTS:

- ❖ Chapter 1: Introduction to Dynamics
- ➔ ❖ Chapter 2: **Kinematics of Particles**
- ❖ Chapter 3: Kinetics of Particles
- ❖ Chapter 4: Kinetics of Systems of Particles
- ❖ Chapter 5: Plane Kinetics of Rigid Bodies
- ❖ Chapter 6: Plane Kinematics of Rigid Bodies

## 2/1 Introduction

- Kinematics is the branch of dynamics which describes the motion of bodies without reference to the forces which either cause the motion or are generated as a result of the motion.
- Kinematics is often described as the “geometry of motion.”
- A thorough working knowledge of kinematics is a prerequisite to kinetics, which is the study of the relationships between motion and the corresponding forces which cause or accompany the motion.



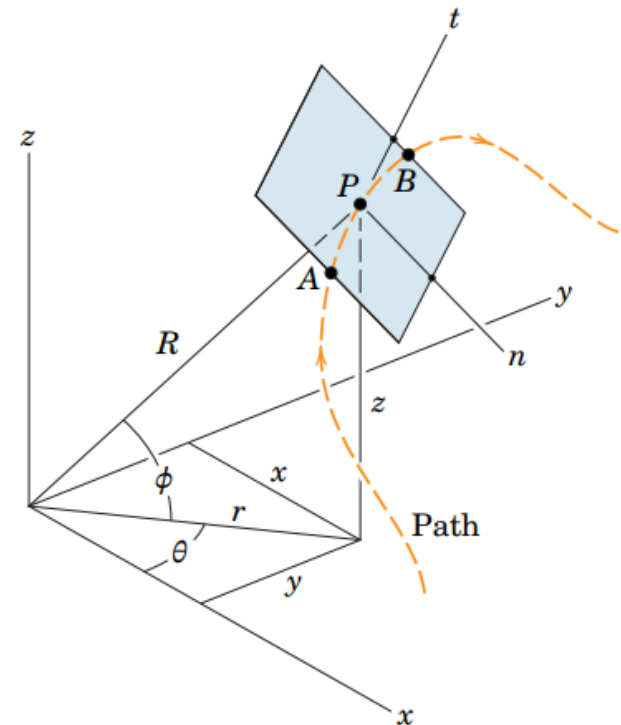
## 2/1 Introduction

### Particle Motion

- ❖ A particle is a body whose physical dimensions are so small compared with the radius of curvature of its path that we may treat the motion of the particle as that of a point.

### Choice of Coordinates

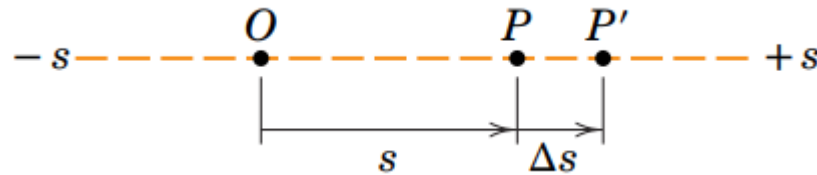
- ❖ Rectangular coordinates  $x, y, z$  (Cartesian)
- ❖ Cylindrical coordinates  $r, \theta, z$
- ❖ Spherical coordinates  $R, \theta, \phi$
- ❖ Tangent  $t$  and normal  $n$  to the curve (path variables)





## 2/2 Rectilinear Motion

- Particle P moving along a straight line



## Velocity and Acceleration

- Average velocity of the particle during the interval  $\Delta t$  is the displacement divided by the time interval or  $v_{av} = \Delta s / \Delta t$

## 2/2 Rectilinear Motion

### Velocity and Acceleration

❖ *instantaneous velocity*

$$v = \frac{ds}{dt} = \dot{s}$$

❖ The average acceleration of the particle during the interval  $\Delta t$  is the change in its velocity divided by the time interval or  $a_{av} = \Delta v / \Delta t$ .

❖ *instantaneous acceleration*

$$a = \frac{dv}{dt} = \dot{v}$$

or

$$a = \frac{d^2s}{dt^2} = \ddot{s}$$



$$v dv = a ds$$

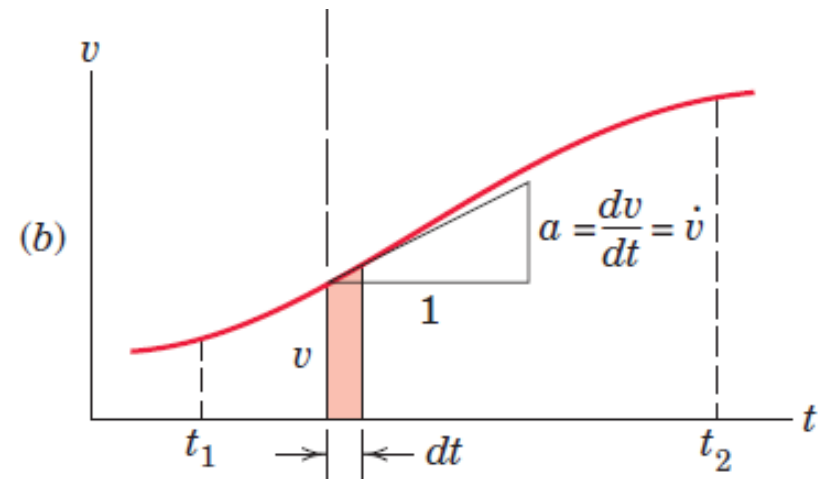
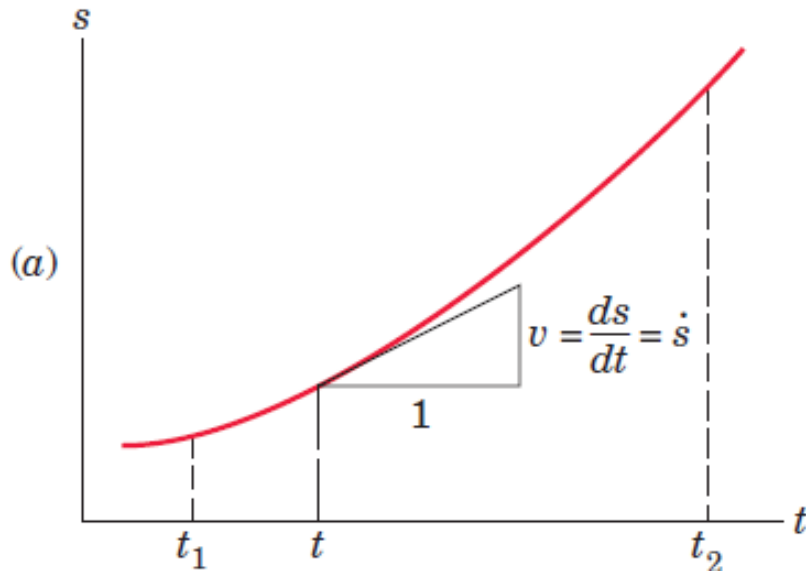
or

$$\dot{s} d\dot{s} = \ddot{s} ds$$



## 2/2 Rectilinear Motion

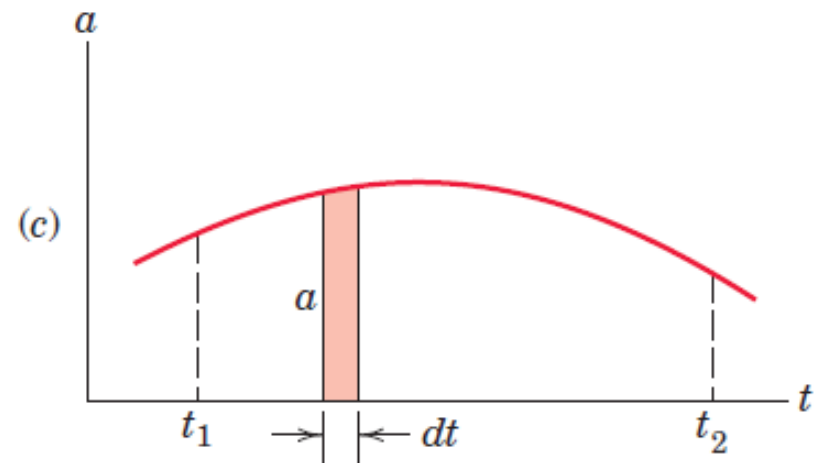
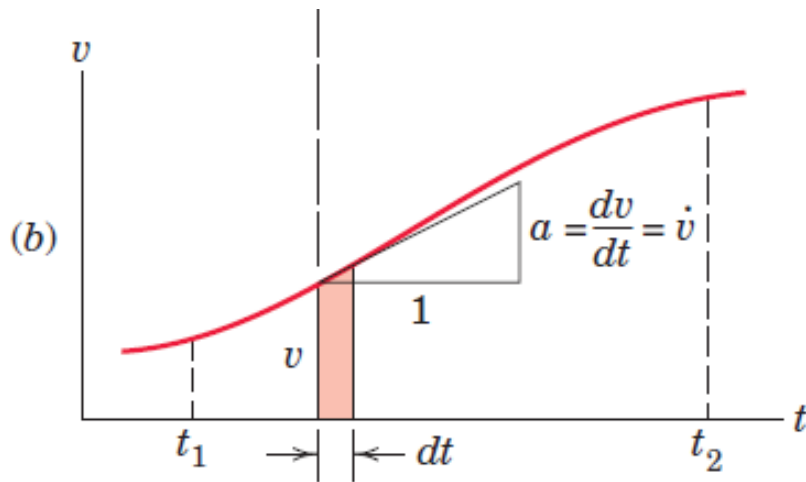
## Graphical Interpretations



$$\int_{s_1}^{s_2} ds = \int_{t_1}^{t_2} v dt \quad \text{or} \quad s_2 - s_1 = (\text{area under } v\text{-}t \text{ curve})$$

## 2/2 Rectilinear Motion

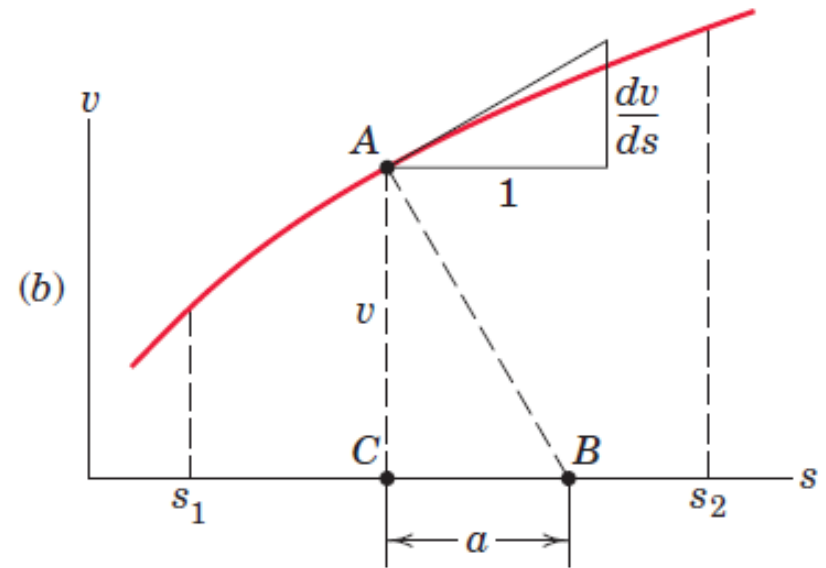
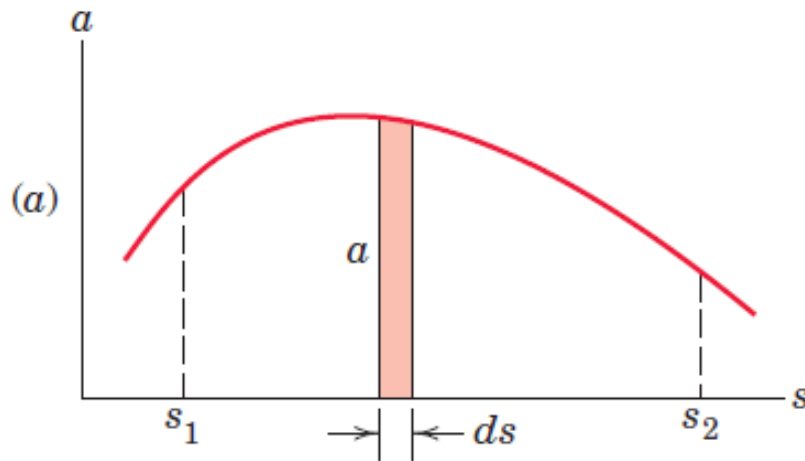
## Graphical Interpretations



$$\int_{v_1}^{v_2} dv = \int_{t_1}^{t_2} a dt \quad \text{or} \quad v_2 - v_1 = (\text{area under } a-t \text{ curve})$$

## 2/2 Rectilinear Motion

## Graphical Interpretations



$$\int_{v_1}^{v_2} v \, dv = \int_{s_1}^{s_2} a \, ds \quad \text{or} \quad \frac{1}{2}(v_2^2 - v_1^2) = (\text{area under } a\text{-}s \text{ curve})$$



2/2

## Rectilinear Motion

## Analytical Integration

□ (a) *Constant Acceleration*

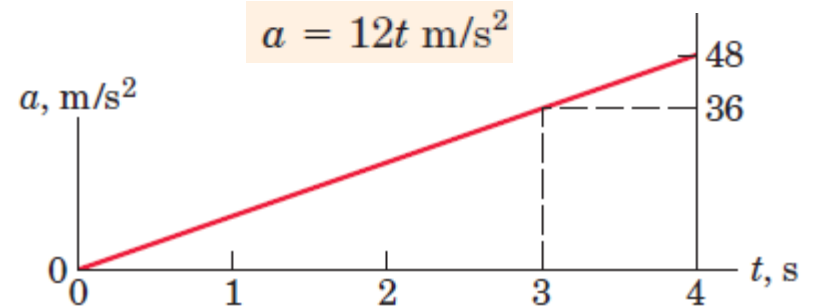
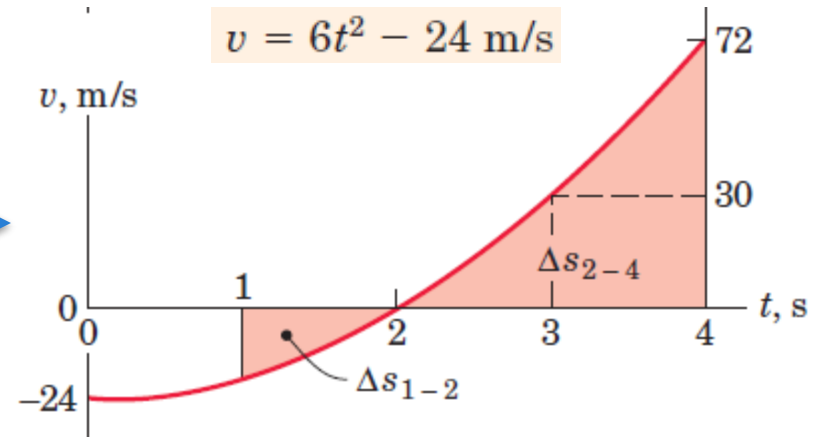
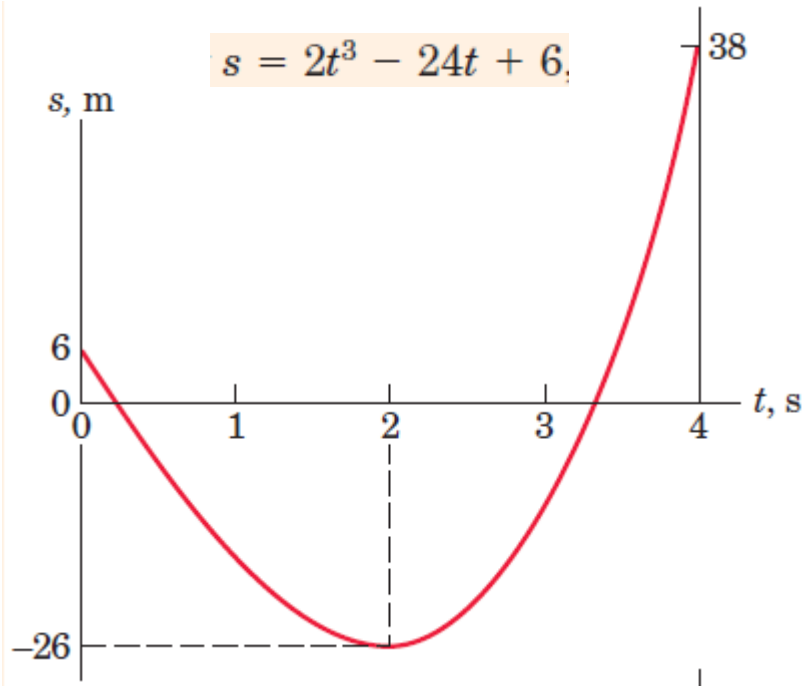
$$\int_{v_0}^v dv = a \int_0^t dt \quad \text{or} \quad v = v_0 + at$$

$$\int_{v_0}^v v dv = a \int_{s_0}^s ds \quad \text{or} \quad v^2 = v_0^2 + 2a(s - s_0)$$

$$\int_{s_0}^s ds = \int_0^t (v_0 + at) dt \quad \text{or} \quad s = s_0 + v_0 t + \frac{1}{2}at^2$$

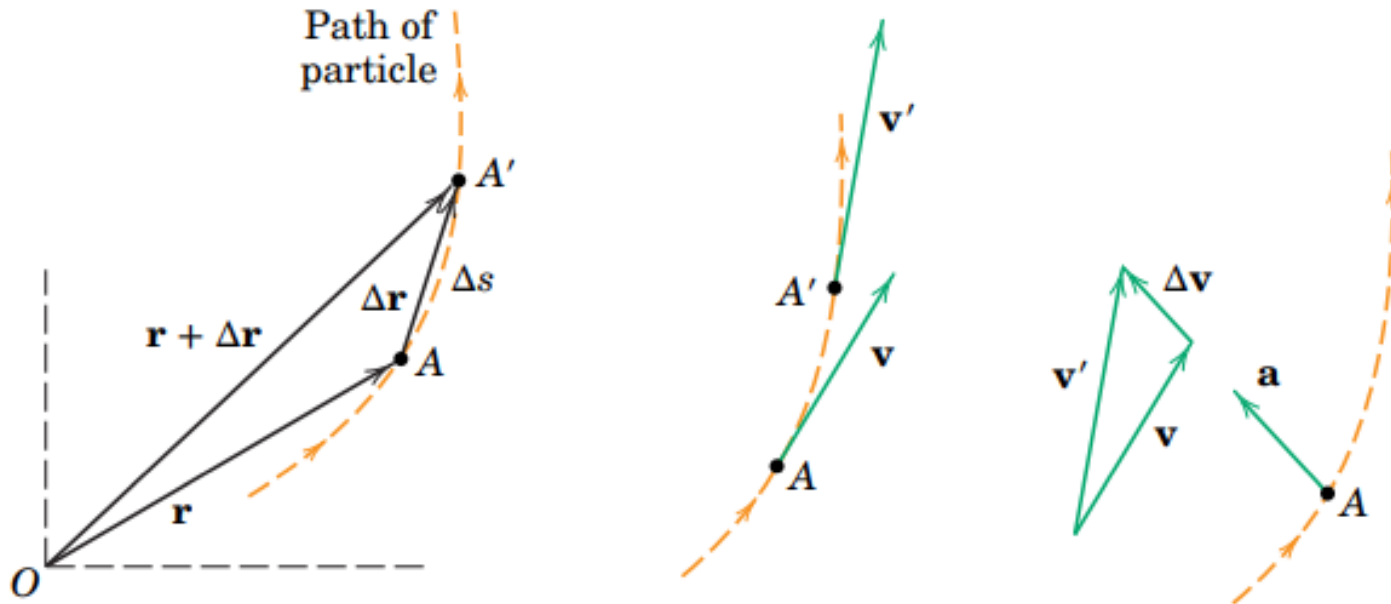


## SAMPLE PROBLEM 2/1



## 2/3 Plane Curvilinear Motion

### □ Position Vector $\mathbf{r}$





## 2/3 Plane Curvilinear Motion

### Velocity

- *Average velocity* of the particle between A and A' is defined as  $\mathbf{v}_{av} = \Delta \mathbf{r} / \Delta t$
- *instantaneous velocity* (approaches tangent to the path)

$$\mathbf{v} = \frac{d\mathbf{r}}{dt} = \dot{\mathbf{r}}$$

- The magnitude of  $\mathbf{v}$  is called the *speed* and is the scalar

$$v = |\mathbf{v}| = \frac{ds}{dt} = \dot{s}$$



## 2/3 Plane Curvilinear Motion

### Acceleration

- The *average acceleration* of the particle between  $A$  and  $A'$  is defined as  $\Delta \mathbf{v} / \Delta t$
- *instantaneous acceleration*

$$\mathbf{a} = \frac{d\mathbf{v}}{dt} = \dot{\mathbf{v}}$$

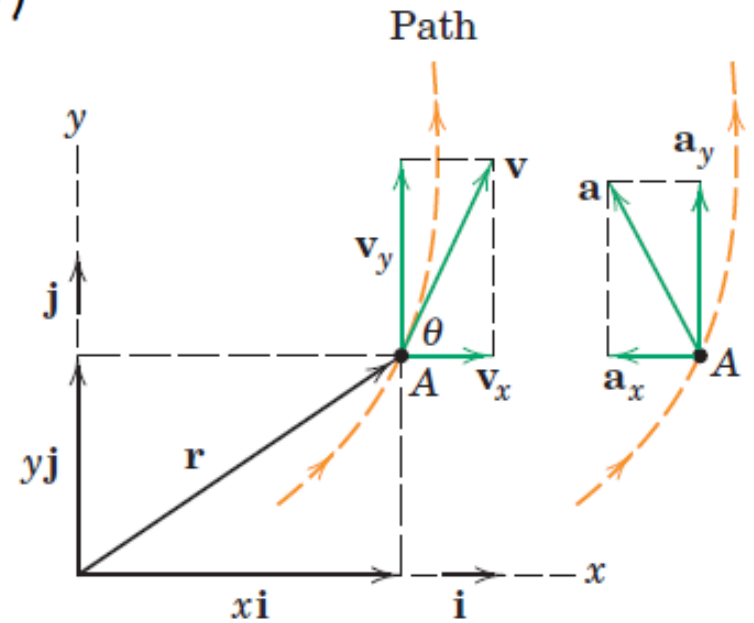
## 2/4 Rectangular Coordinates (x-y)

## Vector Representation

$$\mathbf{r} = x\mathbf{i} + y\mathbf{j}$$

$$\mathbf{v} = \dot{\mathbf{r}} = \dot{x}\mathbf{i} + \dot{y}\mathbf{j}$$

$$\mathbf{a} = \dot{\mathbf{v}} = \ddot{\mathbf{r}} = \ddot{x}\mathbf{i} + \ddot{y}\mathbf{j}$$



$$v^2 = v_x^2 + v_y^2 \quad v = \sqrt{v_x^2 + v_y^2} \quad \tan \theta = \frac{v_y}{v_x}$$

$$a^2 = a_x^2 + a_y^2 \quad a = \sqrt{a_x^2 + a_y^2}$$

## 2/4 Rectangular Coordinates (x-y)

### Projectile Motion

- An important application of two-dimensional kinematic theory is the problem of projectile motion.
- For a first treatment of the subject, we neglect aerodynamic drag and the curvature and rotation of the earth, and we assume that the altitude change is small enough so that the acceleration due to gravity can be considered constant.

$$a_x = 0 \quad a_y = -g$$



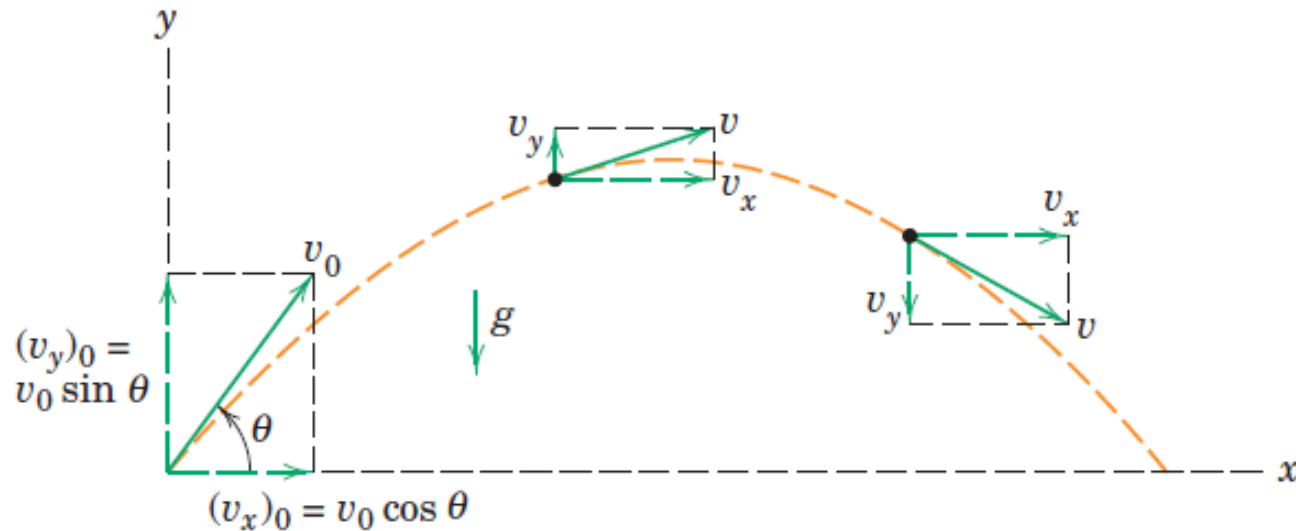
## 2/4 Rectangular Coordinates (x-y)

### Projectile Motion

$$v_x = (v_x)_0 \qquad v_y = (v_y)_0 - gt$$

$$x = x_0 + (v_x)_0 t \qquad y = y_0 + (v_y)_0 t - \frac{1}{2}gt^2$$

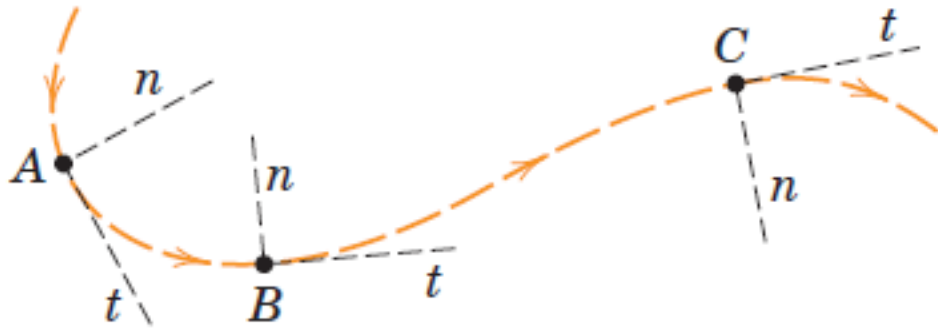
$$v_y^2 = (v_y)_0^2 - 2g(y - y_0)$$



2/5

## Normal and Tangential Coordinates ( $n-t$ )

- Measurements made along the tangent  $t$  and normal  $n$  to the path of the particle

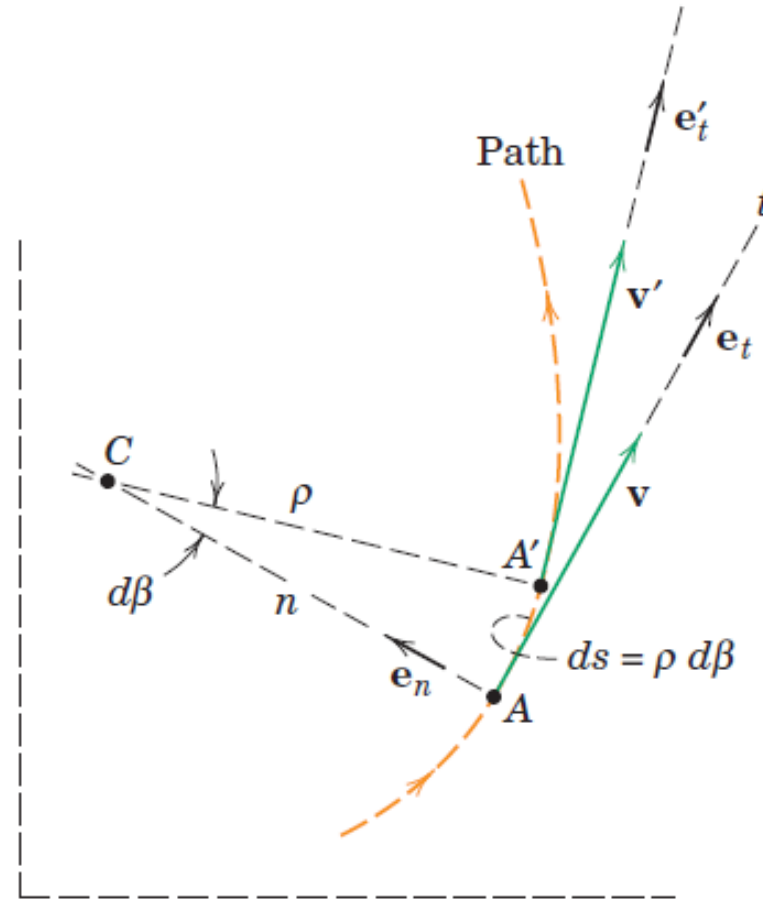


2/5

Normal and Tangential Coordinates ( $n-t$ )

## Velocity and Acceleration

$$\mathbf{v} = v\mathbf{e}_t = \rho\dot{\beta}\mathbf{e}_t$$



2/5

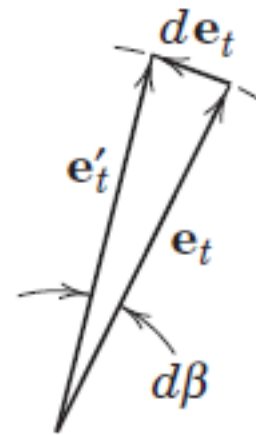
Normal and Tangential Coordinates ( $n-t$ )

## Velocity and Acceleration

$$\mathbf{a} = \frac{d\mathbf{v}}{dt} = \frac{d(v\mathbf{e}_t)}{dt} = v\dot{\mathbf{e}}_t + \dot{v}\mathbf{e}_t$$

$$d\mathbf{e}_t = \mathbf{e}_n d\beta$$

$$\dot{\mathbf{e}}_t = \dot{\beta}\mathbf{e}_n$$





2/5

Normal and Tangential Coordinates ( $n-t$ )

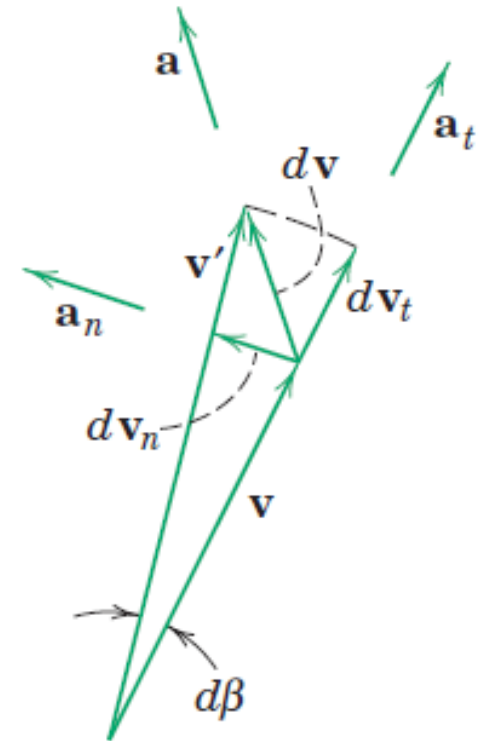
## Velocity and Acceleration

$$\mathbf{a} = \frac{v^2}{\rho} \mathbf{e}_n + \dot{v} \mathbf{e}_t$$

$$a_n = \frac{v^2}{\rho} = \rho \dot{\beta}^2 = v \dot{\beta}$$

$$a_t = \dot{v} = \ddot{s} \quad \rightarrow \quad a_t = \dot{v} = d(\rho \dot{\beta})/dt = \rho \ddot{\beta} + \dot{\rho} \dot{\beta}$$

$$a = \sqrt{a_n^2 + a_t^2}$$

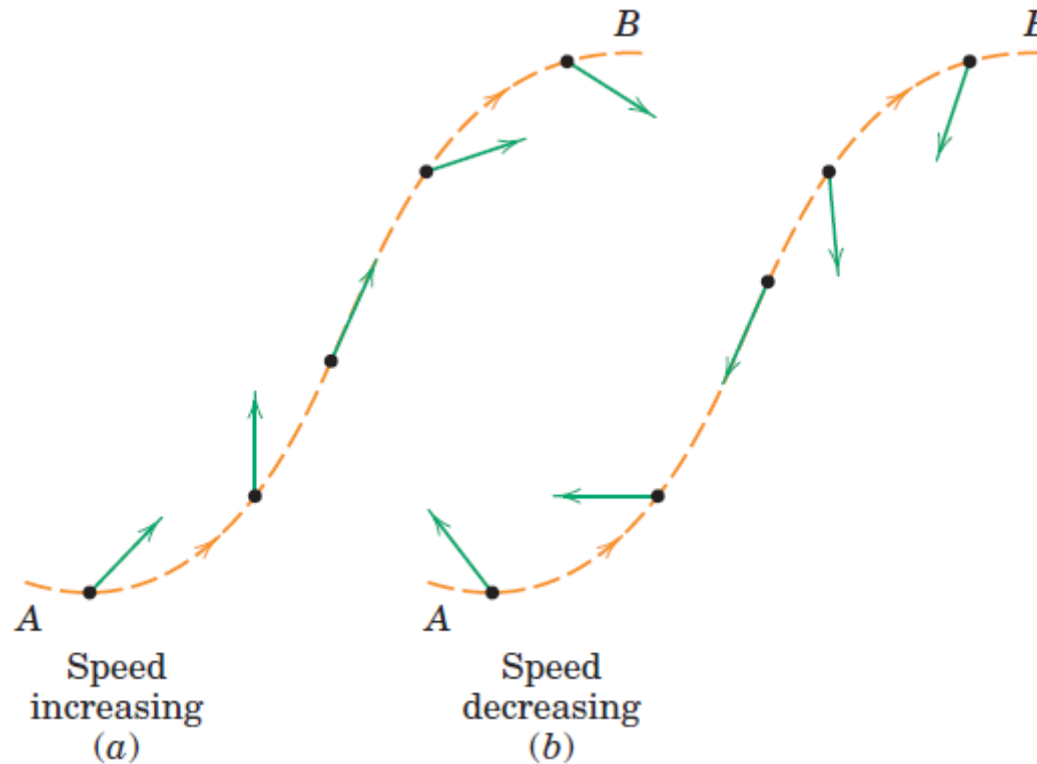


2/5

## Normal and Tangential Coordinates ( $n-t$ )

### Geometric Interpretation

❖  $a_n$  is always directed toward the center of curvature  $C$



2/5

## Normal and Tangential Coordinates ( $n-t$ )

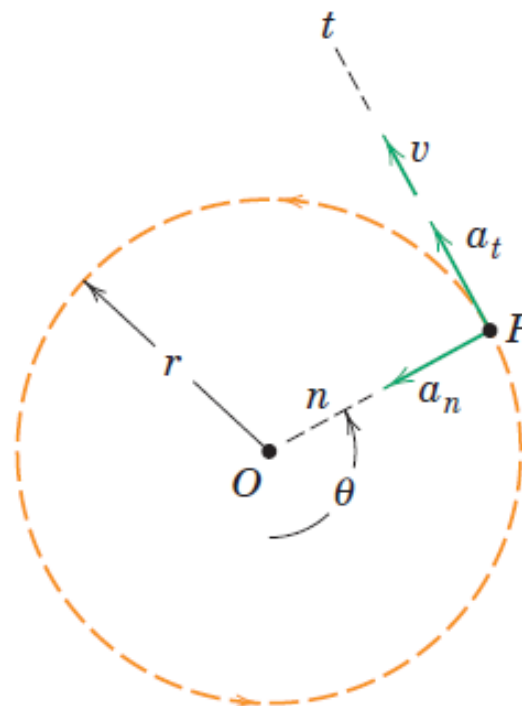
### Circular Motion

- ❖ Circular motion is an important special case of plane curvilinear motion where the radius of curvature  $\rho$  becomes the constant radius  $r$  of the circle.

$$v = r\dot{\theta}$$

$$a_n = v^2/r = r\dot{\theta}^2 = v\dot{\theta}$$

$$a_t = \dot{v} = r\ddot{\theta}$$

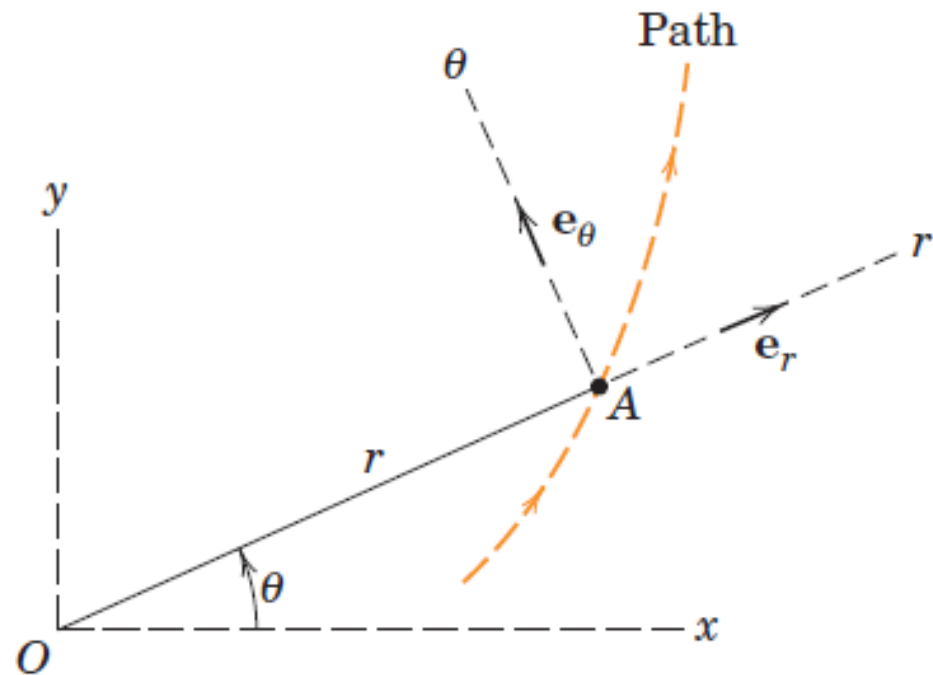


2/6

## Polar Coordinates ( $r-\theta$ )

- ❖ Particle is located by the radial distance  $r$  from a fixed point and by an angular measurement  $\theta$  to the radial line

$$\mathbf{r} = r\mathbf{e}_r$$



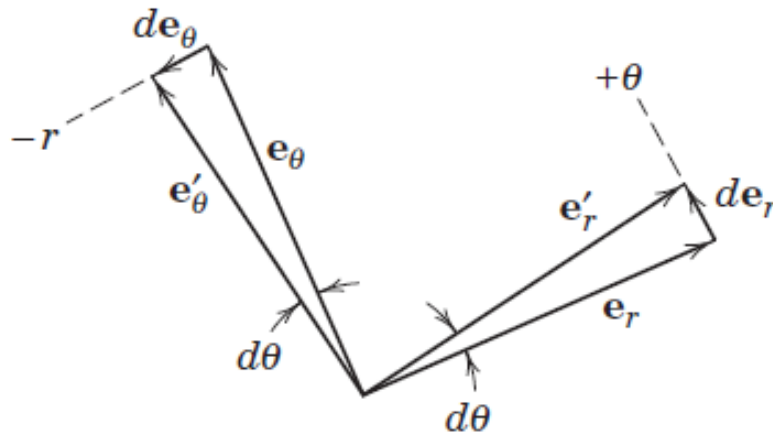
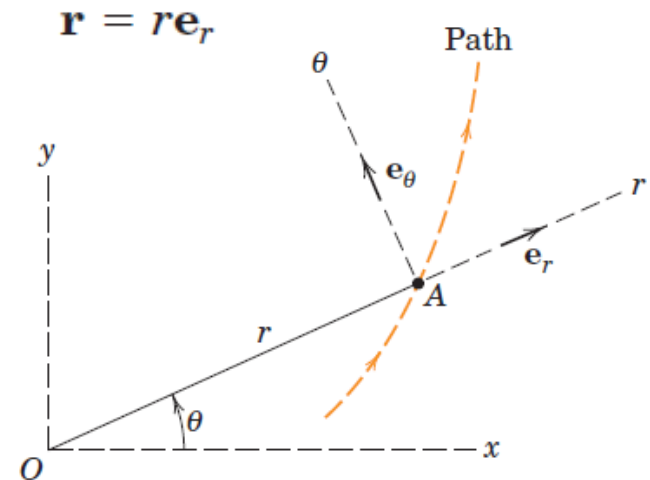
2/6

Polar Coordinates ( $r-\theta$ )

## Time Derivatives of the Unit Vectors

$$\frac{d\mathbf{e}_r}{d\theta} = \mathbf{e}_\theta \quad \text{and} \quad \frac{d\mathbf{e}_\theta}{d\theta} = -\mathbf{e}_r$$

$$\dot{\mathbf{e}}_r = \dot{\theta}\mathbf{e}_\theta \quad \text{and} \quad \dot{\mathbf{e}}_\theta = -\dot{\theta}\mathbf{e}_r$$



2/6

## Polar Coordinates ( $r$ - $\theta$ )

### Velocity

$$\mathbf{v} = \dot{\mathbf{r}} = \dot{r}\mathbf{e}_r + r\dot{\mathbf{e}}_r$$

$$\mathbf{v} = \dot{r}\mathbf{e}_r + r\dot{\theta}\mathbf{e}_\theta$$

$$v_r = \dot{r}$$

$$v_\theta = r\dot{\theta}$$

$$v = \sqrt{v_r^2 + v_\theta^2}$$

### Acceleration

$$\mathbf{a} = \dot{\mathbf{v}} = (\ddot{r}\mathbf{e}_r + \dot{r}\dot{\mathbf{e}}_r) + (\dot{r}\dot{\theta}\mathbf{e}_\theta + r\ddot{\theta}\mathbf{e}_\theta + r\dot{\theta}\dot{\mathbf{e}}_\theta)$$

$$\mathbf{a} = (\ddot{r} - r\dot{\theta}^2)\mathbf{e}_r + (r\ddot{\theta} + 2\dot{r}\dot{\theta})\mathbf{e}_\theta$$

$$a_r = \ddot{r} - r\dot{\theta}^2$$

$$a_\theta = r\ddot{\theta} + 2\dot{r}\dot{\theta}$$

$$a = \sqrt{a_r^2 + a_\theta^2}$$

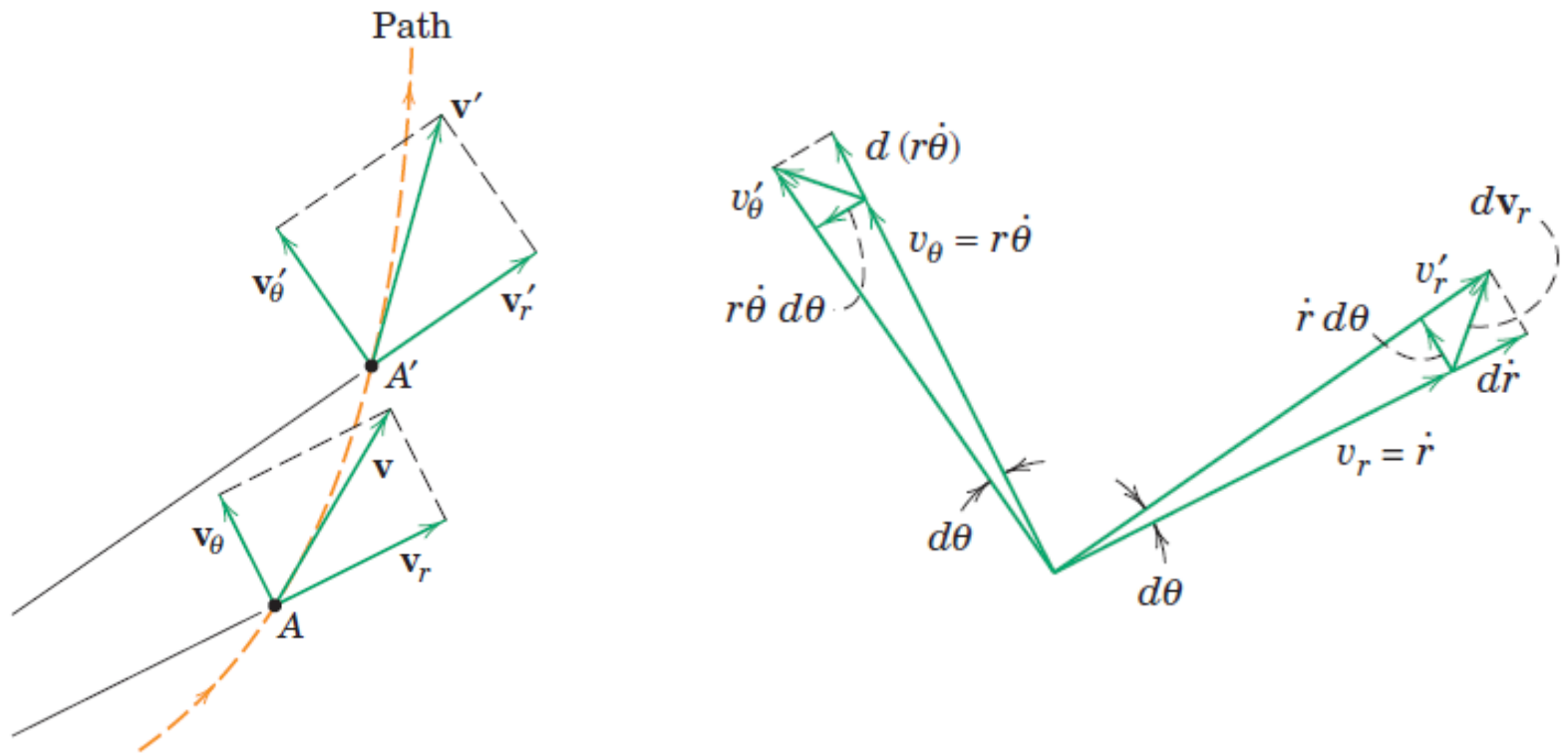
$$a_\theta = \frac{1}{r} \frac{d}{dt}(r^2\dot{\theta})$$



2/6

Polar Coordinates ( $r-\theta$ )

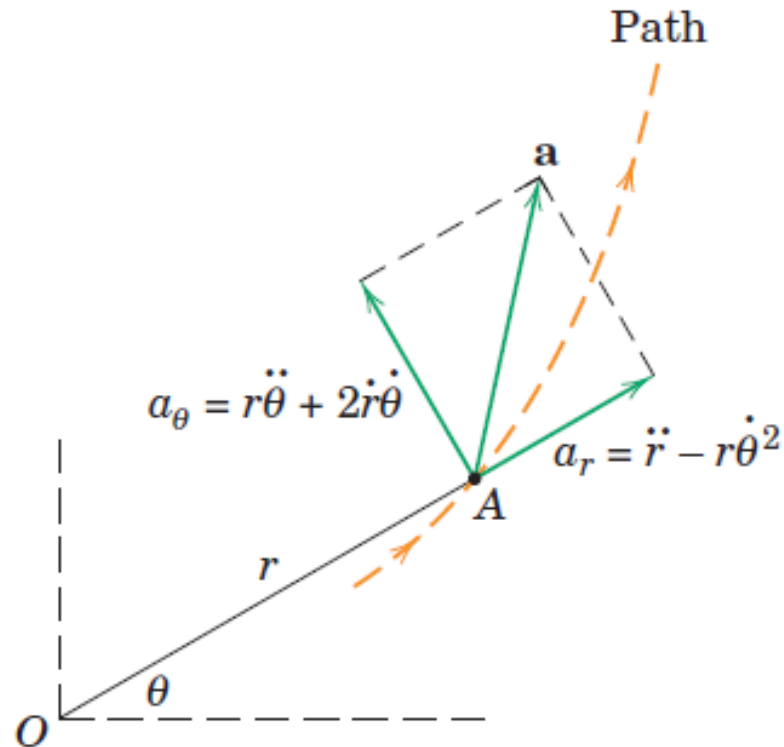
## Geometric Interpretation



2/6

# Polar Coordinates ( $r-\theta$ )

## Geometric Interpretation





2/6

## Polar Coordinates ( $r$ - $\theta$ )

### Circular Motion

- Same as that obtained with  $n$ - and  $t$ -components, where the  $n$ - and  $t$ -directions coincide but the positive  $r$ -direction is in the negative  $n$ -direction.

$$v = r\dot{\theta}$$

$$a_n = v^2/r = r\dot{\theta}^2 = v\dot{\theta}$$

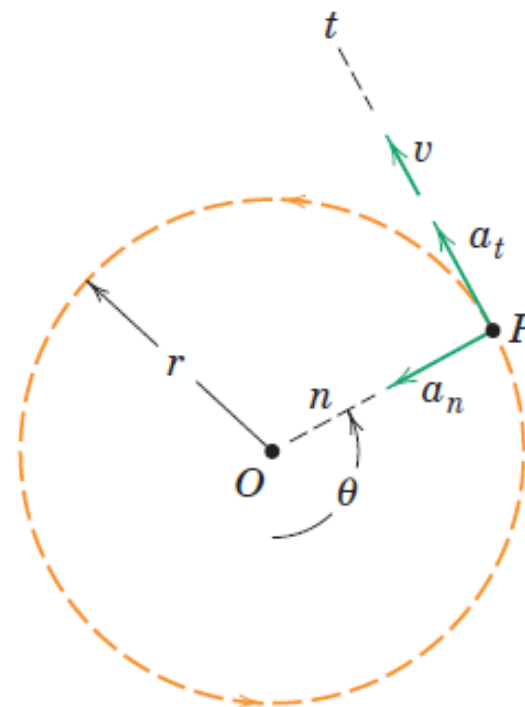
$$a_t = \dot{v} = r\ddot{\theta}$$

$$v_r = 0$$

$$v_\theta = r\dot{\theta}$$

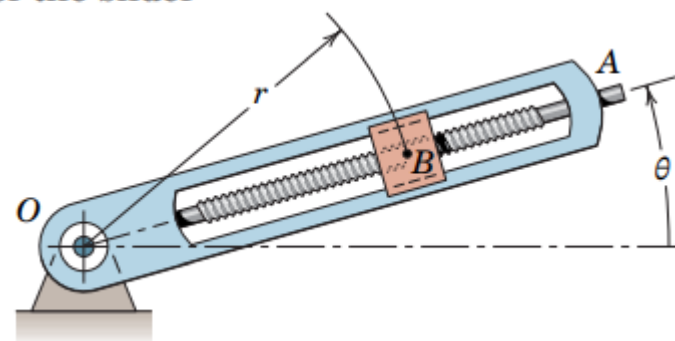
$$a_r = -r\dot{\theta}^2$$

$$a_\theta = r\ddot{\theta}$$



**SAMPLE PROBLEM 2/9**

Rotation of the radially slotted arm is governed by  $\theta = 0.2t + 0.02t^3$ , where  $\theta$  is in radians and  $t$  is in seconds. Simultaneously, the power screw in the arm engages the slider  $B$  and controls its distance from  $O$  according to  $r = 0.2 + 0.04t^2$ , where  $r$  is in meters and  $t$  is in seconds. Calculate the magnitudes of the velocity and acceleration of the slider for the instant when  $t = 3$  s.



$$r = 0.2 + 0.04t^2 \quad r_3 = 0.2 + 0.04(3^2) = 0.56 \text{ m}$$

$$\dot{r} = 0.08t \quad \dot{r}_3 = 0.08(3) = 0.24 \text{ m/s}$$

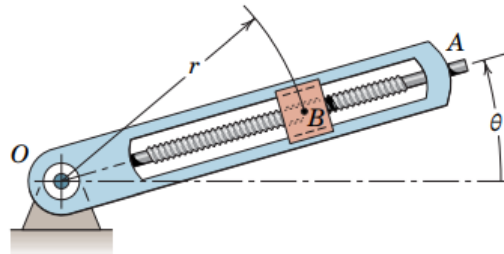
$$\ddot{r} = 0.08 \quad \ddot{r}_3 = 0.08 \text{ m/s}^2$$

$$[v_\theta = r\dot{\theta}] \quad v_r = 0.24 \text{ m/s}$$

$$[v_\theta = r\dot{\theta}] \quad v_\theta = 0.56(0.74) = 0.414 \text{ m/s}$$

$$[v = \sqrt{v_r^2 + v_\theta^2}] \quad v = \sqrt{(0.24)^2 + (0.414)^2} = 0.479 \text{ m/s}$$

## SAMPLE PROBLEM 2/9

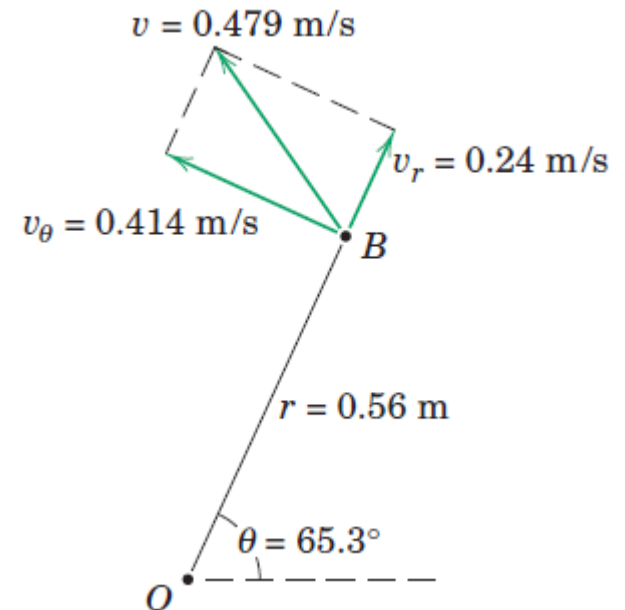


$$\theta = 0.2t + 0.02t^3 \quad \theta_3 = 0.2(3) + 0.02(3^3) = 1.14 \text{ rad}$$

$$\text{or } \theta_3 = 1.14(180/\pi) = 65.3^\circ$$

$$\dot{\theta} = 0.2 + 0.06t^2 \quad \dot{\theta}_3 = 0.2 + 0.06(3^2) = 0.74 \text{ rad/s}$$

$$\ddot{\theta} = 0.12t \quad \ddot{\theta}_3 = 0.12(3) = 0.36 \text{ rad/s}^2$$



## SAMPLE PROBLEM 2/9

$$[a_r = \ddot{r} - r\dot{\theta}^2]$$

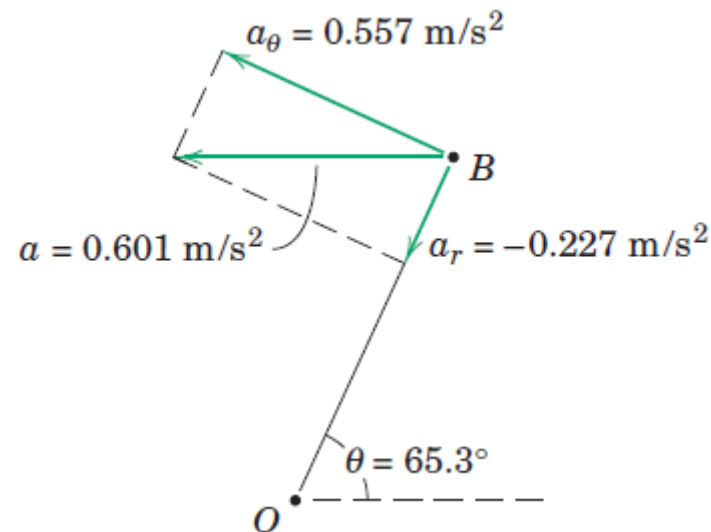
$$a_r = 0.08 - 0.56(0.74)^2 = -0.227 \text{ m/s}^2$$

$$[a_\theta = r\ddot{\theta} + 2\dot{r}\dot{\theta}]$$

$$a_\theta = 0.56(0.36) + 2(0.24)(0.74) = 0.557 \text{ m/s}^2$$

$$[a = \sqrt{a_r^2 + a_\theta^2}]$$

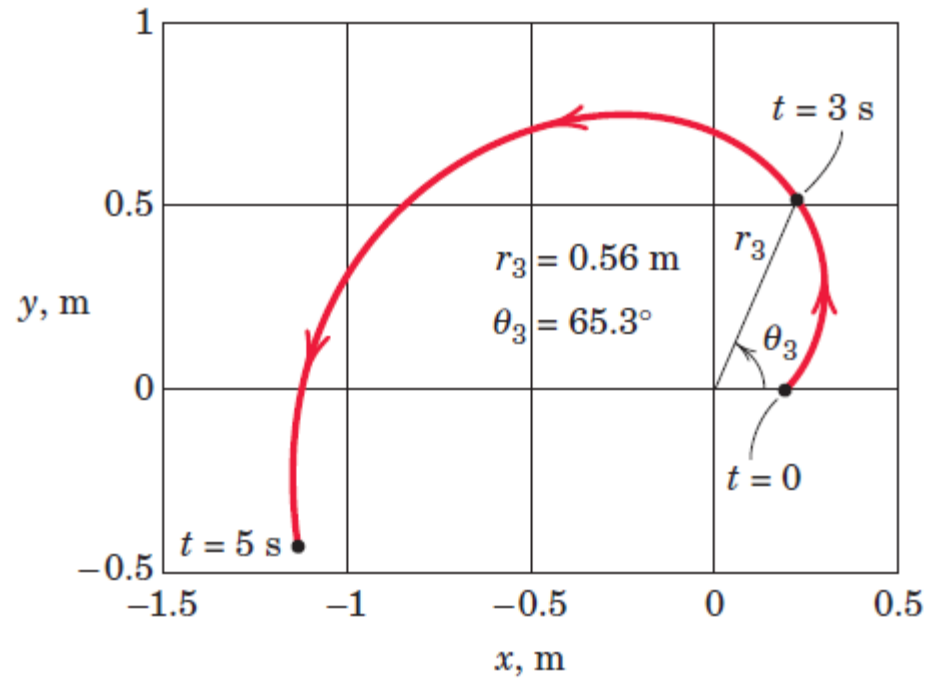
$$a = \sqrt{(-0.227)^2 + (0.557)^2} = 0.601 \text{ m/s}^2$$



## SAMPLE PROBLEM 2/9

- ❖ Conversion from polar to rectangular coordinates

$$x = r \cos \theta \quad y = r \sin \theta$$



2/8

## Relative Motion (Translating Axes)

- ❑ It is not always possible or convenient, however, to use a fixed set of axes to describe or to measure motion.
- ❑ In addition, there are many engineering problems for which the analysis of motion is simplified by using measurements made with respect to a moving reference system.
- ❑ These measurements, when combined with the absolute motion of the moving coordinate system, enable us to determine the absolute motion in question.
- ❑ This approach is called a relative-motion analysis.



2/8

## Relative Motion (Translating Axes)

### Choice of Coordinate System

- ❖ The motion of the moving coordinate system is specified with respect to a fixed coordinate system.

### Vector Representation

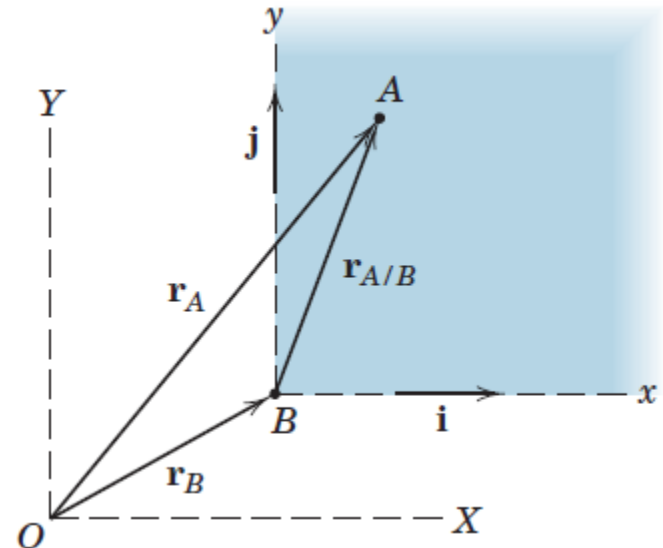
$$\mathbf{r}_A = \mathbf{r}_B + \mathbf{r}_{A/B}$$

$$\dot{\mathbf{r}}_A = \dot{\mathbf{r}}_B + \dot{\mathbf{r}}_{A/B} \quad \text{or}$$

$$\mathbf{v}_A = \mathbf{v}_B + \mathbf{v}_{A/B}$$

$$\ddot{\mathbf{r}}_A = \ddot{\mathbf{r}}_B + \ddot{\mathbf{r}}_{A/B} \quad \text{or}$$

$$\mathbf{a}_A = \mathbf{a}_B + \mathbf{a}_{A/B}$$



2/9

## Constrained Motion of Connected Particles

- Sometimes the motions of particles are interrelated because of the constraints imposed by interconnecting members.

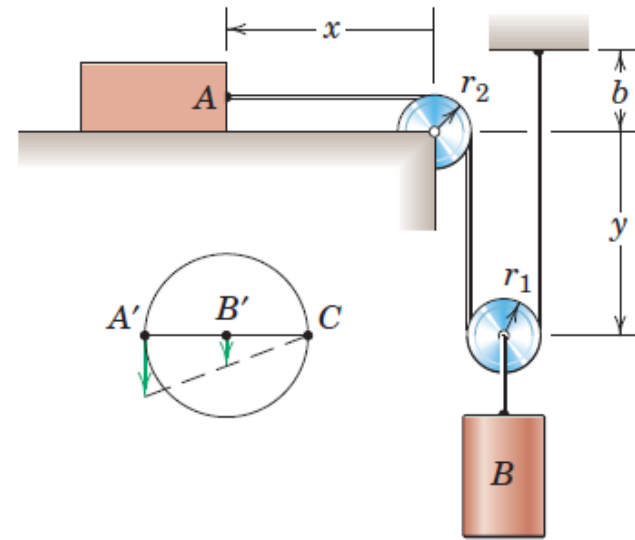
### One Degree of Freedom

- ❖ One degree of freedom: only one variable, either  $x$  or  $y$ , is needed to specify the positions of all parts of the system.

$$L = x + \frac{\pi r_2}{2} + 2y + \pi r_1 + b$$

$$0 = \dot{x} + 2\dot{y} \quad \text{or} \quad 0 = v_A + 2v_B$$

$$0 = \ddot{x} + 2\ddot{y} \quad \text{or} \quad 0 = a_A + 2a_B$$





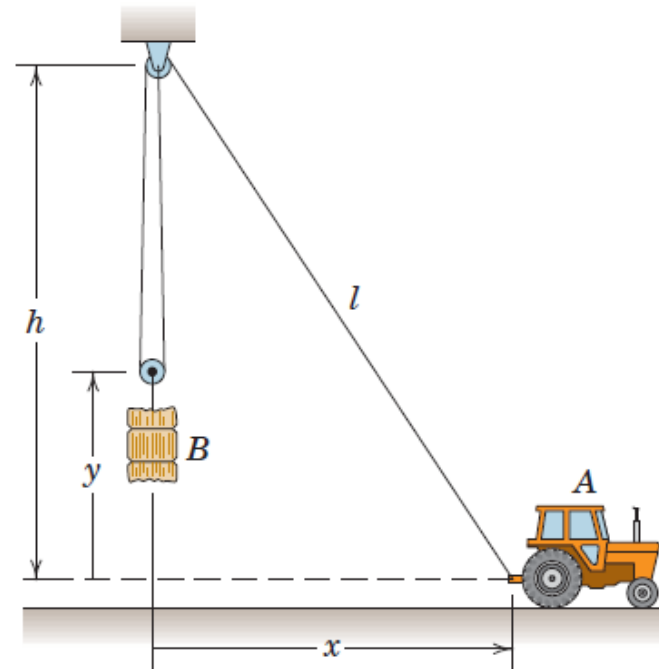
**SAMPLE PROBLEM 2/16**

The tractor  $A$  is used to hoist the bale  $B$  with the pulley arrangement shown. If  $A$  has a forward velocity  $v_A$ , determine an expression for the upward velocity  $v_B$  of the bale in terms of  $x$ .

$$L = 2(h - y) + l = 2(h - y) + \sqrt{h^2 + x^2}$$

$$0 = -2\dot{y} + \frac{x\dot{x}}{\sqrt{h^2 + x^2}}$$

$$v_B = \frac{1}{2} \frac{xv_A}{\sqrt{h^2 + x^2}}$$



❑ CONTENTS:

- ❖ Chapter 1: Introduction to Dynamics
- ❖ Chapter 2: Kinematics of Particles
- ➔ ❖ Chapter 3: **Kinetics of Particles**
- ❖ Chapter 4: Kinetics of Systems of Particles
- ❖ Chapter 5: Plane Kinetics of Rigid Bodies
- ❖ Chapter 6: Plane Kinematics of Rigid Bodies

3/1

## Introduction

- According to Newton's second law, a particle will accelerate when it is subjected to unbalanced forces.
- Kinetics is the study of the relations between unbalanced forces and the resulting changes in motion.
- We combine our knowledge of the properties of forces, which we developed in statics, and the kinematics of particle motion, and solve engineering problems involving force, mass, and motion.



3/1

## Introduction

- General approaches to the solution of kinetics problems:
  - ❖ (A) Direct application of Newton's second law  
(called the force-mass-acceleration method)
  - ❖ (B) Use of work and energy principles
  - ❖ (C) Solution by impulse and momentum methods.



## SECTION A Force, Mass, and Acceleration

3/2

### Newton's Second Law

- The ratios of applied force to corresponding acceleration all equal the same number, provided the units used for measurement are not changed in the experiments.

$$\frac{F_1}{a_1} = \frac{F_2}{a_2} = \dots = \frac{F}{a} = C, \quad \text{a constant}$$

$$\Sigma \mathbf{F} = m \mathbf{a}$$

- We conclude that the constant  $C$  is a measure of some invariable property of the particle. This property is the inertia of the particle, which is its resistance to rate of change of velocity.



3/3

## Equation of Motion and Solution of Problems

### Free-Body Diagram

- The only reliable way to account accurately and consistently for every force is to isolate the particle under consideration from all contacting and influencing bodies and replace the bodies removed by the forces they exert on the particle isolated.
- The resulting free body diagram is the means by which every force, known and unknown, which acts on the particle is represented and thus accounted for.
- Only after this vital step has been completed should you write the appropriate equation or equations of motion.



### 3/4 Rectilinear Motion

- If we choose the x-direction, for example, as the direction of the rectilinear motion of a particle:

$$\Sigma F_x = ma_x$$

$$\Sigma F_y = 0$$

$$\Sigma F_z = 0$$

- For cases where we are not free to choose a coordinate direction along the motion:

$$\Sigma F_x = ma_x$$

$$\Sigma F_y = ma_y$$

$$\Sigma F_z = ma_z$$



**3/4** Rectilinear Motion

- Acceleration and resultant force:

$$\mathbf{a} = a_x \mathbf{i} + a_y \mathbf{j} + a_z \mathbf{k}$$

$$a = \sqrt{a_x^2 + a_y^2 + a_z^2}$$

$$\Sigma \mathbf{F} = \Sigma F_x \mathbf{i} + \Sigma F_y \mathbf{j} + \Sigma F_z \mathbf{k}$$

$$|\Sigma \mathbf{F}| = \sqrt{(\Sigma F_x)^2 + (\Sigma F_y)^2 + (\Sigma F_z)^2}$$





### SAMPLE PROBLEM 3/3

The 250-lb concrete block  $A$  is released from rest in the position shown and pulls the 400-lb log up the  $30^\circ$  ramp. If the coefficient of kinetic friction between the log and the ramp is 0.5, determine the velocity of the block as it hits the ground at  $B$ .

$$L = 2s_C + s_A + \text{constant}$$

$$0 = 2a_C + a_A$$

$$[\Sigma F_y = 0]$$

$$N - 400 \cos 30^\circ = 0 \quad N = 346 \text{ lb}$$

$$[\Sigma F_x = ma_x]$$

$$0.5(346) - 2T + 400 \sin 30^\circ = \frac{400}{32.2} a_C$$

$$[+\downarrow \Sigma F = ma]$$

$$250 - T = \frac{250}{32.2} a_A$$

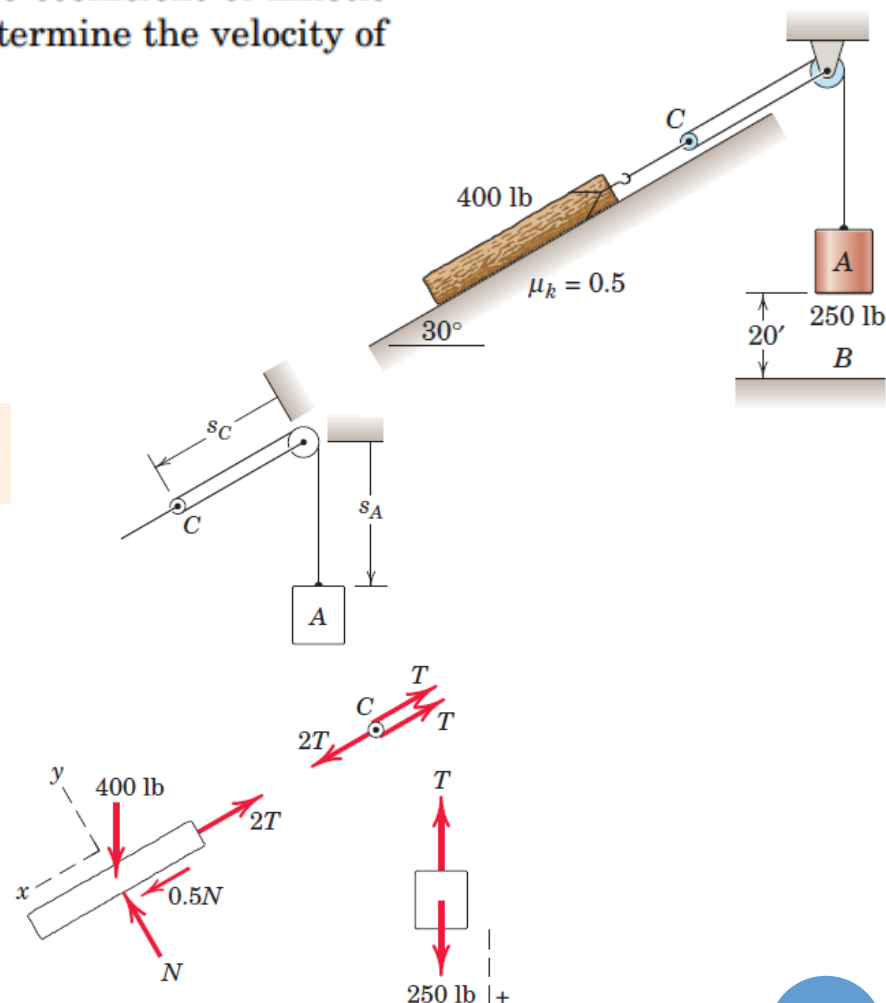
$$a_A = 5.83 \text{ ft/sec}^2$$

$$a_C = -2.92 \text{ ft/sec}^2$$

$$T = 205 \text{ lb}$$

$$[v^2 = 2ax]$$

$$v_A = \sqrt{2(5.83)(20)} = 15.27 \text{ ft/sec}$$



## 3/5 Curvilinear Motion

- Rectangular coordinates:

$$\begin{array}{l} \Sigma F_x = ma_x \\ \Sigma F_y = ma_y \end{array} \quad a_x = \ddot{x} \quad \text{and} \quad a_y = \ddot{y}$$

- Normal and tangential coordinates:

$$\begin{array}{l} \Sigma F_n = ma_n \\ \Sigma F_t = ma_t \end{array} \quad a_n = \rho \dot{\beta}^2 = v^2 / \rho = v \dot{\beta}, \quad a_t = \dot{v}, \quad \text{and} \quad v = \rho \dot{\beta}$$

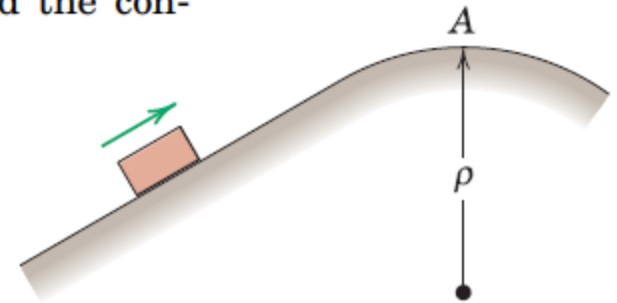
- Polar coordinates:

$$\begin{array}{l} \Sigma F_r = ma_r \\ \Sigma F_\theta = ma_\theta \end{array} \quad a_r = \ddot{r} - r\dot{\theta}^2 \quad \text{and} \quad a_\theta = r\ddot{\theta} + 2\dot{r}\dot{\theta}$$



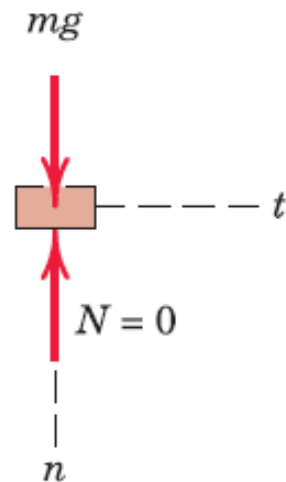
### SAMPLE PROBLEM 3/6

Determine the maximum speed  $v$  which the sliding block may have as it passes the topmost point  $A$  without losing contact with the lower surface. Assume a slightly loose fit between the slider and the constraint surfaces.



$$[\Sigma F_n = ma_n]$$

$$mg = m \frac{v^2}{\rho} \quad v = \sqrt{g\rho}$$



## SECTION B Work and Energy

---

3/6

### Work and Kinetic Energy

- There are two general classes of problems:
  - ❖ (1) Integration of the forces with respect to the displacement of the particle
  - ❖ (2) Integration of the forces with respect to the time they are applied.
  
- Integration with respect to displacement leads to the equations of work and energy.



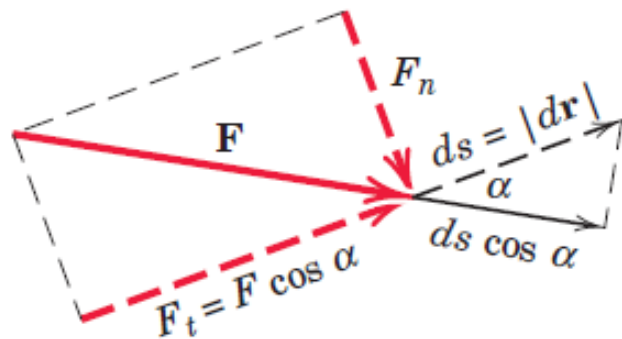
## 3/6 Work and Kinetic Energy

### Definition of Work

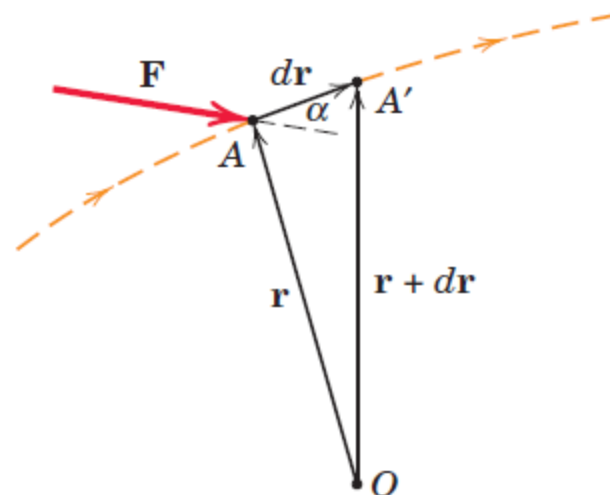
- The work done by the force  $F$  during the displacement  $dr$ :

$$\rightarrow dU = F ds \cos \alpha$$

$$F_t = F \cos \alpha \rightarrow dU = F_t ds$$



$$dU = \mathbf{F} \cdot d\mathbf{r}$$

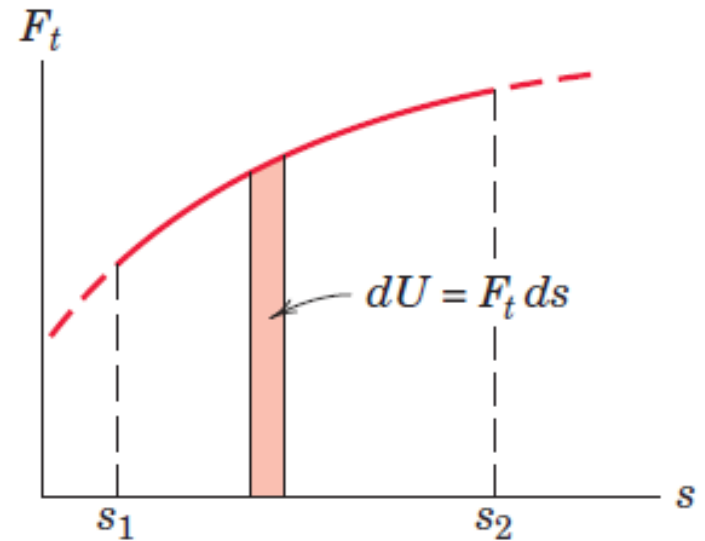


## 3/6 Work and Kinetic Energy

### Calculation of Work

$$U = \int_1^2 \mathbf{F} \cdot d\mathbf{r} = \int_1^2 (F_x dx + F_y dy + F_z dz)$$

$$U = \int_{s_1}^{s_2} F_t ds$$

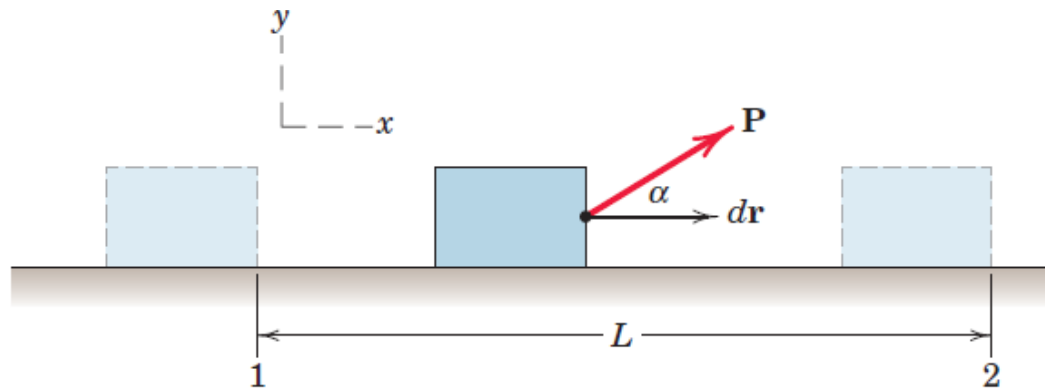


## 3/6 Work and Kinetic Energy

### Examples of Work

- Work Associated with a Constant External Force

$$\begin{aligned}
 U_{1-2} &= \int_1^2 \mathbf{F} \cdot d\mathbf{r} = \int_1^2 [(P \cos \alpha)\mathbf{i} + (P \sin \alpha)\mathbf{j}] \cdot dx \mathbf{i} \\
 &= \int_{x_1}^{x_2} P \cos \alpha dx = P \cos \alpha (x_2 - x_1) = PL \cos \alpha
 \end{aligned}$$

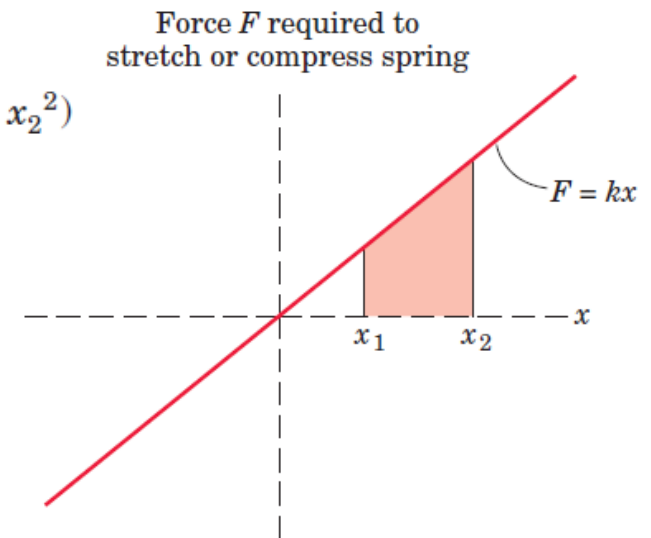
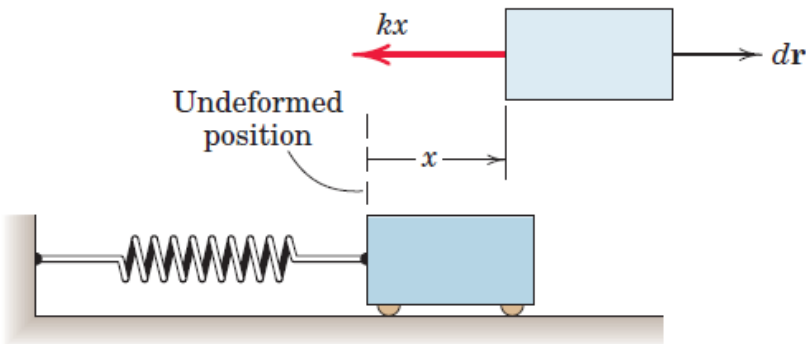


## 3/6 Work and Kinetic Energy

### Examples of Work

- Work Associated with a Spring Force

$$U_{1-2} = \int_1^2 \mathbf{F} \cdot d\mathbf{r} = \int_1^2 (-kx\mathbf{i}) \cdot dx\mathbf{i} = -\int_{x_1}^{x_2} kx dx = \frac{1}{2}k(x_1^2 - x_2^2)$$



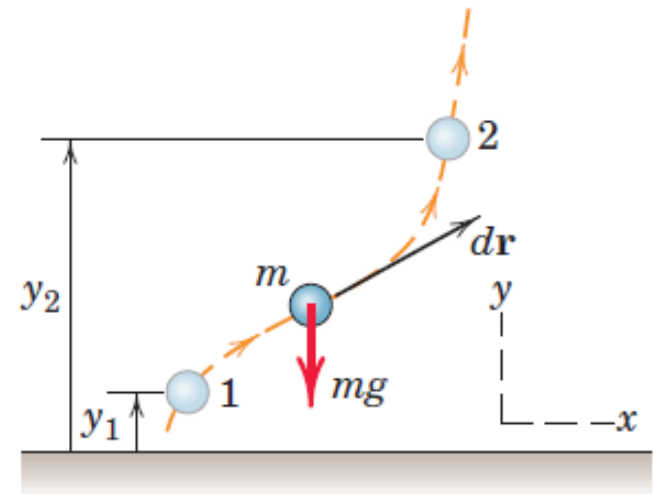


## 3/6 Work and Kinetic Energy

### Examples of Work

- Work Associated with Weight

$$\begin{aligned}
 U_{1-2} &= \int_1^2 \mathbf{F} \cdot d\mathbf{r} = \int_1^2 (-mg\mathbf{j}) \cdot (dx\mathbf{i} + dy\mathbf{j}) \\
 &= -mg \int_{y_1}^{y_2} dy = -mg(y_2 - y_1)
 \end{aligned}$$



## 3/6 Work and Kinetic Energy

### Principle of Work and Kinetic Energy

- The *kinetic energy*  $T$  of the particle:

$$T = \frac{1}{2}mv^2$$

- The *work-energy equation* for a particle:

$$T_1 + U_{1-2} = T_2$$

- ❖ The equation states that the total work done by all forces acting on a particle as it moves from point 1 to point 2 equals the corresponding change in kinetic energy of the particle.



3/6

## Work and Kinetic Energy

### Power

- The capacity of a machine is rated by its power, which is defined as the time rate of doing work

$$P = dU/dt = \mathbf{F} \cdot d\mathbf{r}/dt$$



$$P = \mathbf{F} \cdot \mathbf{v}$$

$$1 \text{ W} = 1 \text{ J/s}$$

$$1 \text{ hp} = 550 \text{ ft-lb/sec} = 33,000 \text{ ft-lb/min}$$

$$1 \text{ hp} = 746 \text{ W} = 0.746 \text{ kW}$$



3/6

## Work and Kinetic Energy

### Efficiency

- ❖ The ratio of the work done by a machine to the work done on the machine during the same time interval is called the mechanical efficiency  $e_m$  of the machine.

$$e_m = \frac{P_{\text{output}}}{P_{\text{input}}}$$

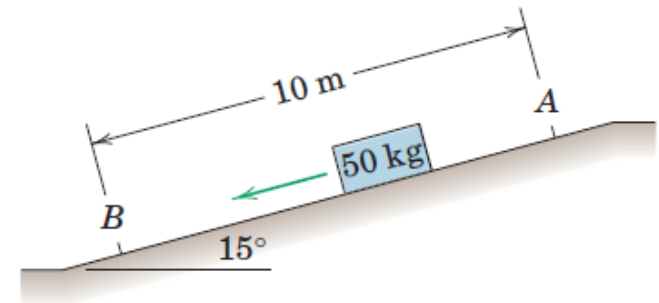
- ❖ In addition to energy loss by mechanical friction, there may also be electrical and thermal energy loss, in which case, the electrical efficiency  $e_e$  and thermal efficiency  $e_t$  are also involved. The overall efficiency  $e$  in such instances is:

$$e = e_m e_e e_t$$



**SAMPLE PROBLEM 3/11**

Calculate the velocity  $v$  of the 50-kg crate when it reaches the bottom of the chute at  $B$  if it is given an initial velocity of 4 m/s down the chute at  $A$ . The coefficient of kinetic friction is 0.30.

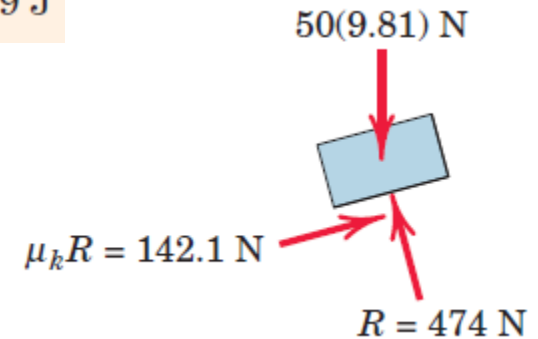


$$[U = Fs] \quad U_{1-2} = 50(9.81)(10 \sin 15^\circ) - 142.1(10) = -151.9 \text{ J}$$

$$[T_1 + U_{1-2} = T_2] \quad \frac{1}{2}mv_1^2 + U_{1-2} = \frac{1}{2}mv_2^2$$

$$\frac{1}{2}(50)(4)^2 - 151.9 = \frac{1}{2}(50)v_2^2$$

$$v_2 = 3.15 \text{ m/s}$$



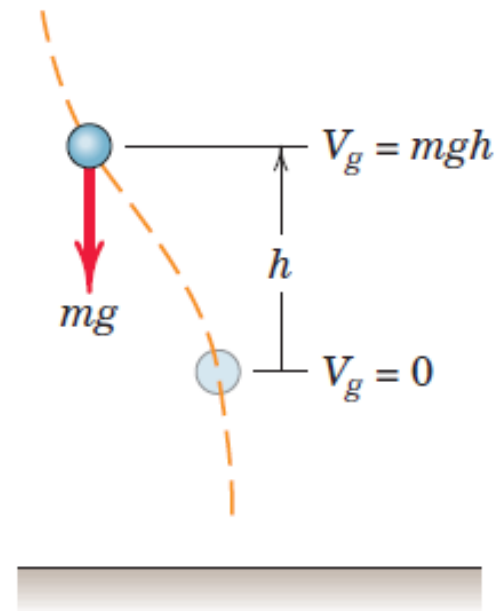
## 3/7 Potential Energy

### Gravitational Potential Energy

- The *gravitational potential energy*  $V_g$  of the particle is defined as the work  $mgh$  done against the gravitational field to elevate the particle a distance  $h$  above some arbitrary reference plane.

$$V_g = mgh$$

$$\Delta V_g = mg(h_2 - h_1) = mg\Delta h$$



## 3/7 Potential Energy

### Elastic Potential Energy

- The work which is done on the spring to deform it is stored in the spring and is called its *elastic potential energy*  $V_e$ .

$$V_e = \int_0^x kx \, dx = \frac{1}{2} kx^2$$

$$\Delta V_e = \frac{1}{2} k(x_2^2 - x_1^2)$$



## 3/7 Potential Energy

### Work-Energy Equation

- ❖ Work-energy equation modification to account for the potential-energy terms

$$U'_{1-2} + (-\Delta V_g) + (-\Delta V_e) = \Delta T$$

$$U'_{1-2} = \Delta T + \Delta V$$

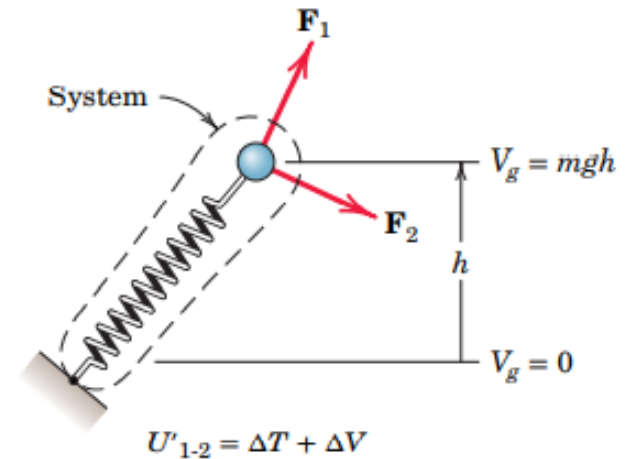
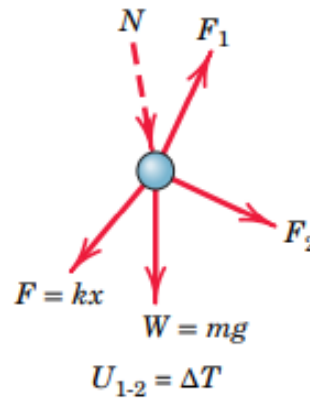
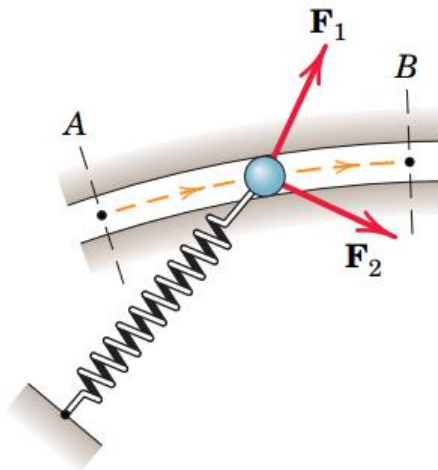
$$T_1 + V_1 + U'_{1-2} = T_2 + V_2$$





## 3/7 Potential Energy

### Work-Energy Equation



3/7

## Potential Energy

### Work-Energy Equation

- ❖ For problems where the only forces are gravitational, elastic, and nonworking constraint forces:

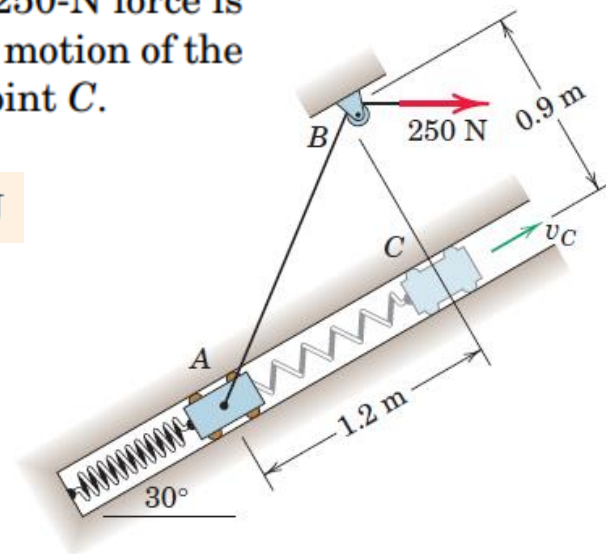
$$T_1 + V_1 = T_2 + V_2 \quad \text{or} \quad E_1 = E_2$$

- ❖  $E=T+V$  is the total mechanical energy of the particle and its attached spring.



### SAMPLE PROBLEM 3/17

The 10-kg slider moves with negligible friction up the inclined guide. The attached spring has a stiffness of 60 N/m and is stretched 0.6 m in position A, where the slider is released from rest. The 250-N force is constant and the pulley offers negligible resistance to the motion of the cord. Calculate the velocity  $v_C$  of the slider as it passes point C.



$$\overline{AB} - \overline{BC} \text{ or } 1.5 - 0.9 = 0.6 \text{ m.}$$

$$U'_{A-C} = 250(0.6) = 150 \text{ J}$$

$$V_A = 0 \quad V_C = mgh = 10(9.81)(1.2 \sin 30^\circ) = 58.9 \text{ J}$$

$$V_A = \frac{1}{2} kx_A^2 = \frac{1}{2} (60)(0.6)^2 = 10.8 \text{ J}$$

$$V_C = \frac{1}{2} kx_B^2 = \frac{1}{2} 60(0.6 + 1.2)^2 = 97.2 \text{ J}$$

$$[T_A + V_A + U'_{A-C} = T_C + V_C] \quad 0 + 0 + 10.8 + 150 = \frac{1}{2}(10)v_C^2 + 58.9 + 97.2$$

$$v_C = 0.974 \text{ m/s}$$

# SECTION C Impulse and Momentum

---

3/8

## Introduction

- We can integrate the equation of motion with respect to time rather than displacement.
- This approach leads to the equations of impulse and momentum.
- These equations greatly facilitate the solution of many problems in which the applied forces act during extremely short periods of time (as in impact problems) or over specified intervals of time.



3/9

## Linear Impulse and Linear Momentum

- Linear momentum of the particle:

$$\diamond \mathbf{G} = m\mathbf{v}$$

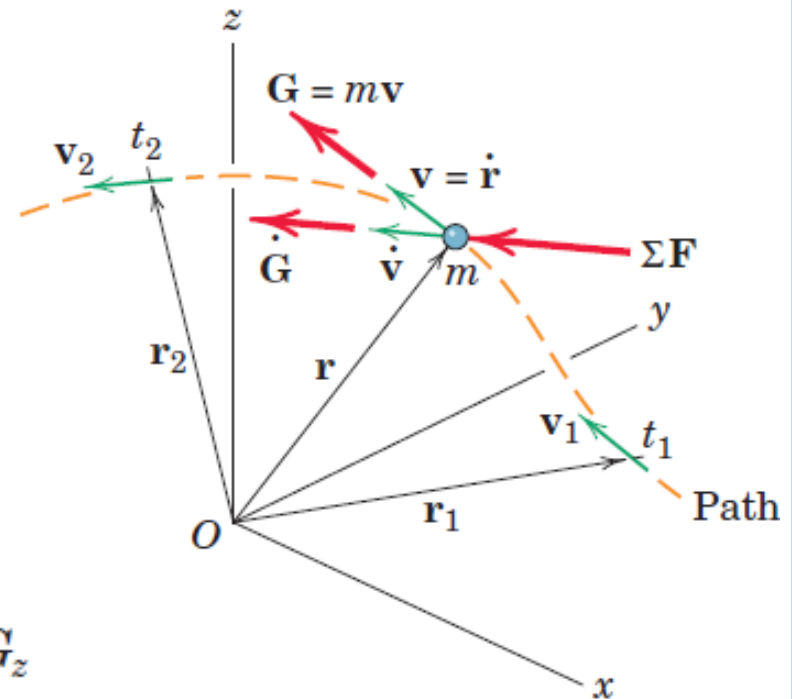
$$\Sigma \mathbf{F} = m\dot{\mathbf{v}} = \frac{d}{dt}(m\mathbf{v})$$

$$\rightarrow \boxed{\Sigma \mathbf{F} = \dot{\mathbf{G}}}$$

$$\Sigma F_x = \dot{G}_x$$

$$\Sigma F_y = \dot{G}_y$$

$$\Sigma F_z = \dot{G}_z$$

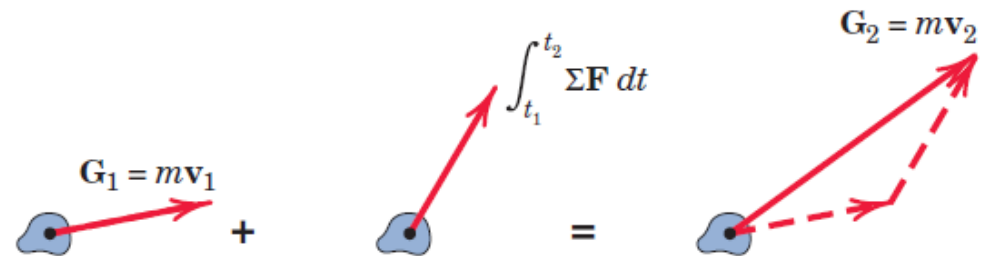


3/9

## Linear Impulse and Linear Momentum

### The Linear Impulse-Momentum Principle

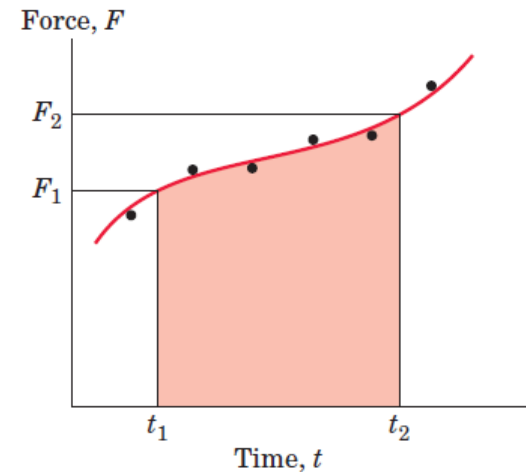
$$\mathbf{G}_1 + \int_{t_1}^{t_2} \Sigma \mathbf{F} dt = \mathbf{G}_2$$



$$m(v_1)_x + \int_{t_1}^{t_2} \Sigma F_x dt = m(v_2)_x$$

$$m(v_1)_y + \int_{t_1}^{t_2} \Sigma F_y dt = m(v_2)_y$$

$$m(v_1)_z + \int_{t_1}^{t_2} \Sigma F_z dt = m(v_2)_z$$



3/9

## Linear Impulse and Linear Momentum

### Conservation of Linear Momentum

- If the resultant force on a particle is zero during an interval of time, its linear momentum  $\mathbf{G}$  remain constant.
- In this case, the linear momentum of the particle is said to be *conserved*.

$$\Delta \mathbf{G} = \mathbf{0} \quad \text{or} \quad \mathbf{G}_1 = \mathbf{G}_2$$

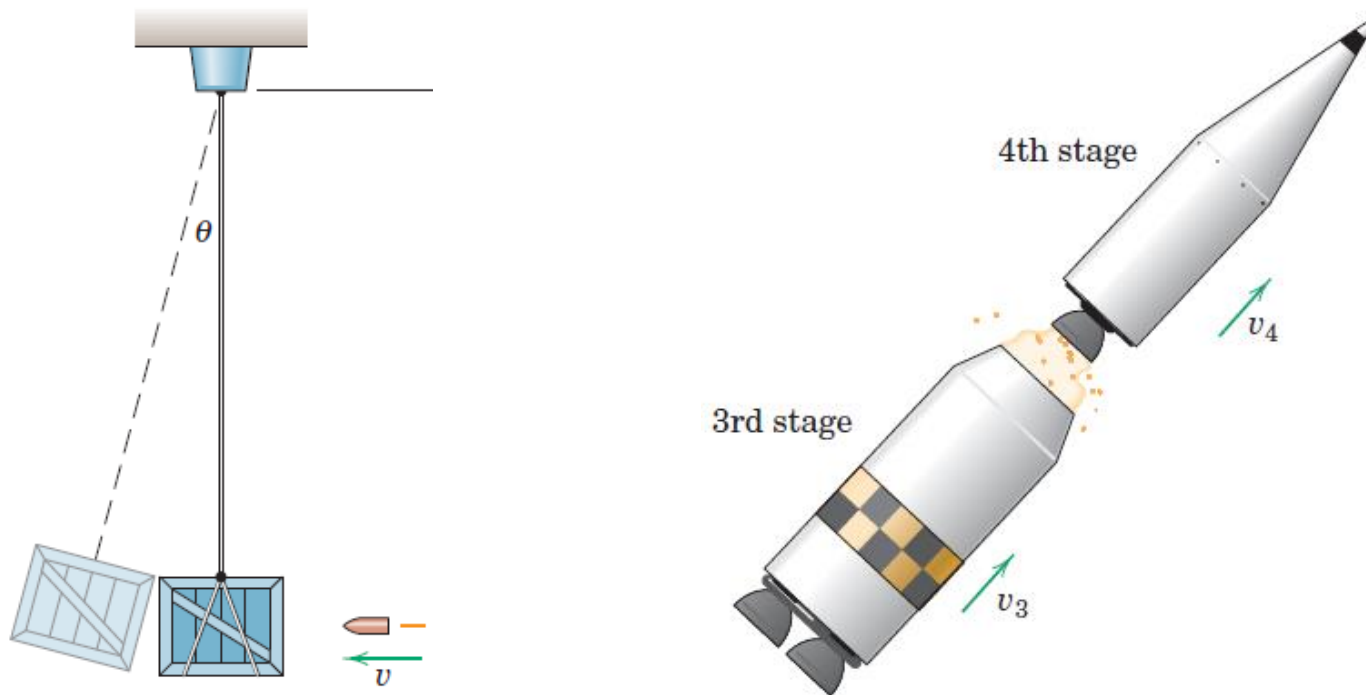
*Principle of conservation of linear momentum*



3/9

# Linear Impulse and Linear Momentum

## Conservation of Linear Momentum



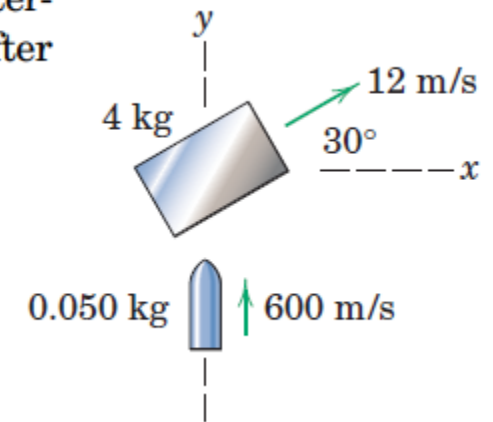


### SAMPLE PROBLEM 3/23

The 50-g bullet traveling at 600 m/s strikes the 4-kg block centrally and is embedded within it. If the block slides on a smooth horizontal plane with a velocity of 12 m/s in the direction shown prior to impact, determine the velocity  $\mathbf{v}_2$  of the block and embedded bullet immediately after impact.

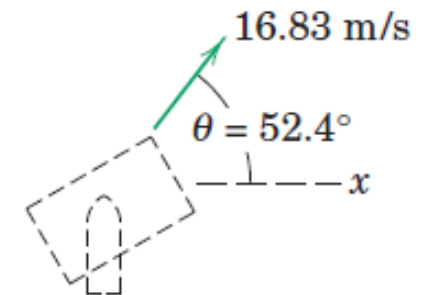
$$[\mathbf{G}_1 = \mathbf{G}_2] \quad 0.050(600\mathbf{j}) + 4(12)(\cos 30^\circ\mathbf{i} + \sin 30^\circ\mathbf{j}) = (4 + 0.050)\mathbf{v}_2$$

$$\mathbf{v}_2 = 10.26\mathbf{i} + 13.33\mathbf{j} \text{ m/s}$$



$$[v = \sqrt{v_x^2 + v_y^2}] \quad v_2 = \sqrt{(10.26)^2 + (13.33)^2} = 16.83 \text{ m/s}$$

$$[\tan \theta = v_y/v_x] \quad \tan \theta = \frac{13.33}{10.26} = 1.299 \quad \theta = 52.4^\circ$$

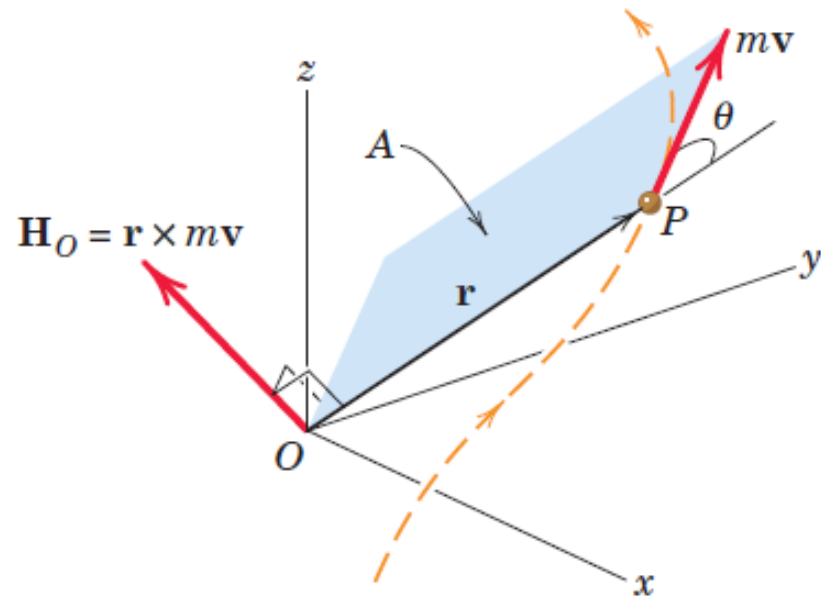


3/10

## Angular Impulse and Angular Momentum

- A particle P of mass  $m$  moving along a curve in space.
- The velocity of the particle is  $v$ , and its linear momentum is  $G = mv$ .
- The moment of the linear momentum vector  $mv$  about the origin  $O$  is defined as the angular momentum  $H_O$  of P about  $O$ .

$$\mathbf{H}_O = \mathbf{r} \times m\mathbf{v}$$



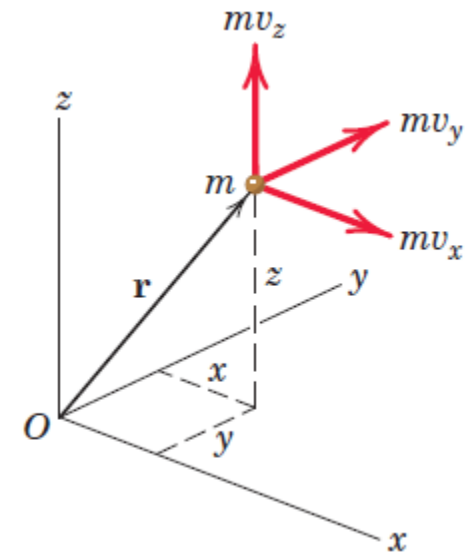
3/10

## Angular Impulse and Angular Momentum

- The scalar components of angular momentum:

$$\mathbf{H}_O = \mathbf{r} \times m\mathbf{v} = m(v_z y - v_y z)\mathbf{i} + m(v_x z - v_z x)\mathbf{j} + m(v_y x - v_x y)\mathbf{k}$$

$$\mathbf{H}_O = m \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ x & y & z \\ v_x & v_y & v_z \end{vmatrix}$$

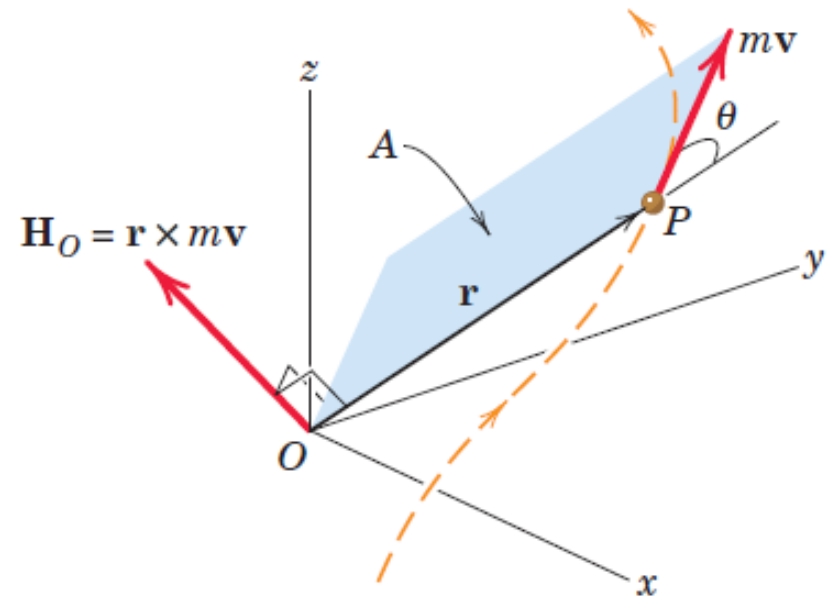
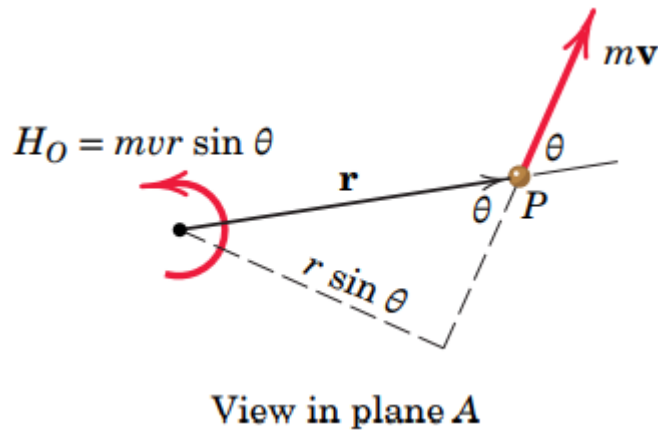


$$H_x = m(v_z y - v_y z) \quad H_y = m(v_x z - v_z x) \quad H_z = m(v_y x - v_x y)$$

3/10

## Angular Impulse and Angular Momentum

- Two- dimensional representation:



3/10

## Angular Impulse and Angular Momentum

### Rate of Change of Angular Momentum

- The moment of the forces acting on the particle P to its angular momentum relation:

$$\Sigma \mathbf{M}_O = \mathbf{r} \times \Sigma \mathbf{F} = \mathbf{r} \times m \dot{\mathbf{v}}$$

- $H_O$  Differentiation leads to:

$$\dot{H}_O = \dot{\mathbf{r}} \times m \mathbf{v} + \mathbf{r} \times m \dot{\mathbf{v}} = \mathbf{v} \times m \mathbf{v} + \mathbf{r} \times m \dot{\mathbf{v}}$$

- So:

$$\Sigma \mathbf{M}_O = \dot{H}_O$$

$$\Sigma M_{O_x} = \dot{H}_{O_x}$$

$$\Sigma M_{O_y} = \dot{H}_{O_y}$$

$$\Sigma M_{O_z} = \dot{H}_{O_z}$$

- ❖ The moment about the fixed point O of all forces acting on m equals the time rate of change of angular momentum of m about O.



3/10

## Angular Impulse and Angular Momentum

### The Angular Impulse-Momentum Principle

- The total angular impulse on  $m$  about the fixed point  $O$  equals the corresponding change in angular momentum of  $m$  about  $O$ .

$$\int_{t_1}^{t_2} \Sigma \mathbf{M}_O dt = (\mathbf{H}_O)_2 - (\mathbf{H}_O)_1 = \Delta \mathbf{H}_O$$

$$(\mathbf{H}_O)_1 + \int_{t_1}^{t_2} \Sigma \mathbf{M}_O dt = (\mathbf{H}_O)_2$$

$$(H_{O_x})_1 + \int_{t_1}^{t_2} \Sigma M_{O_x} dt = (H_{O_x})_2$$

$$m(v_z y - v_y z)_1 + \int_{t_1}^{t_2} \Sigma M_{O_x} dt = m(v_z y - v_y z)_2$$



3/10

## Angular Impulse and Angular Momentum

### Conservation of Angular Momentum

- If the resultant moment about a fixed point O of all forces acting on a particle is zero during an interval of time, its angular momentum  $H_O$  about that point remain constant.
- In this case, the angular momentum of the particle is said to be conserved.

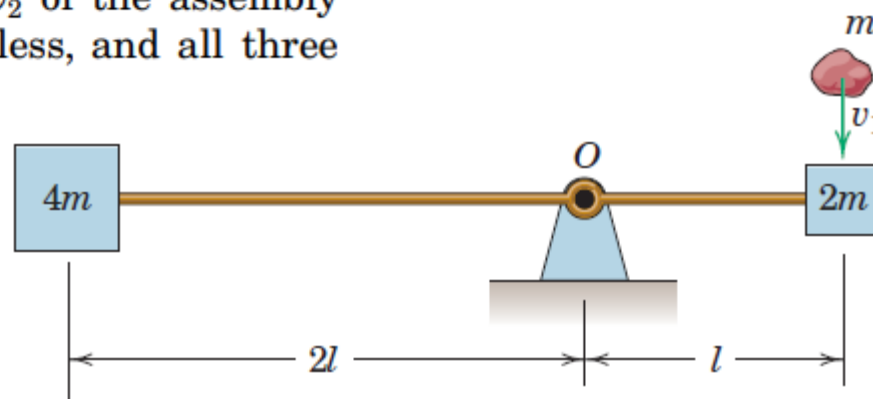
$$\Delta \mathbf{H}_O = \mathbf{0} \quad \text{or} \quad (\mathbf{H}_O)_1 = (\mathbf{H}_O)_2$$

*Principle of conservation of angular momentum*



**SAMPLE PROBLEM 3/26**

The assembly of the light rod and two end masses is at rest when it is struck by the falling wad of putty traveling with speed  $v_1$  as shown. The putty adheres to and travels with the right-hand end mass. Determine the angular velocity  $\dot{\theta}_2$  of the assembly just after impact. The pivot at  $O$  is frictionless, and all three masses may be assumed to be particles.



$$(H_O)_1 = (H_O)_2$$

$$mv_1l = (m + 2m)(l\dot{\theta}_2)l + 4m(2l\dot{\theta}_2)2l$$

$$\dot{\theta}_2 = \frac{v_1}{19l} \text{ CW}$$



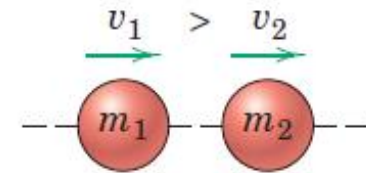
## 3/12 Impact

### Direct Central Impact

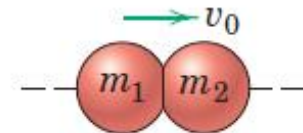
- Collision of two spheres with collinear motion
  - Conservation of linear momentum:

$$m_1 v_1 + m_2 v_2 = m_1 v_1' + m_2 v_2'$$

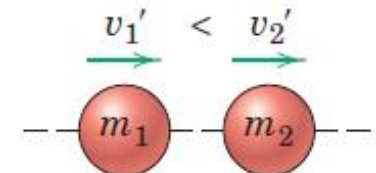
Before impact



Maximum deformation during impact



After impact



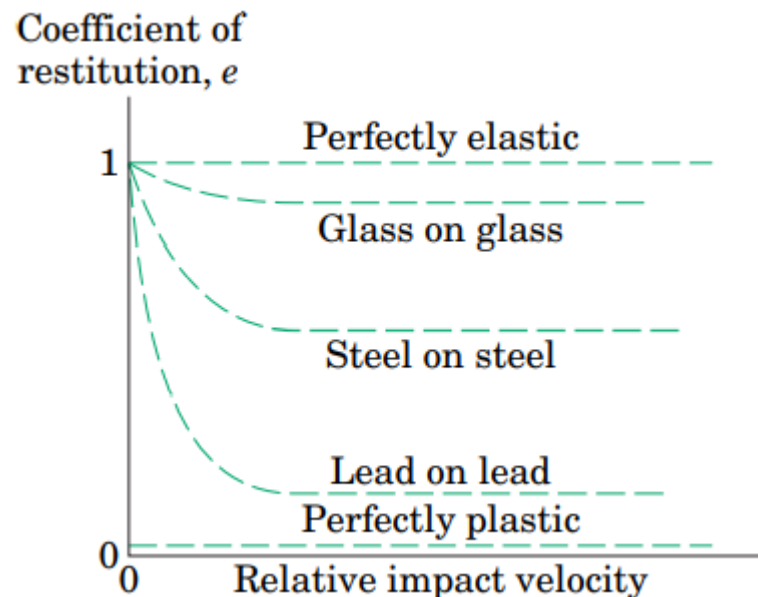
### Coefficient of Restitution

$$e = \frac{v_2' - v_1'}{v_1 - v_2}$$

## 3/12 Impact

### Energy Loss During Impact

- Impact phenomena are almost always accompanied by energy loss, which may be calculated by subtracting the kinetic energy of the system just after impact from that just before impact.



## 3/12 Impact

## Oblique Central Impact

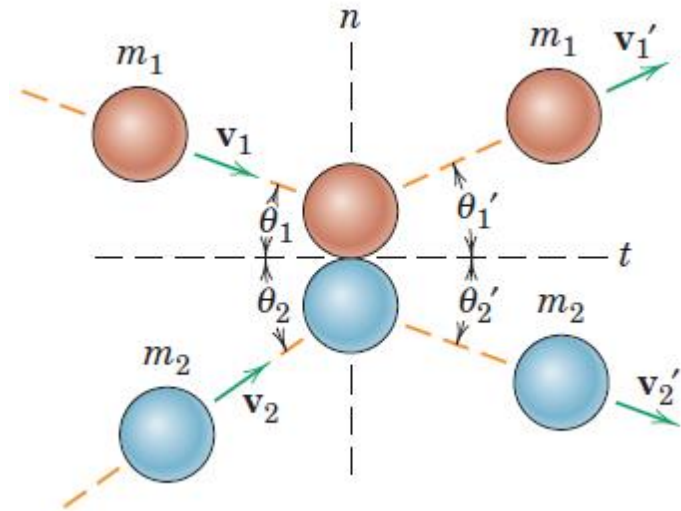
- Tangent and normal directions

$$m_1(v_1)_t = m_1(v_1')_t$$

$$m_2(v_2)_t = m_2(v_2')_t$$

$$m_1(v_1)_n + m_2(v_2)_n = m_1(v_1')_n + m_2(v_2')_n$$

$$e = \frac{(v_2')_n - (v_1')_n}{(v_1)_n - (v_2)_n}$$



### SAMPLE PROBLEM 3/29

A ball is projected onto the heavy plate with a velocity of 50 ft/sec at the  $30^\circ$  angle shown. If the effective coefficient of restitution is 0.5, compute the rebound velocity  $v'$  and its angle  $\theta'$ .

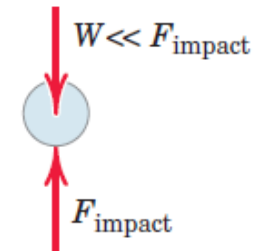
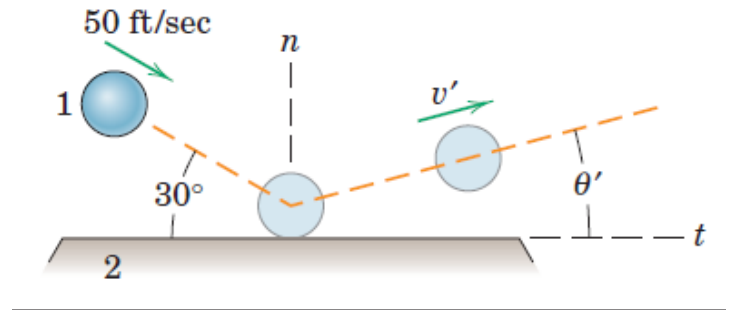
$$e = \frac{(v_2')_n - (v_1')_n}{(v_1)_n - (v_2)_n} \quad 0.5 = \frac{0 - (v_1')_n}{-50 \sin 30^\circ - 0}$$

$$(v_1')_n = 12.5 \text{ ft/sec}$$

$$m(v_1)_t = m(v_1')_t \quad (v_1')_t = (v_1)_t = 50 \cos 30^\circ = 43.3 \text{ ft/sec}$$

$$v' = \sqrt{(v_1')_n^2 + (v_1')_t^2} = \sqrt{12.5^2 + 43.3^2} = 45.1 \text{ ft/sec}$$

$$\theta' = \tan^{-1} \left( \frac{(v_1')_n}{(v_1')_t} \right) = \tan^{-1} \left( \frac{12.5}{43.3} \right) = 16.10^\circ$$



3/14

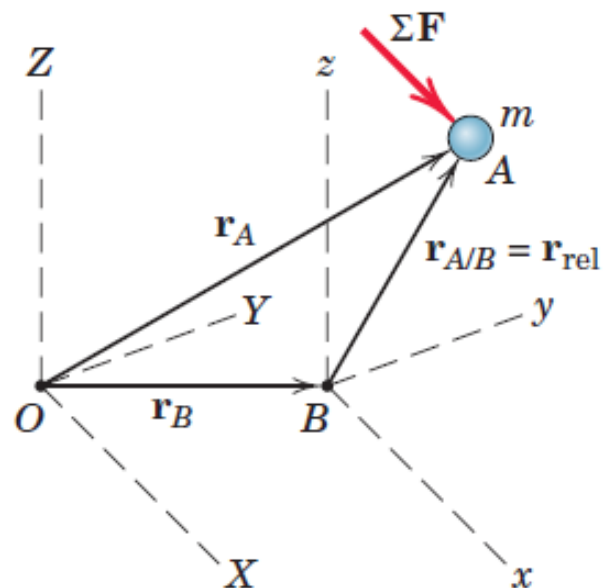
## Relative Motion

## Relative-Motion Equation

- A particle **A** of mass  $m$  whose motion is observed from a set of axes  $x$ - $y$ - $z$  which translate with respect to a fixed reference frame  $X$ - $Y$ - $Z$ .

$$\mathbf{a}_A = \mathbf{a}_B + \mathbf{a}_{\text{rel}}$$

$$\Sigma \mathbf{F} = m(\mathbf{a}_B + \mathbf{a}_{\text{rel}})$$

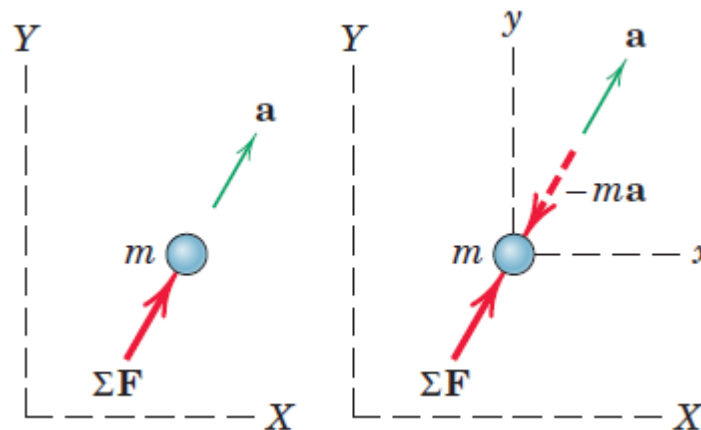


3/14

## Relative Motion

## D'Alembert's Principle

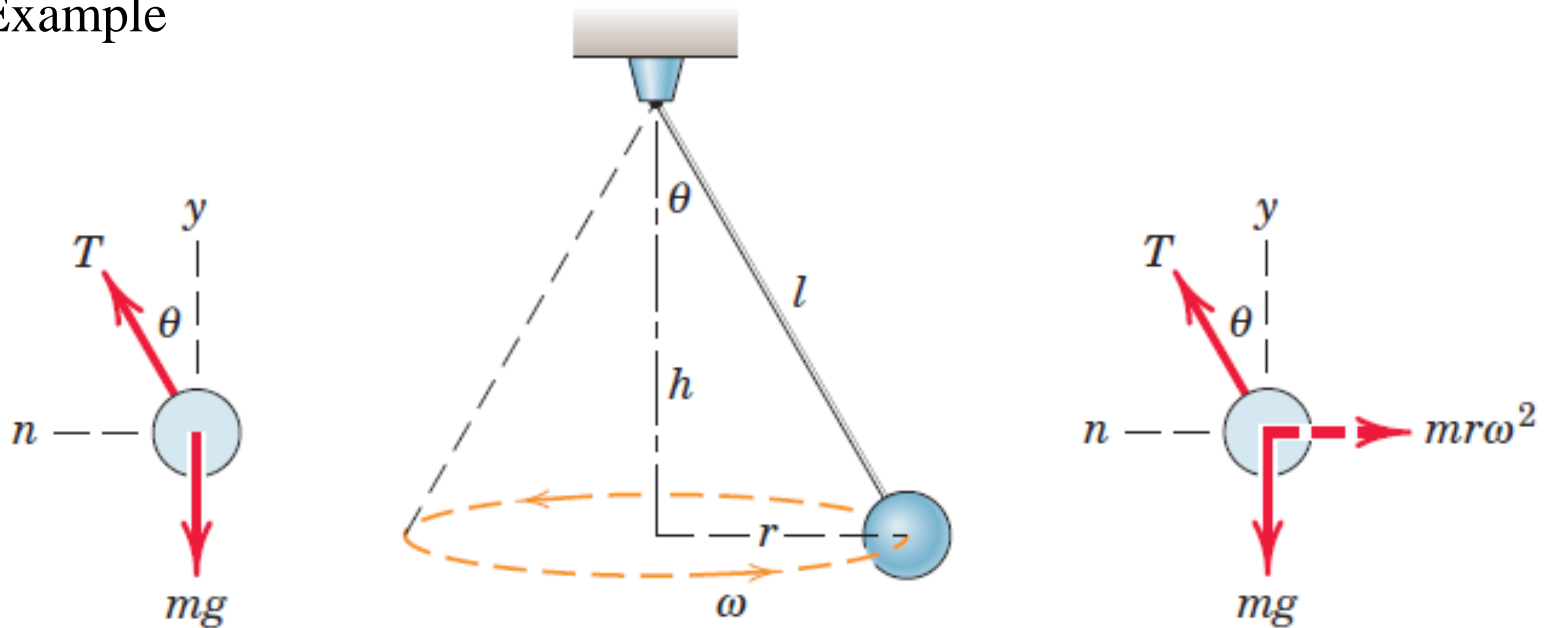
- ❖ The particle acceleration we measure from a fixed set of axes X-Y-Z, is its absolute acceleration  $a$ . In this case the familiar relation  $\Sigma F = ma$  applies.
- ❖ When we observe the particle from a moving system x-y-z attached to the particle, the particle necessarily appears to be at rest or in equilibrium in x-y-z.
- ❖ Thus, the observer who is accelerating with x-y-z concludes that a force  $-ma$  acts on the particle to balance  $\Sigma F$ .



## 3/14 Relative Motion

## D'Alembert's Principle

## □ Example



❑ CONTENTS:

- ❖ Chapter 1: Introduction to Dynamics
- ❖ Chapter 2: Kinematics of Particles
- ❖ Chapter 3: Kinetics of Particles
- ➔ ❖ Chapter 4: **Kinetics of Systems of Particles**
- ❖ Chapter 5: Plane Kinematics of Rigid Bodies
- ❖ Chapter 6: Plane Kinetics of Rigid Bodies



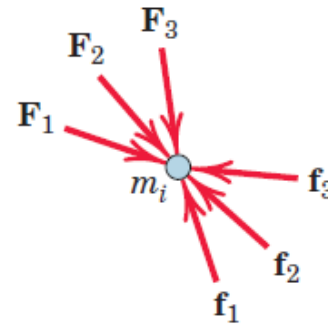
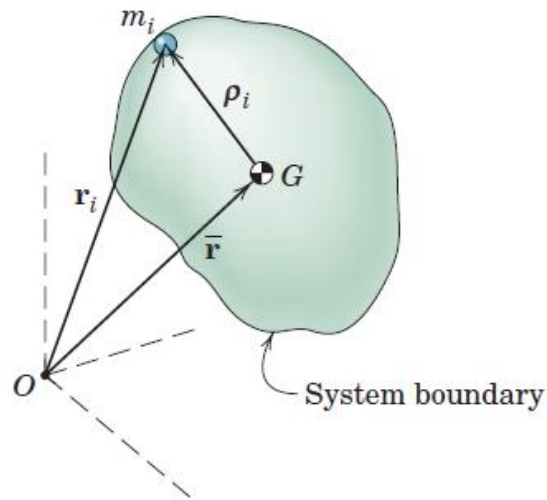
## 4/1 Introduction

- In the previous two chapters, we have applied the principles of dynamics to the motion of a particle.
- Our next major step in the development of dynamics is to extend these principles, which we applied to a single particle, to describe the motion of a general system of particles.
- Recall that a rigid body is a solid system of particles wherein the distances between particles remain essentially unchanged.



## 4/2 Generalized Newton's Second Law

- Considering  $n$  mass particles bounded by a closed surface in space



- Forces  $F_1, F_2, F_3, \dots$  acting on  $m_i$  from sources external to the envelop
- Forces  $f_1, f_2, f_3, \dots$  acting on  $m_i$  from sources internal to the system boundary

## 4/2 Generalized Newton's Second Law

- The center of mass  $G$  of the isolated system of particles

$$m\bar{\mathbf{r}} = \sum m_i \mathbf{r}_i \qquad m = \sum m_i$$

- Newton's second law when applied to  $m_i$  gives:

$$\mathbf{F}_1 + \mathbf{F}_2 + \mathbf{F}_3 + \cdots + \mathbf{f}_1 + \mathbf{f}_2 + \mathbf{f}_3 + \cdots = m_i \ddot{\mathbf{r}}_i$$

$$\rightarrow \Sigma \mathbf{F} + \Sigma \mathbf{f} = \Sigma m_i \ddot{\mathbf{r}}_i$$

- Substitution into the summation of the equations of motion gives:

$$\Sigma \mathbf{F} = m \ddot{\mathbf{r}} \quad \text{or} \quad \mathbf{F} = m \bar{\mathbf{a}}$$

$$\Sigma F_x = m \bar{a}_x \qquad \Sigma F_y = m \bar{a}_y \qquad \Sigma F_z = m \bar{a}_z$$



4/3

## Work-Energy

## Work-Energy Relation

$$(U_{1-2})_i = \Delta T_i$$

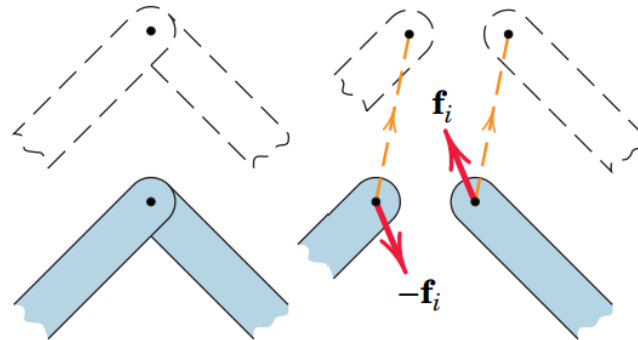


$$U_{1-2} = \Delta T \quad \text{or} \quad T_1 + U_{1-2} = T_2$$

$$U'_{1-2} = \Delta T + \Delta V$$

$$T_1 + V_1 + U'_{1-2} = T_2 + V_2$$

- ❖ For a rigid body or a system of rigid bodies joined by ideal frictionless connections, no net work is done by the internal interacting forces or moments in the connections.



4/3

## Work-Energy

## Kinetic Energy Expression

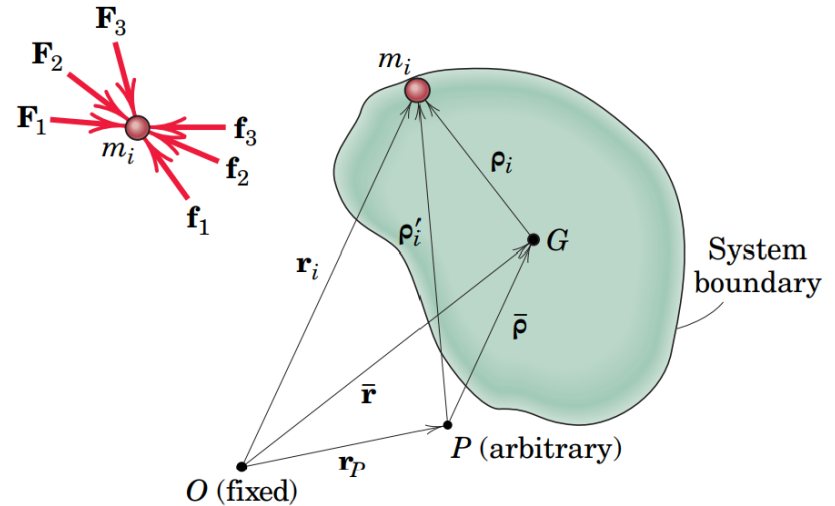
$$T = \Sigma \frac{1}{2} m_i v_i^2$$

$$\mathbf{v}_i = \bar{\mathbf{v}} + \dot{\boldsymbol{\rho}}_i$$

$$\begin{aligned} \rightarrow T &= \Sigma \frac{1}{2} m_i \mathbf{v}_i \cdot \mathbf{v}_i = \Sigma \frac{1}{2} m_i (\bar{\mathbf{v}} + \dot{\boldsymbol{\rho}}_i) \cdot (\bar{\mathbf{v}} + \dot{\boldsymbol{\rho}}_i) \\ &= \Sigma \frac{1}{2} m_i \bar{v}^2 + \Sigma \frac{1}{2} m_i |\dot{\boldsymbol{\rho}}_i|^2 + \Sigma m_i \bar{\mathbf{v}} \cdot \dot{\boldsymbol{\rho}}_i \end{aligned}$$

Because  $\boldsymbol{\rho}_i$  is measured from the mass center,  $\Sigma m_i \boldsymbol{\rho}_i = \mathbf{0}$

$$\rightarrow T = \frac{1}{2} m \bar{v}^2 + \Sigma \frac{1}{2} m_i |\dot{\boldsymbol{\rho}}_i|^2$$



## 4/4 Impulse-Momentum

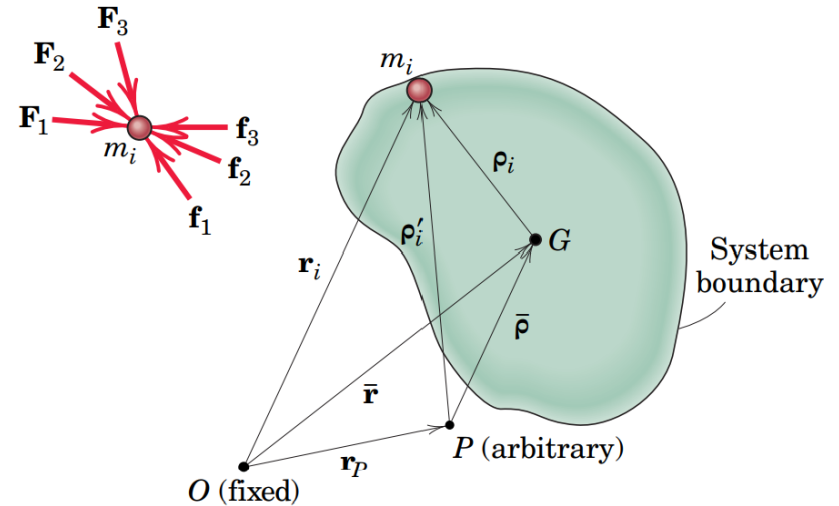
### Linear Momentum

$$\begin{aligned}\mathbf{G} &= \Sigma m_i(\bar{\mathbf{v}} + \dot{\boldsymbol{\rho}}_i) = \Sigma m_i \bar{\mathbf{v}} + \frac{d}{dt} \Sigma m_i \boldsymbol{\rho}_i \\ &= \bar{\mathbf{v}} \Sigma m_i + \frac{d}{dt} (\mathbf{0})\end{aligned}$$



$$\mathbf{G} = m \bar{\mathbf{v}}$$

$$\Sigma \mathbf{F} = \dot{\mathbf{G}}$$



## 4/4 Impulse-Momentum

## Angular Momentum

*About a Fixed Point O.*

$$\mathbf{H}_O = \Sigma(\mathbf{r}_i \times m_i \mathbf{v}_i)$$

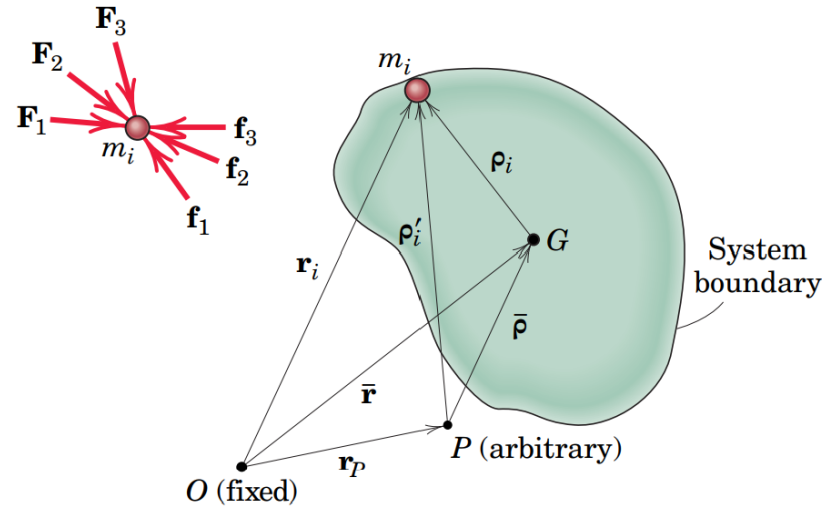
$$\dot{\mathbf{H}}_O = \Sigma(\dot{\mathbf{r}}_i \times m_i \mathbf{v}_i) + \Sigma(\mathbf{r}_i \times m_i \dot{\mathbf{v}}_i)$$



$$\Sigma(\mathbf{r}_i \times m_i \mathbf{a}_i) = \Sigma(\mathbf{r}_i \times \mathbf{F}_i)$$



$$\Sigma \mathbf{M}_O = \dot{\mathbf{H}}_O$$



## 4/4 Impulse-Momentum

## Angular Momentum

*About the Mass Center G.*

$$\mathbf{H}_G = \Sigma \boldsymbol{\rho}_i \times m_i \dot{\mathbf{r}}_i$$

$$\mathbf{H}_G = \Sigma \boldsymbol{\rho}_i \times m_i (\dot{\mathbf{r}} + \dot{\boldsymbol{\rho}}_i) = \Sigma \boldsymbol{\rho}_i \times m_i \dot{\mathbf{r}} + \Sigma \boldsymbol{\rho}_i \times m_i \dot{\boldsymbol{\rho}}_i$$



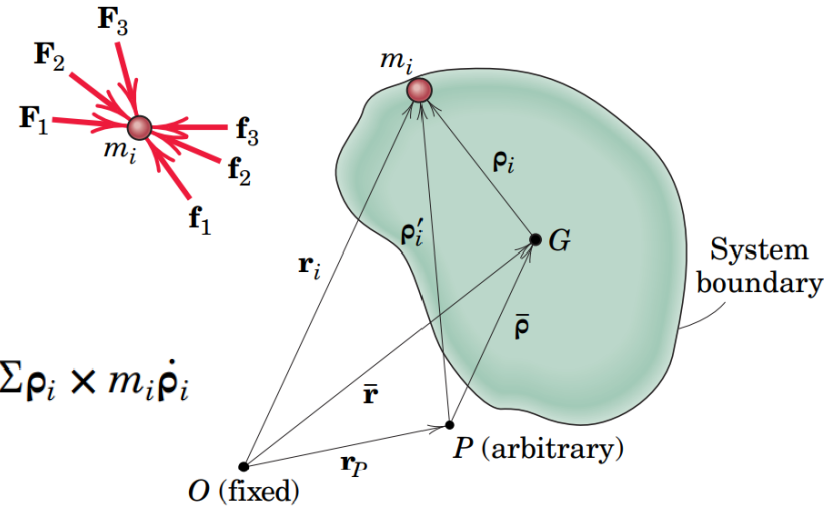
$$-\dot{\mathbf{r}} \times \Sigma m_i \boldsymbol{\rho}_i$$

$$\rightarrow \mathbf{H}_G = \Sigma \boldsymbol{\rho}_i \times m_i \dot{\boldsymbol{\rho}}_i$$

$$\dot{\mathbf{H}}_G = \Sigma \dot{\boldsymbol{\rho}}_i \times m_i (\dot{\mathbf{r}} + \dot{\boldsymbol{\rho}}_i) + \Sigma \boldsymbol{\rho}_i \times m_i \ddot{\mathbf{r}}_i$$



$$\Sigma \mathbf{M}_G = \dot{\mathbf{H}}_G$$





## 4/4 Impulse-Momentum

## Angular Momentum

*About an Arbitrary Point P.*

$$\mathbf{H}_P = \Sigma \rho'_i \times m_i \dot{\mathbf{r}}_i = \Sigma (\bar{\rho} + \rho_i) \times m_i \dot{\mathbf{r}}_i$$

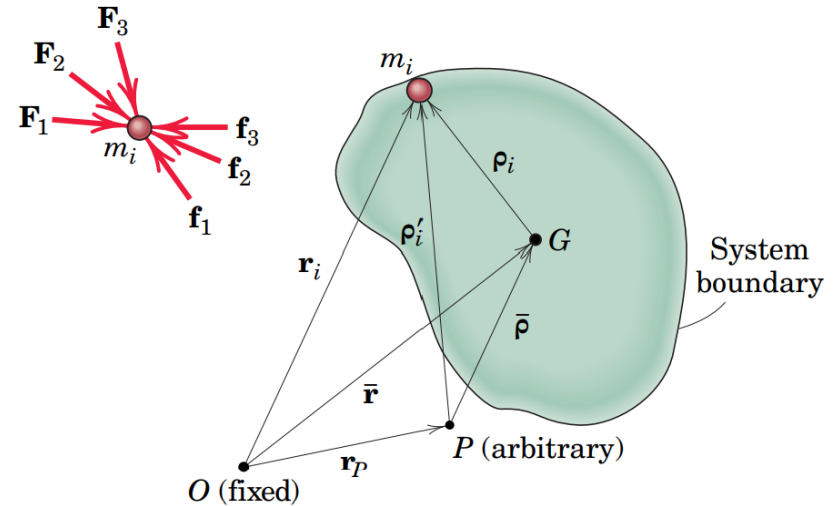


$$\mathbf{H}_P = \mathbf{H}_G + \bar{\rho} \times m \bar{\mathbf{v}}$$

$$\Sigma \mathbf{M}_P = \Sigma \mathbf{M}_G + \bar{\rho} \times \Sigma \mathbf{F}$$



$$\Sigma \mathbf{M}_P = \dot{\mathbf{H}}_G + \bar{\rho} \times m \bar{\mathbf{a}}$$



## 4/4 Impulse-Momentum

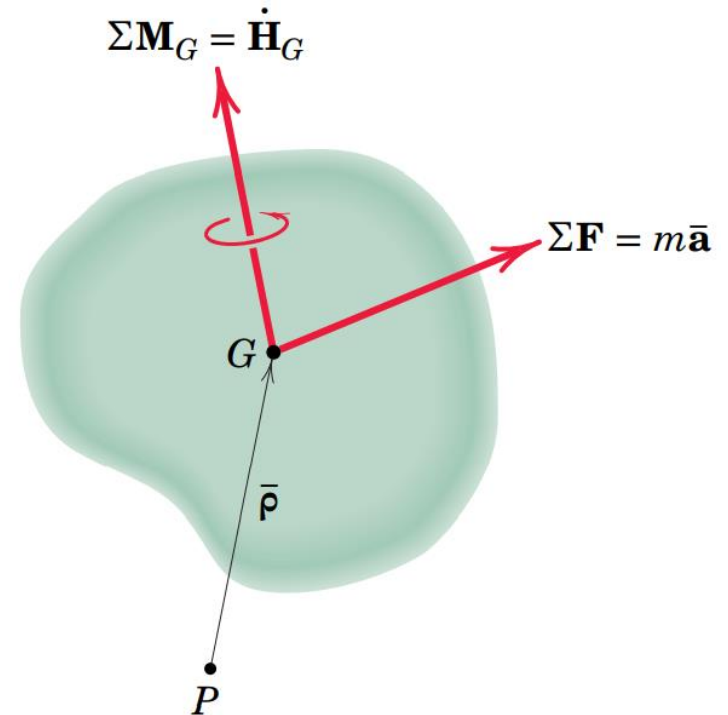
## Angular Momentum

*About an Arbitrary Point P.*

$$\mathbf{H}_P = \mathbf{H}_G + \bar{\rho} \times m\bar{\mathbf{v}}$$



$$\Sigma \mathbf{M}_P = \dot{\mathbf{H}}_G + \bar{\rho} \times m\bar{\mathbf{a}}$$



4/5

## Conservation of Energy and Momentum

### Conservation of Energy

- ❖ A mass system is said to be conservative if it does not lose energy by virtue of internal friction forces which do negative work or by virtue of inelastic members which dissipate energy upon cycling.
- ❖ If no work is done on a conservative system during an interval of motion by external forces (other than gravity or other potential forces), then none of the energy of the system is lost.

$$\Delta T + \Delta V = 0$$

$$T_1 + V_1 = T_2 + V_2$$



4/5

## Conservation of Energy and Momentum

### Conservation of Momentum

- The principle of conservation of linear momentum

$$\mathbf{G}_1 = \mathbf{G}_2$$

- The principle of conservation of angular momentum

$$(\mathbf{H}_O)_1 = (\mathbf{H}_O)_2 \quad \text{or} \quad (\mathbf{H}_G)_1 = (\mathbf{H}_G)_2$$



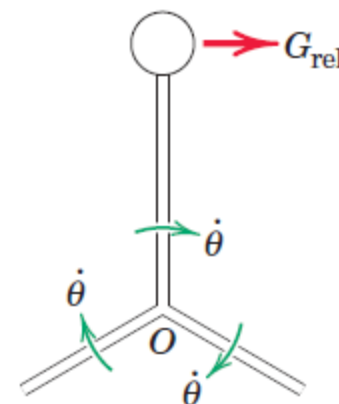
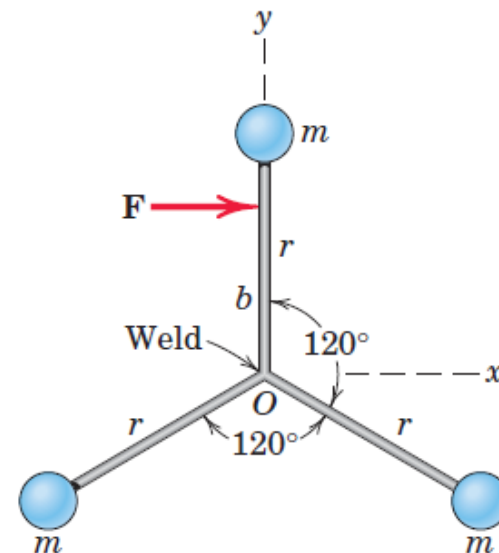
### SAMPLE PROBLEM 4/2

Each of the three balls has a mass  $m$  and is welded to the rigid equian-gular frame of negligible mass. The assembly rests on a smooth horizontal surface. If a force  $\mathbf{F}$  is suddenly applied to one bar as shown, deter-mine (a) the acceleration of point  $O$  and (b) the angular acceleration  $\ddot{\theta}$  of the frame.

$$[\Sigma \mathbf{F} = m\bar{\mathbf{a}}] \quad F\mathbf{i} = 3m\bar{\mathbf{a}} \quad \bar{\mathbf{a}} = \mathbf{a}_O = \frac{F}{3m}\mathbf{i}$$

$$H_O = H_G = 3(mr\dot{\theta})r = 3mr^2\dot{\theta}$$

$$[\Sigma \mathbf{M}_G = \dot{\mathbf{H}}_G] \quad Fb = \frac{d}{dt}(3mr^2\dot{\theta}) = 3mr^2\ddot{\theta} \quad \text{so} \quad \ddot{\theta} = \frac{Fb}{3mr^2}$$



❑ CONTENTS:

- ❖ Chapter 1: Introduction to Dynamics
- ❖ Chapter 2: Kinematics of Particles
- ❖ Chapter 3: Kinetics of Particles
- ❖ Chapter 4: Kinetics of Systems of Particles
- ➔ ❖ Chapter 5: **Plane Kinematics of Rigid Bodies**
- ❖ Chapter 6: Plane Kinetics of Rigid Bodies

5/1

## Introduction

### Rigid-Body Assumption

- A rigid body:
  - ❖ A system of particles for which the distances between the particles remain unchanged.
  - ❖ If the movements associated with the changes in shape are very small compared with the movements of the body as a whole, then the assumption of rigidity is usually acceptable.

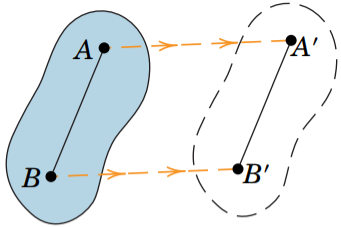
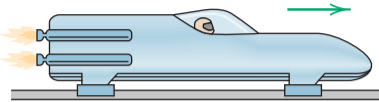
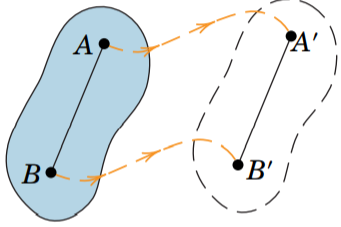
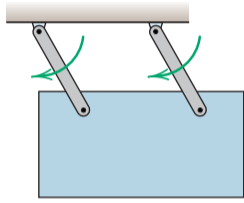
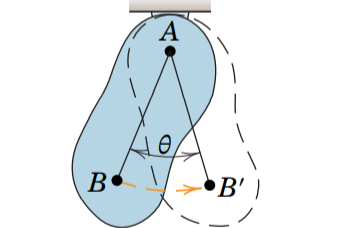
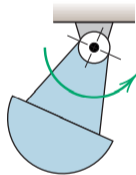
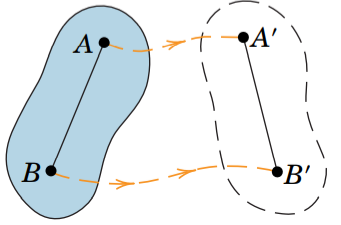
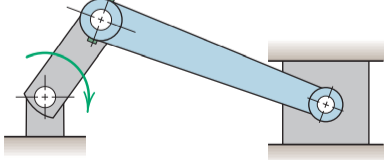


5/1

## Introduction

## Plane Motion

- ❖ Translation
- ❖ Rotation
- ❖ General plane motion

	Type of Rigid-Body Plane Motion	Example
(a) Rectilinear translation		 Rocket test sled
(b) Curvilinear translation		 Parallel-link swinging plate
(c) Fixed-axis rotation		 Compound pendulum
(d) General plane motion		 Connecting rod in a reciprocating engine



5/2

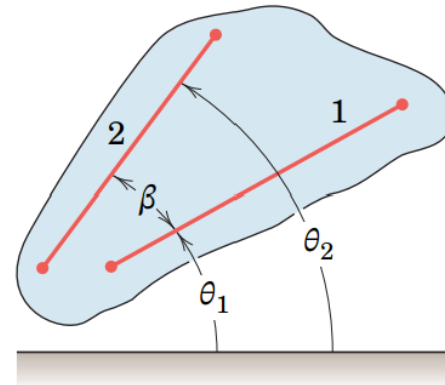
## Rotation

- The rotation of a rigid body is described by its angular motion.

$$\theta_2 = \theta_1 + \beta$$

$$\dot{\theta}_2 = \dot{\theta}_1$$

$$\ddot{\theta}_2 = \ddot{\theta}_1$$



- ❖ All lines on a rigid body in its plane of motion have the same angular displacement, the same angular velocity, and the same angular acceleration

5/2

## Rotation

## Angular-Motion Relations

The angular velocity  $\omega$  and angular acceleration  $\alpha$

$$\omega = \frac{d\theta}{dt} = \dot{\theta}$$

$$\alpha = \frac{d\omega}{dt} = \dot{\omega} \quad \text{or} \quad \alpha = \frac{d^2\theta}{dt^2} = \ddot{\theta}$$

$$\omega d\omega = \alpha d\theta \quad \text{or} \quad \dot{\theta} d\dot{\theta} = \ddot{\theta} d\theta$$

❖ For rotation with constant angular acceleration:

$$\omega = \omega_0 + \alpha t$$

$$\omega^2 = \omega_0^2 + 2\alpha(\theta - \theta_0)$$

$$\theta = \theta_0 + \omega_0 t + \frac{1}{2}\alpha t^2$$



5/2

## Rotation

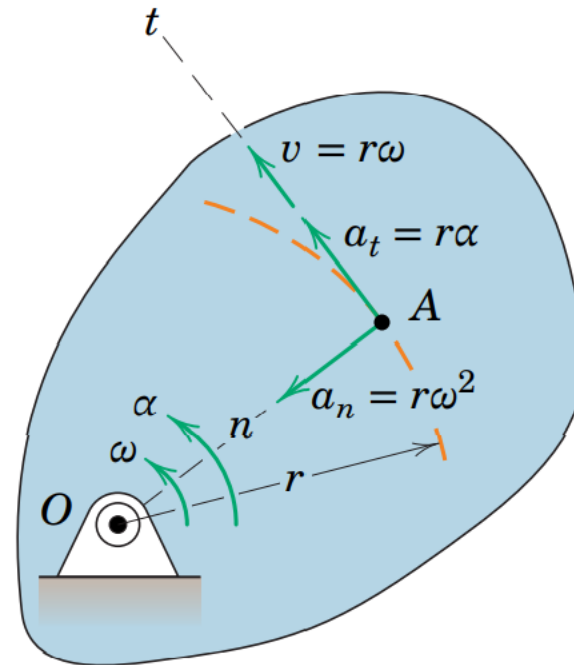
## Angular-Motion Relations

## Rotation about a Fixed Axis

$$v = r\omega$$

$$a_n = r\omega^2 = v^2/r = v\omega$$

$$a_t = r\alpha$$



5/2

## Rotation

## Angular-Motion Relations

- Using the cross-product relationship

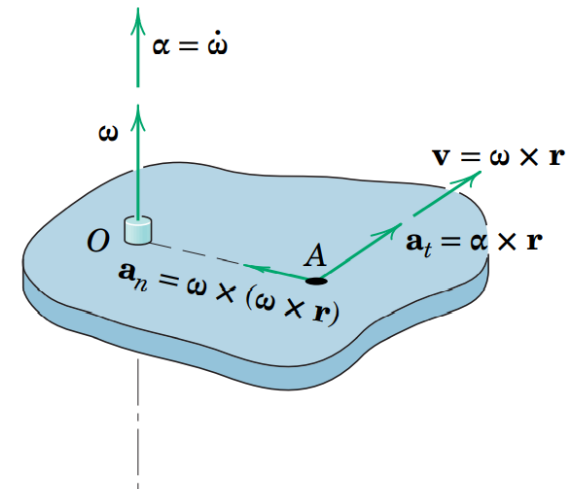
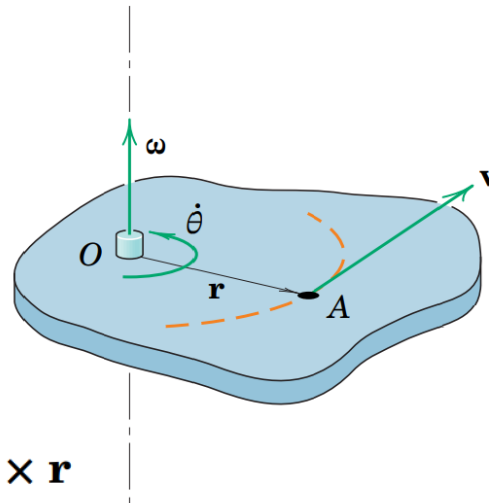
$$\mathbf{v} = \dot{\mathbf{r}} = \boldsymbol{\omega} \times \mathbf{r}$$

$$\mathbf{r} \times \boldsymbol{\omega} = -\mathbf{v}$$

$$\mathbf{a} = \dot{\mathbf{v}} = \boldsymbol{\omega} \times \dot{\mathbf{r}} + \dot{\boldsymbol{\omega}} \times \mathbf{r}$$

$$= \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) + \dot{\boldsymbol{\omega}} \times \mathbf{r}$$

$$= \boldsymbol{\omega} \times \mathbf{v} + \boldsymbol{\alpha} \times \mathbf{r}$$



$$\mathbf{v} = \boldsymbol{\omega} \times \mathbf{r}$$

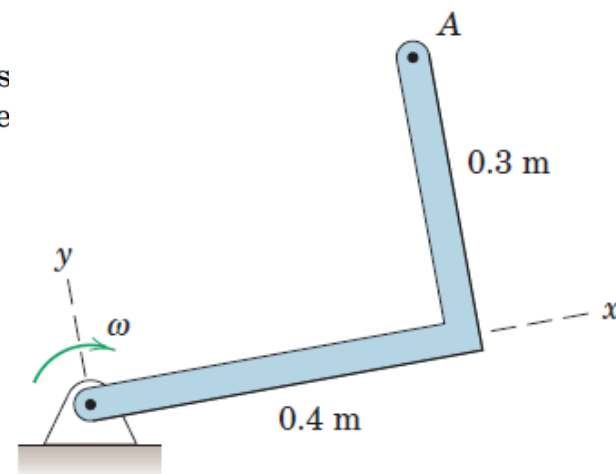


$$\mathbf{a}_n = \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})$$

$$\mathbf{a}_t = \boldsymbol{\alpha} \times \mathbf{r}$$

### SAMPLE PROBLEM 5/3

The right-angle bar rotates clockwise with an angular velocity which is decreasing at the rate of  $4 \text{ rad/s}^2$ . Write the vector expressions for the velocity and acceleration of point  $A$  when  $\omega = 2 \text{ rad/s}$ .



$$\boldsymbol{\omega} = -2\mathbf{k} \text{ rad/s} \quad \text{and} \quad \boldsymbol{\alpha} = +4\mathbf{k} \text{ rad/s}^2$$

$$[\mathbf{v} = \boldsymbol{\omega} \times \mathbf{r}] \quad \mathbf{v} = -2\mathbf{k} \times (0.4\mathbf{i} + 0.3\mathbf{j}) = 0.6\mathbf{i} - 0.8\mathbf{j} \text{ m/s}$$

$$[\mathbf{a}_n = \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})] \quad \mathbf{a}_n = -2\mathbf{k} \times (0.6\mathbf{i} - 0.8\mathbf{j}) = -1.6\mathbf{i} - 1.2\mathbf{j} \text{ m/s}^2$$

$$[\mathbf{a}_t = \boldsymbol{\alpha} \times \mathbf{r}] \quad \mathbf{a}_t = 4\mathbf{k} \times (0.4\mathbf{i} + 0.3\mathbf{j}) = -1.2\mathbf{i} + 1.6\mathbf{j} \text{ m/s}^2$$

$$[\mathbf{a} = \mathbf{a}_n + \mathbf{a}_t] \quad \mathbf{a} = -2.8\mathbf{i} + 0.4\mathbf{j} \text{ m/s}^2$$

$$v = \sqrt{0.6^2 + 0.8^2} = 1 \text{ m/s} \quad \text{and} \quad a = \sqrt{2.8^2 + 0.4^2} = 2.83 \text{ m/s}^2$$

5/3

## Absolute Motion

- ❖ We now develop the approach of absolute-motion analysis to describe the plane kinematics of rigid bodies.
- ❖ In this approach, we make use of the geometric relations which define the configuration of the body involved and then proceed to take the time derivatives of the defining geometric relations to obtain velocities and accelerations.
- ❖ The absolute-motion approach to rigid-body kinematics is quite straightforward, provided the configuration lends itself to a geometric description which is not overly complex. If the geometric configuration is awkward or complex, analysis by the principles of relative motion may be preferable.



### SAMPLE PROBLEM 5/4

A wheel of radius  $r$  rolls on a flat surface without slipping. Determine the angular motion of the wheel in terms of the linear motion of its center  $O$ . Also determine the acceleration of a point on the rim of the wheel as the point comes into contact with the surface on which the wheel rolls.

$$s = r\theta$$

$$v_O = r\omega$$

$$a_O = r\alpha$$

$$x = s - r \sin \theta = r(\theta - \sin \theta)$$

$$y = r - r \cos \theta = r(1 - \cos \theta)$$

$$\dot{x} = r\dot{\theta}(1 - \cos \theta) = v_O(1 - \cos \theta)$$

$$\dot{y} = r\dot{\theta} \sin \theta = v_O \sin \theta$$

$$\ddot{x} = \dot{v}_O(1 - \cos \theta) + v_O\dot{\theta} \sin \theta$$

$$\ddot{y} = \dot{v}_O \sin \theta + v_O\dot{\theta} \cos \theta$$

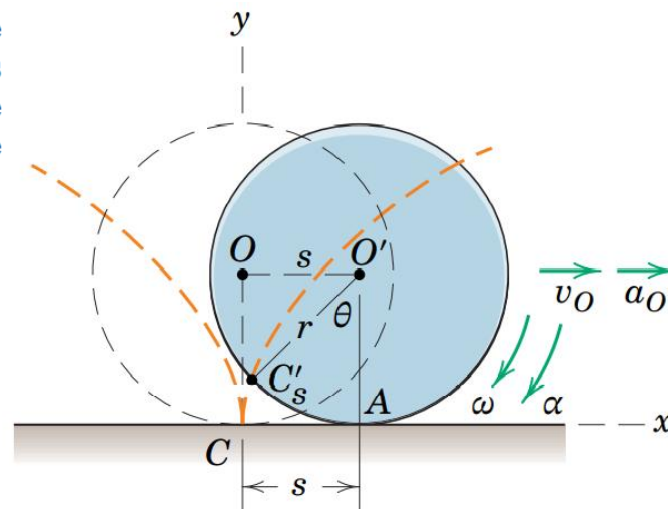
$$= a_O(1 - \cos \theta) + r\omega^2 \sin \theta$$

$$= a_O \sin \theta + r\omega^2 \cos \theta$$

$$\theta = 0 \quad \rightarrow \quad \ddot{x} = 0 \quad \text{and} \quad \ddot{y} = r\omega^2$$

$$\mathbf{v} = \dot{x}\mathbf{i} + \dot{y}\mathbf{j}$$

$$\mathbf{a} = \ddot{x}\mathbf{i} + \ddot{y}\mathbf{j}$$

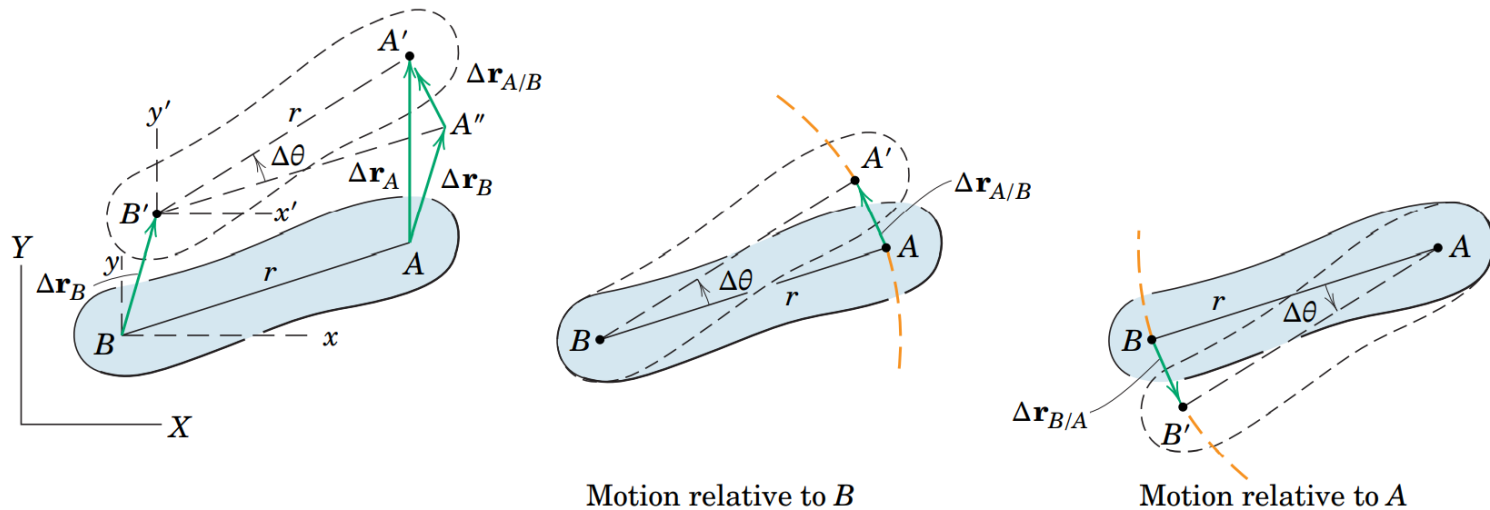


## 5/4 Relative Velocity

- The principles of relative motion:

$$\mathbf{v}_A = \mathbf{v}_B + \mathbf{v}_{A/B}$$

### Relative Velocity Due to Rotation



$$\Delta \mathbf{r}_A = \Delta \mathbf{r}_B + \Delta \mathbf{r}_{A/B}$$



$$\mathbf{v}_A = \mathbf{v}_B + \mathbf{v}_{A/B}$$



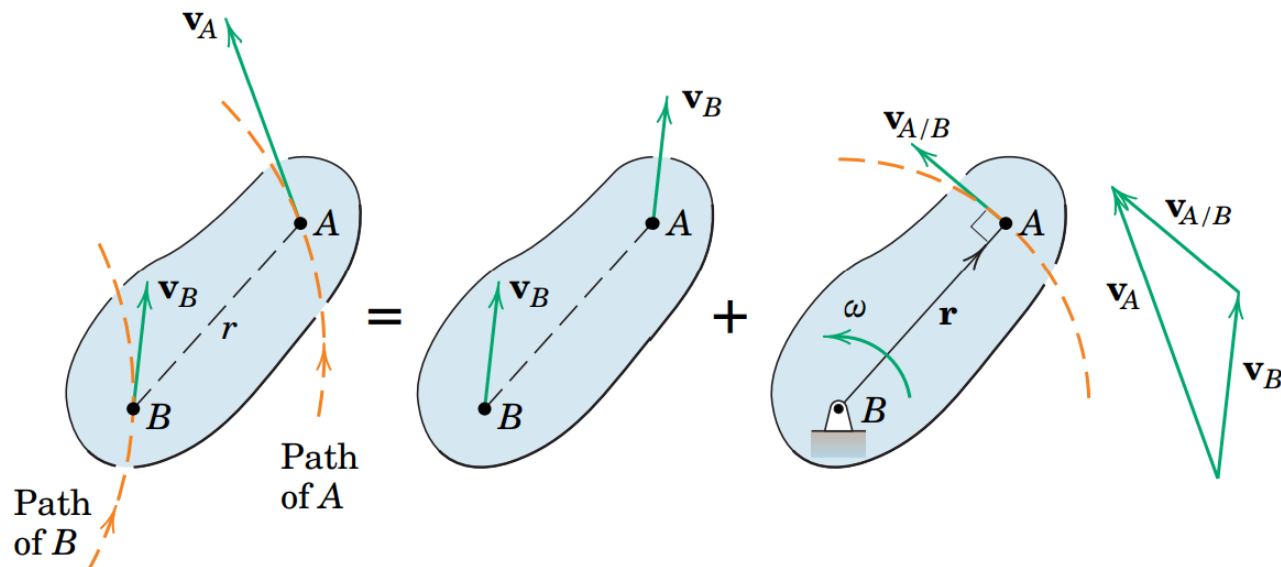
$$v_{A/B} = r\omega$$

$$\mathbf{v}_{A/B} = \boldsymbol{\omega} \times \mathbf{r}$$



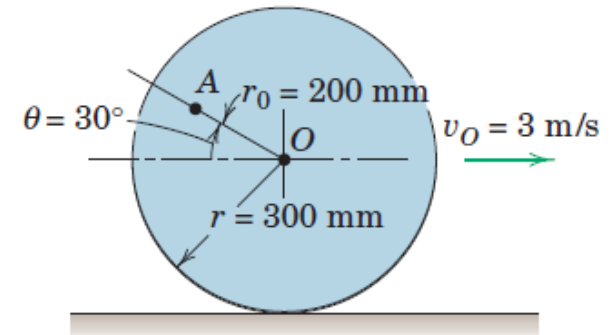
## 5/4 Relative Velocity

### Interpretation of the Relative-Velocity Equation



**SAMPLE PROBLEM 5/7**

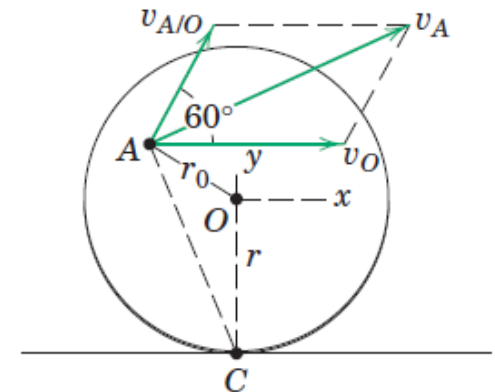
The wheel of radius  $r = 300$  mm rolls to the right without slipping and has a velocity  $v_O = 3$  m/s of its center  $O$ . Calculate the velocity of point  $A$  on the wheel for the instant represented.

**Solution I (Scalar-Geometric)**

$$\mathbf{v}_A = \mathbf{v}_O + \mathbf{v}_{A/O}$$

$$[v_{A/O} = r_0 \dot{\theta}] \quad v_{A/O} = 0.2(10) = 2 \text{ m/s}$$

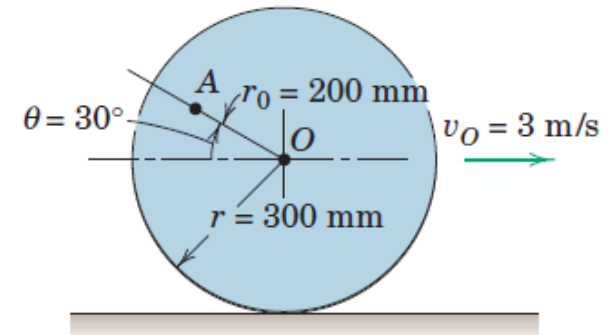
$$v_A^2 = 3^2 + 2^2 + 2(3)(2) \cos 60^\circ = 19 \text{ (m/s)}^2 \quad v_A = 4.36 \text{ m/s}$$



$$v_{A/C} = \overline{AC} \omega = \frac{\overline{AC}}{\overline{OC}} v_O = \frac{0.436}{0.300} (3) = 4.36 \text{ m/s} \quad v_A = v_{A/C} = 4.36 \text{ m/s}$$

**SAMPLE PROBLEM 5/7**

The wheel of radius  $r = 300$  mm rolls to the right without slipping and has a velocity  $v_O = 3$  m/s of its center  $O$ . Calculate the velocity of point  $A$  on the wheel for the instant represented.

**Solution II (Vector)**

$$\mathbf{v}_A = \mathbf{v}_O + \mathbf{v}_{A/O} = \mathbf{v}_O + \boldsymbol{\omega} \times \mathbf{r}_0$$

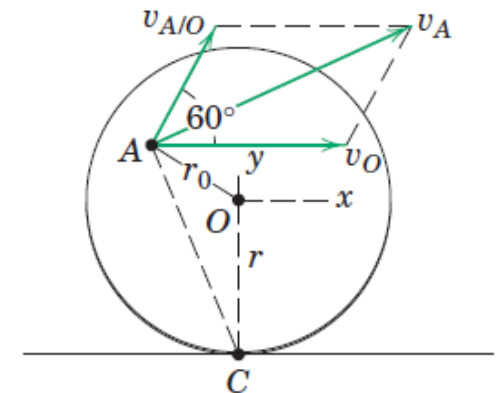
$$\boldsymbol{\omega} = -10\mathbf{k} \text{ rad/s}$$

$$\mathbf{r}_0 = 0.2(-\mathbf{i} \cos 30^\circ + \mathbf{j} \sin 30^\circ) = -0.1732\mathbf{i} + 0.1\mathbf{j} \text{ m}$$

$$\mathbf{v}_O = 3\mathbf{i} \text{ m/s}$$

$$\mathbf{v}_A = 3\mathbf{i} + \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ 0 & 0 & -10 \\ -0.1732 & 0.1 & 0 \end{vmatrix} = 3\mathbf{i} + 1.732\mathbf{j} + \mathbf{i}$$

$$= 4\mathbf{i} + 1.732\mathbf{j} \text{ m/s}$$



$$\rightarrow v_A = \sqrt{4^2 + (1.732)^2} = \sqrt{19} = 4.36 \text{ m/s}$$

5/5

## Instantaneous Center of Zero Velocity

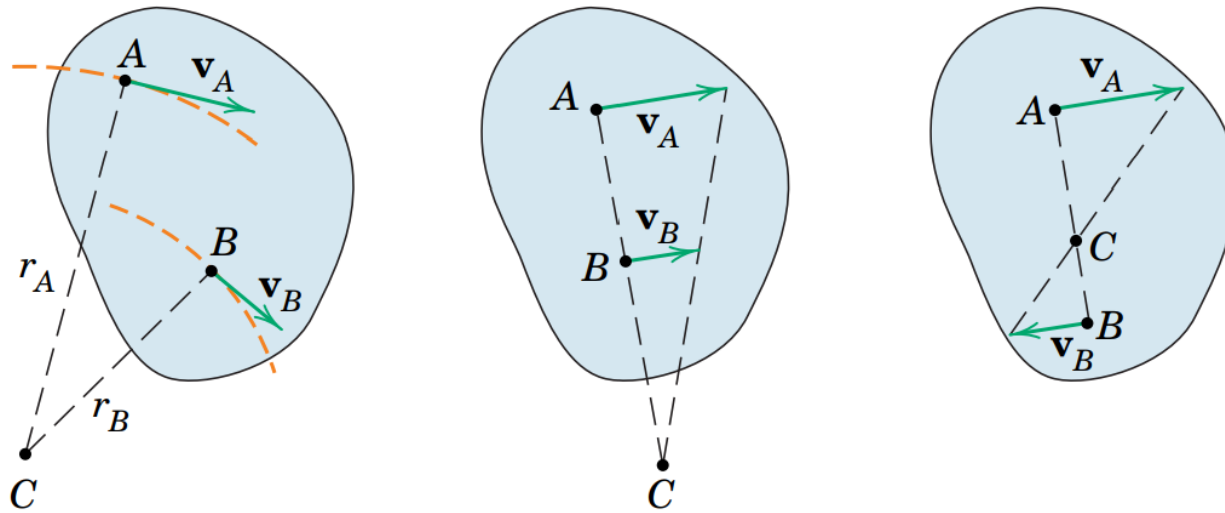
- We can solve the problem by choosing a unique reference point which momentarily has zero velocity.
- As far as velocities are concerned, the body may be considered to be in pure rotation about an axis, normal to the plane of motion, passing through this point.
- This axis is called the instantaneous axis of zero velocity, and the intersection of this axis with the plane of motion is known as the instantaneous center of zero velocity.



5/5

## Instantaneous Center of Zero Velocity

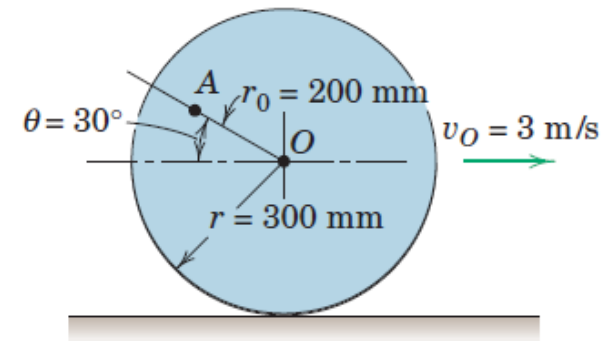
## Locating the Instantaneous Center



$$\omega = \frac{v_A}{r_A}$$

**SAMPLE PROBLEM 5/11**

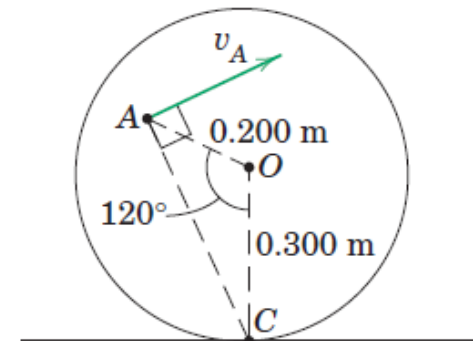
The wheel of Sample Problem 5/7, shown again here, rolls to the right without slipping, with its center  $O$  having a velocity  $v_O = 3$  m/s. Locate the instantaneous center of zero velocity and use it to find the velocity of point  $A$  for the position indicated.



$$[\omega = v/r] \quad \omega = v_O / \overline{OC} = 3/0.300 = 10 \text{ rad/s}$$

$$\overline{AC} = \sqrt{(0.300)^2 + (0.200)^2 - 2(0.300)(0.200) \cos 120^\circ} = 0.436 \text{ m}$$

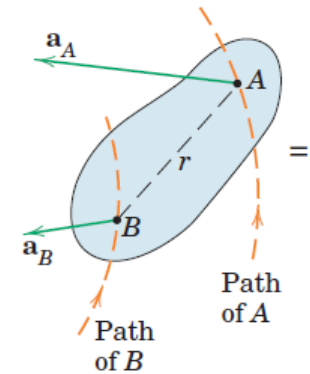
$$[v = r\omega] \quad v_A = \overline{AC}\omega = 0.436(10) = 4.36 \text{ m/s}$$



## 5/6 Relative Acceleration

- The relative-acceleration equation:

$$\mathbf{a}_A = \mathbf{a}_B + \mathbf{a}_{A/B}$$



### Relative Acceleration Due to Rotation

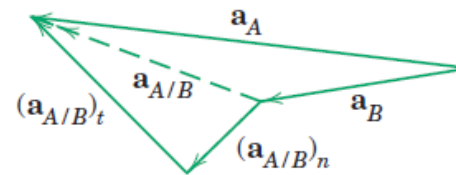
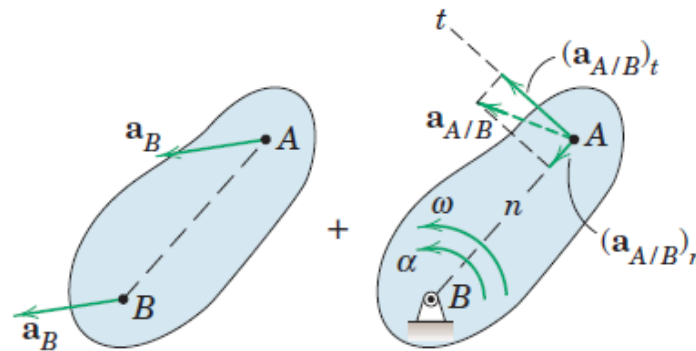
$$\mathbf{a}_A = \mathbf{a}_B + (\mathbf{a}_{A/B})_n + (\mathbf{a}_{A/B})_t$$

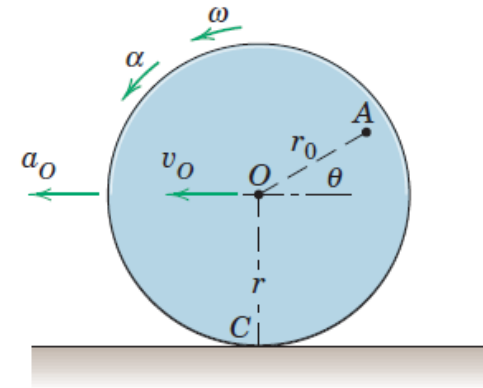
$$(\mathbf{a}_{A/B})_n = v_{A/B}^2 / r = r\omega^2$$

$$(\mathbf{a}_{A/B})_t = \dot{v}_{A/B} = r\alpha$$

$$(\mathbf{a}_{A/B})_n = \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})$$

$$(\mathbf{a}_{A/B})_t = \boldsymbol{\alpha} \times \mathbf{r}$$





### SAMPLE PROBLEM 5/13

The wheel of radius  $r$  rolls to the left without slipping and, at the instant considered, the center  $O$  has a velocity  $\mathbf{v}_O$  and an acceleration  $\mathbf{a}_O$  to the left. Determine the acceleration of points  $A$  and  $C$  on the wheel for the instant considered.

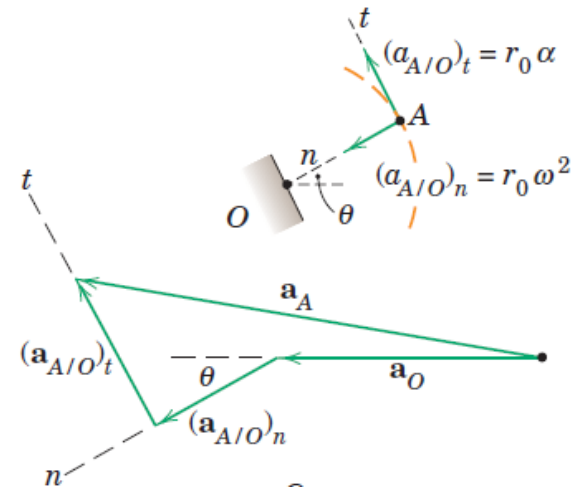
$$\omega = v_O/r \quad \text{and} \quad \alpha = a_O/r$$

$$\mathbf{a}_A = \mathbf{a}_O + \mathbf{a}_{A/O} = \mathbf{a}_O + (\mathbf{a}_{A/O})_n + (\mathbf{a}_{A/O})_t$$

$$(\mathbf{a}_{A/O})_n = r_0 \omega^2 = r_0 \left( \frac{v_O}{r} \right)^2$$

$$(\mathbf{a}_{A/O})_t = r_0 \alpha = r_0 \left( \frac{a_O}{r} \right)$$

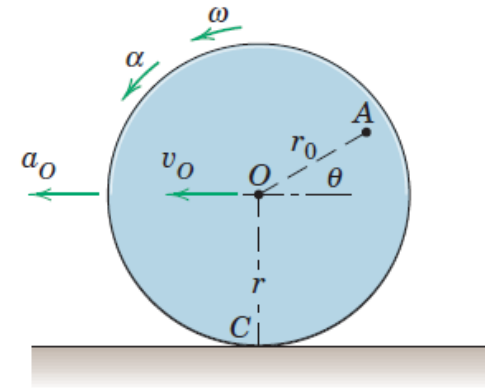
$$\begin{aligned} a_A &= \sqrt{(\mathbf{a}_A)_n^2 + (\mathbf{a}_A)_t^2} \\ &= \sqrt{[a_O \cos \theta + (\mathbf{a}_{A/O})_n]^2 + [a_O \sin \theta + (\mathbf{a}_{A/O})_t]^2} \\ &= \sqrt{(r\alpha \cos \theta + r_0 \omega^2)^2 + (r\alpha \sin \theta + r_0 \alpha)^2} \end{aligned}$$





**SAMPLE PROBLEM 5/13**

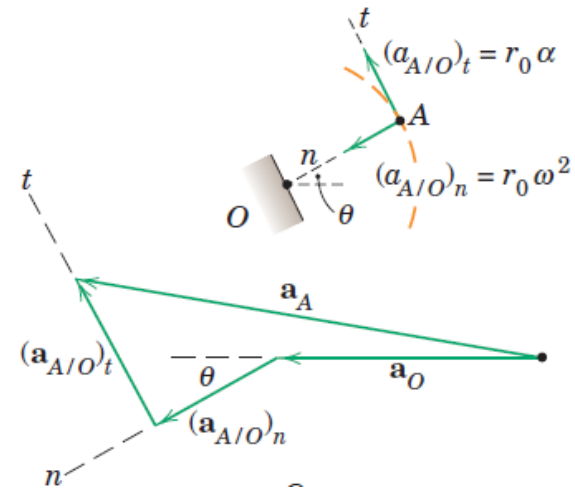
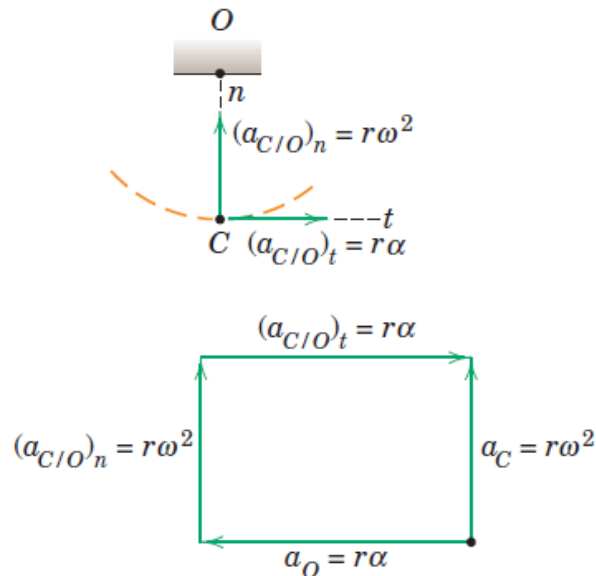
The wheel of radius  $r$  rolls to the left without slipping and, at the instant considered, the center  $O$  has a velocity  $\mathbf{v}_O$  and an acceleration  $\mathbf{a}_O$  to the left. Determine the acceleration of points  $A$  and  $C$  on the wheel for the instant considered.



$$\mathbf{a}_C = \mathbf{a}_O + \mathbf{a}_{C/O}$$



$$a_C = r\omega^2$$



5/7

## Motion Relative to Rotating Axes

## Rotating versus Nonrotating Systems

$$\begin{aligned}
 \mathbf{a}_A &= \mathbf{a}_B + \underbrace{\dot{\boldsymbol{\omega}} \times \mathbf{r} + \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})}_{\mathbf{a}_{P/B}} + \underbrace{2\boldsymbol{\omega} \times \mathbf{v}_{\text{rel}} + \mathbf{a}_{\text{rel}}}_{\mathbf{a}_{A/P}} \\
 \mathbf{a}_A &= \mathbf{a}_B + \mathbf{a}_{P/B} + \mathbf{a}_{A/P} \\
 \mathbf{a}_A &= \mathbf{a}_P + \mathbf{a}_{A/P} \\
 \mathbf{a}_A &= \mathbf{a}_B + \mathbf{a}_{A/B}
 \end{aligned}$$



$$\mathbf{a}_A = \mathbf{a}_P + 2\boldsymbol{\omega} \times \mathbf{v}_{\text{rel}} + \mathbf{a}_{\text{rel}}$$

❑ CONTENTS:

- ❖ Chapter 1: Introduction to Dynamics
- ❖ Chapter 2: Kinematics of Particles
- ❖ Chapter 3: Kinetics of Particles
- ❖ Chapter 4: Kinetics of Systems of Particles
- ❖ Chapter 5: Plane Kinematics of Rigid Bodies
- ➔ ❖ Chapter 6: **Plane Kinetics of Rigid Bodies**



6/1

## Introduction

- The kinetics of rigid bodies treats the relationships between the external forces acting on a body and the corresponding translational and rotational motions of the body
  
- For our purpose, a body which can be approximated as a thin slab with its motion confined to the plane of the slab will be considered to be in plane motion.
  - ✓ Section A: forces and moments to instantaneous linear and angular accelerations relations
  - ✓ Section B: method of work and energy
  - ✓ Section C: methods of impulse and momentum

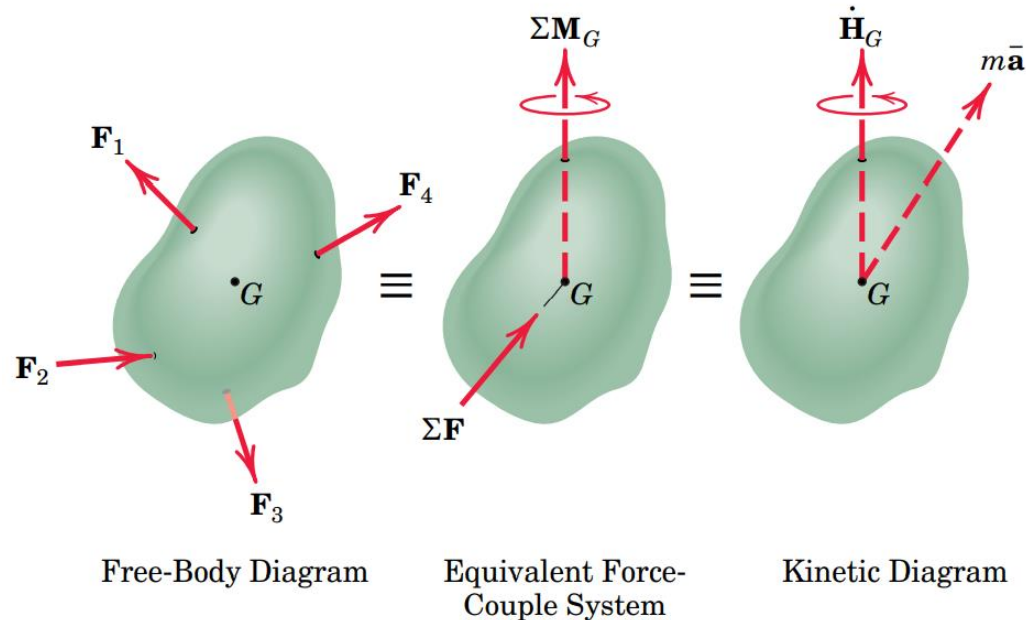


## SECTION A Force, Mass, and Acceleration

### 6/2 General Equations of Motion

□ The force equation:  $\Sigma \mathbf{F} = m\bar{\mathbf{a}}$

□ The moment equation taken about the mass center:  $\Sigma \mathbf{M}_G = \dot{\mathbf{H}}_G$



6/2

## General Equations of Motion

### Plane-Motion Equations

- ❖ A rigid body moving with plane motion in the x-y plane
  - ✓ The mass moment of inertia:

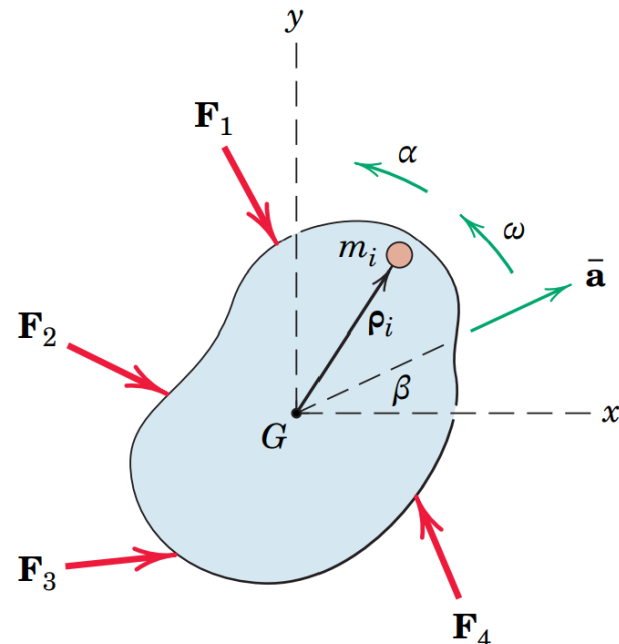
$$H_G = \sum \rho_i^2 m_i \omega = \omega \sum \rho_i^2 m_i \quad \rightarrow \quad \int \rho^2 dm \quad \rightarrow \quad \bar{I}$$

$$H_G = \bar{I} \omega \quad \rightarrow \quad \Sigma M_G = \dot{H}_G = \bar{I} \dot{\omega} = \bar{I} \alpha$$

$$\Sigma \mathbf{F} = m \bar{\mathbf{a}}$$

$$\Sigma M_G = \bar{I} \alpha$$

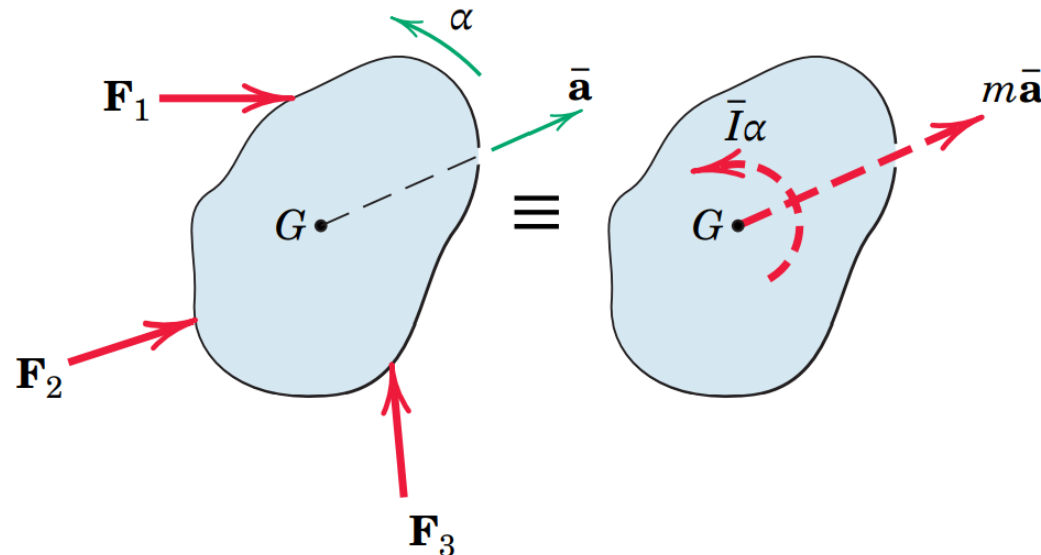
using x-y, n-t, or r-θ coordinates



6/2

## General Equations of Motion

- ✓ The free-body diagram discloses the forces and moments appearing on the lefthand side of equations of motion.
- ✓ The kinetic diagram discloses the resulting dynamic response in terms of the translational term and the rotational term which appear on the right-hand side of equations of motion.



Free-Body Diagram

Kinetic Diagram

## 6/2 General Equations of Motion

### Alternative Moment Equations

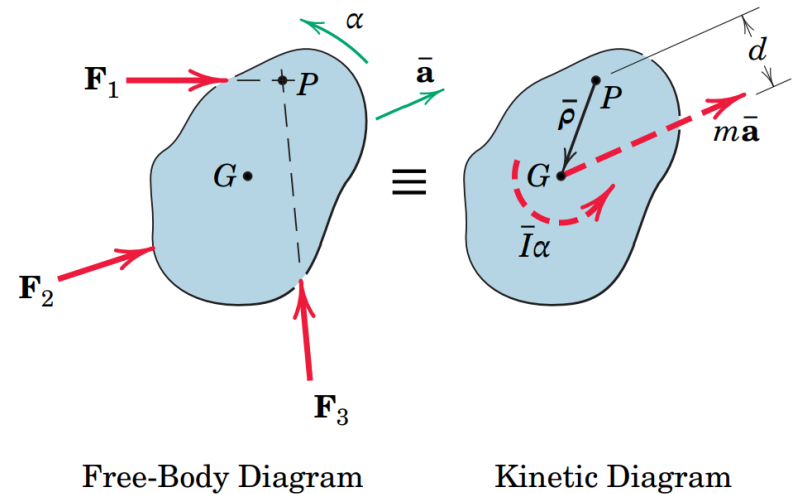
- ❖ General equation for moments about an arbitrary point P

$$\Sigma \mathbf{M}_P = \dot{\mathbf{H}}_G + \bar{\rho} \times m \bar{\mathbf{a}}$$

$$\rightarrow \Sigma M_P = \bar{I} \alpha + m \bar{a} d$$

$$\Sigma \mathbf{M}_P = (\dot{\mathbf{H}}_P)_{\text{rel}} + \bar{\rho} \times m \mathbf{a}_P$$

$$\rightarrow \Sigma \mathbf{M}_P = I_P \alpha + \bar{\rho} \times m \mathbf{a}_P$$



- ✓ When point P becomes a point O fixed in an inertial reference system

$$\rightarrow \Sigma M_O = I_O \alpha$$

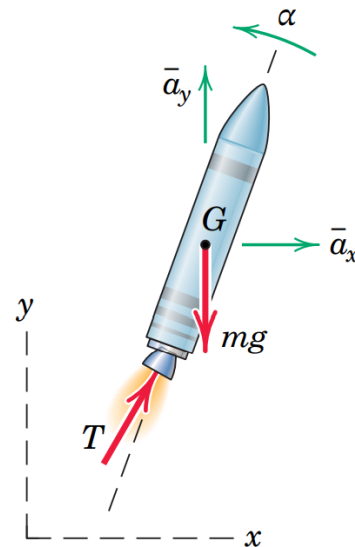


6/2

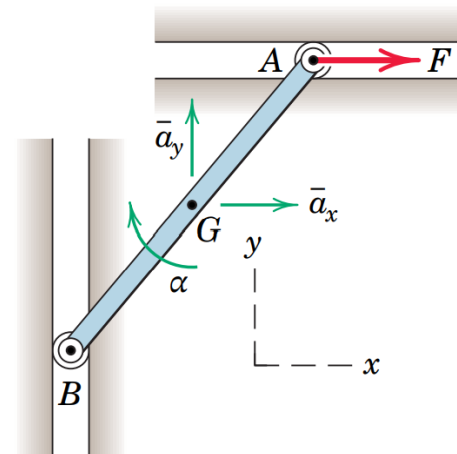
## General Equations of Motion

### Unconstrained and Constrained Motion

- ❖ The motion of a rigid body may be unconstrained or constrained.
- ❖ The two components  $a_x$  and  $a_y$  of the mass center acceleration and the angular acceleration  $\alpha$  may be determined independently or not.
- ❖ In general, dynamics problems which involve physical constraints to motion require a kinematic analysis relating linear to angular acceleration before the force and moment equations of motion can be solved.



(a) Unconstrained Motion



(b) Constrained Motion

دانشکده مهندسی مکانیک - درس مکاترونیک ۱

6/2

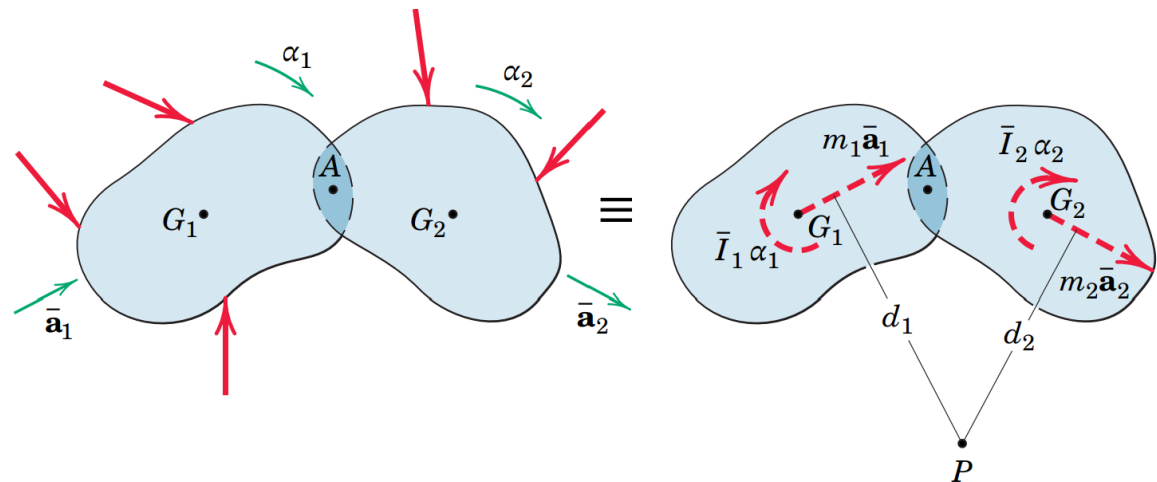
## General Equations of Motion

### Systems of Interconnected Bodies

- ✓ In problems dealing with two or more connected rigid bodies whose motions are related kinematically, it is convenient to analyze the bodies as an entire system.

$$\Sigma \mathbf{F} = \Sigma m \bar{\mathbf{a}}$$

$$\Sigma M_P = \Sigma \bar{I} \alpha + \Sigma m \bar{\mathbf{a}} d$$

Free-Body Diagram  
of SystemKinetic Diagram  
of System

- ✓ If there are more than three remaining unknowns, the E.O.M. is not sufficient to solve the problem.

6/3

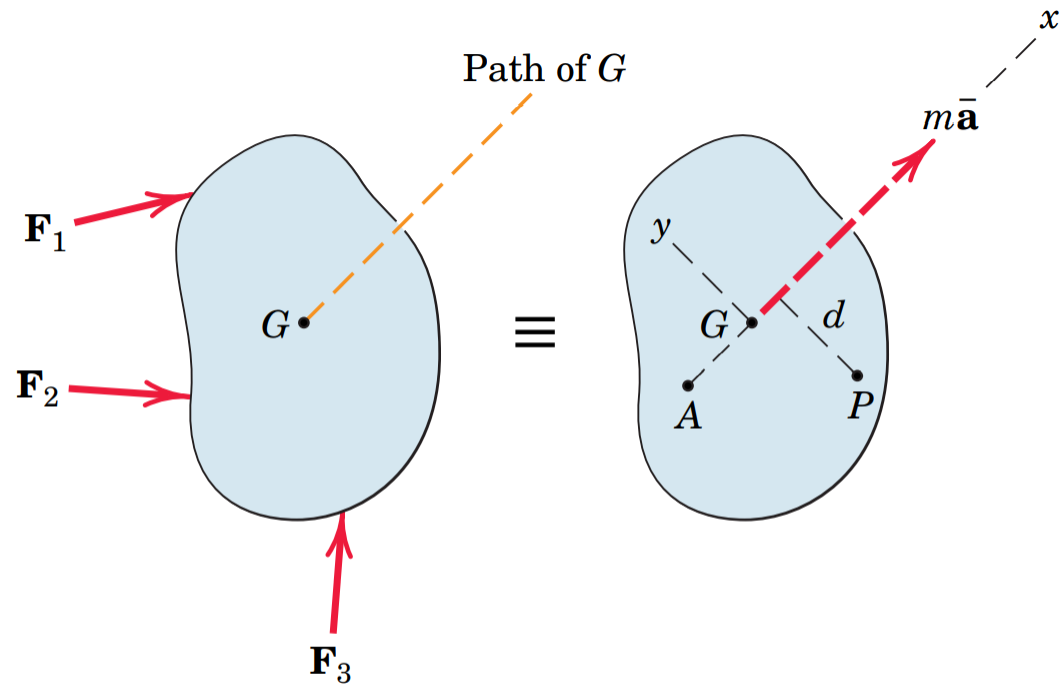
## Translation

- There is no angular motion of the translating body, so that both  $\omega$  and  $\alpha$  are zero.

❖ Rectilinear Translation:

$$\Sigma \mathbf{F} = m\bar{\mathbf{a}}$$

$$\Sigma M_G = \bar{I}\alpha = 0$$



Free-Body Diagram

Kinetic Diagram

(a) Rectilinear Translation

( $\alpha = 0, \omega = 0$ )

6/3

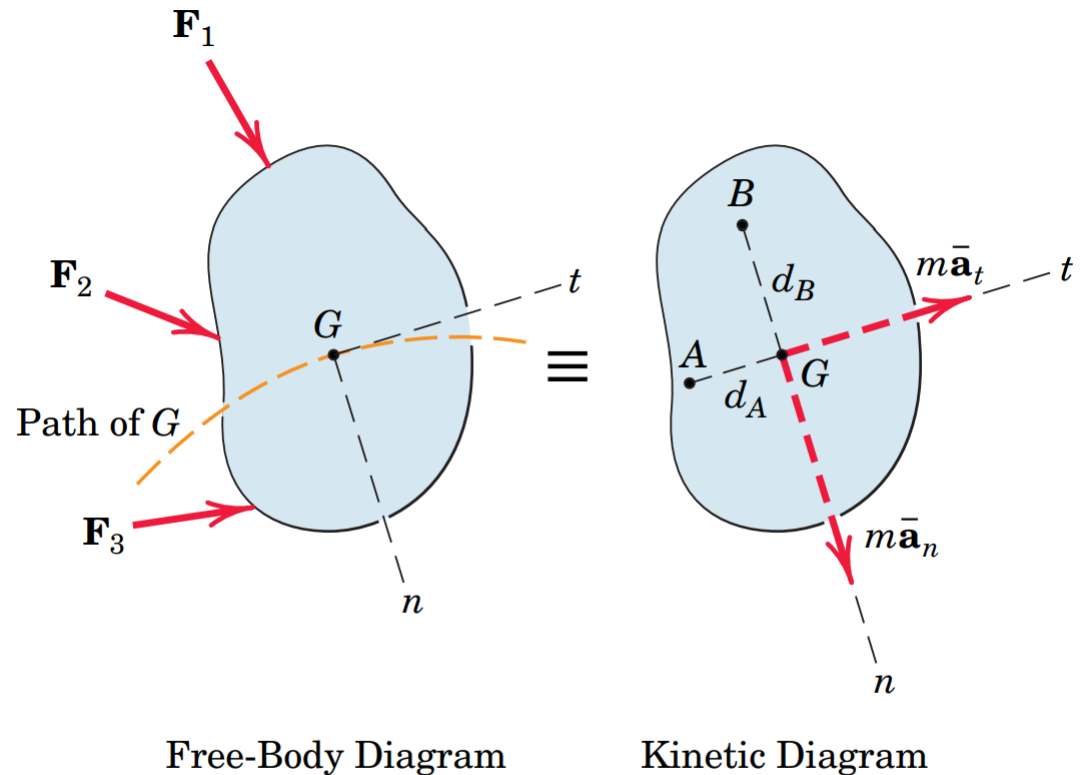
## Translation

- There is no angular motion of the translating body, so that both  $\omega$  and  $\alpha$  are zero.

❖ Curvilinear Translation:

$$\Sigma \mathbf{F} = m\bar{\mathbf{a}}$$

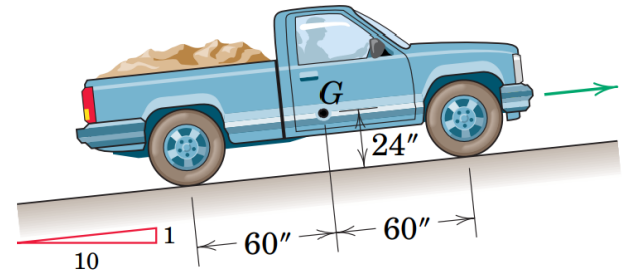
$$\Sigma M_G = \bar{I}\alpha = 0$$



(b) Curvilinear Translation  
 $(\alpha = 0, \omega = 0)$

### SAMPLE PROBLEM 6/1

The pickup truck weighs 3220 lb and reaches a speed of 30 mi/hr from rest in a distance of 200 ft up the 10-percent incline with constant acceleration. Calculate the normal force under each pair of wheels and the friction force under the rear driving wheels. The effective coefficient of friction between the tires and the road is known to be at least 0.80.



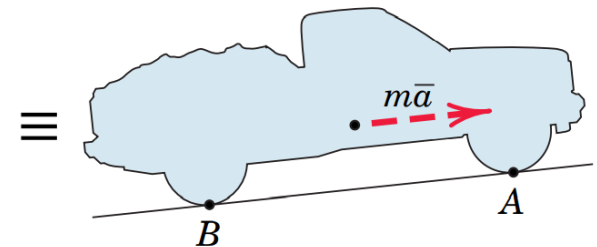
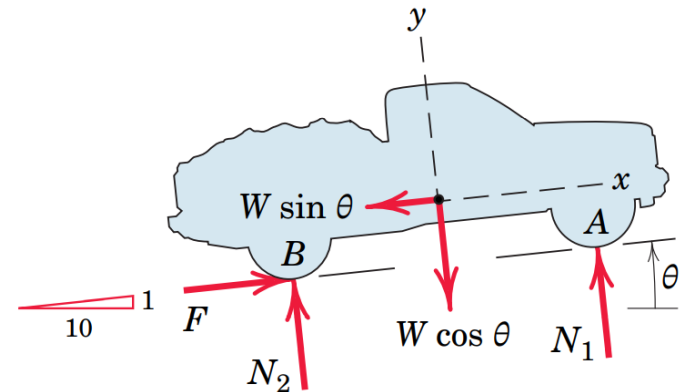
$$m\bar{a} = \frac{3220}{32.2}(4.84) = 484 \text{ lb}$$

$$[\Sigma F_x = m\bar{a}_x] \quad F - 320 = 484 \quad F = 804 \text{ lb}$$

$$[\Sigma F_y = m\bar{a}_y = 0] \quad N_1 + N_2 - 3200 = 0$$

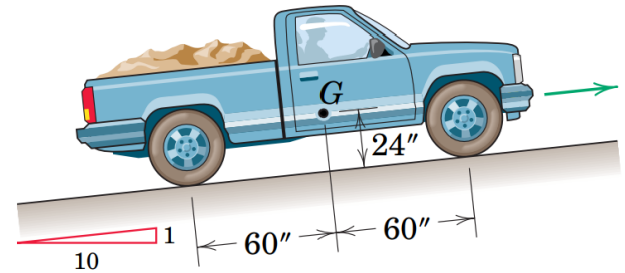
$$[\Sigma M_G = \bar{I}\alpha = 0] \quad 60N_1 + 804(24) - N_2(60) = 0$$

$$\rightarrow N_1 = 1441 \text{ lb} \quad N_2 = 1763 \text{ lb}$$



**SAMPLE PROBLEM 6/1**

The pickup truck weighs 3220 lb and reaches a speed of 30 mi/hr from rest in a distance of 200 ft up the 10-percent incline with constant acceleration. Calculate the normal force under each pair of wheels and the friction force under the rear driving wheels. The effective coefficient of friction between the tires and the road is known to be at least 0.80.

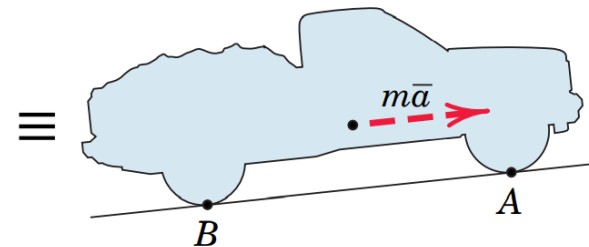
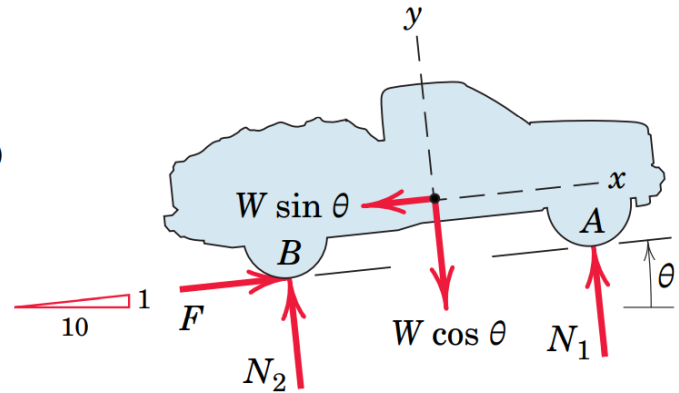
**Alternative Solution**

$$[\Sigma M_A = m\bar{a}d] \quad 120N_2 - 60(3200) - 24(320) = 484(24)$$

$$N_2 = 1763 \text{ lb}$$

$$[\Sigma M_B = m\bar{a}d] \quad 3200(60) - 320(24) - 120N_1 = 484(24)$$

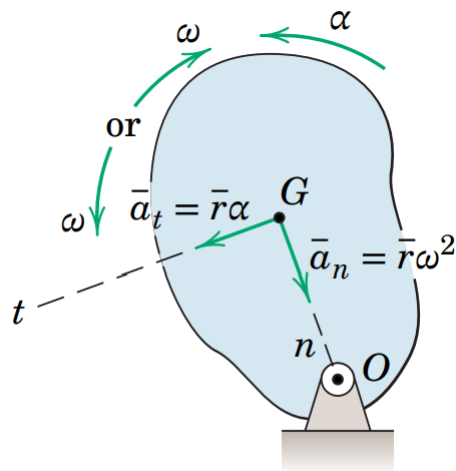
$$N_1 = 1441 \text{ lb}$$



## 6/4 Fixed-Axis Rotation

- ❖ All points in the body describe circles about the rotation axis, and all lines of the body in the plane of motion have the same angular velocity  $\omega$  and angular acceleration  $\alpha$ .

$$\Sigma M_O = \bar{I}\alpha + m\bar{a}_t\bar{r} \quad \rightarrow \quad I_O = \bar{I} + m\bar{r}^2 \quad \rightarrow \quad \Sigma M_O = (I_O - m\bar{r}^2)\alpha + m\bar{r}^2\alpha = I_O\alpha$$

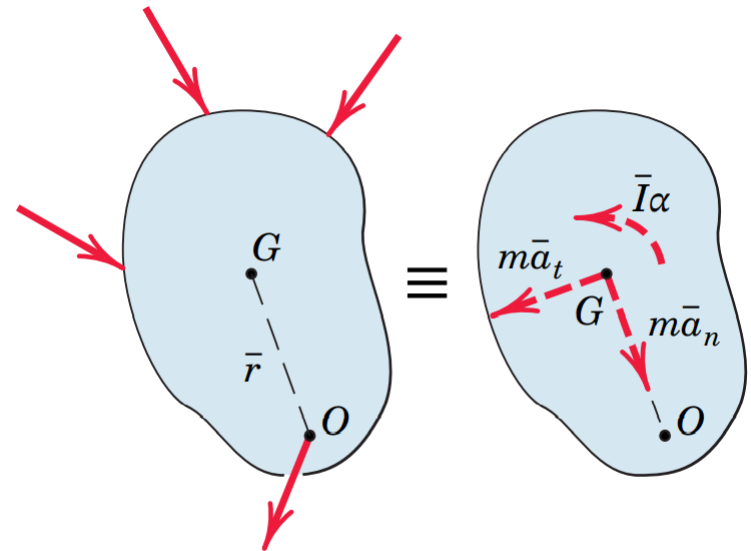


Fixed-Axis Rotation

$$\Sigma \mathbf{F} = m\bar{\mathbf{a}}$$

$$\Sigma M_G = \bar{I}\alpha$$

$$\Sigma M_O = I_O\alpha$$



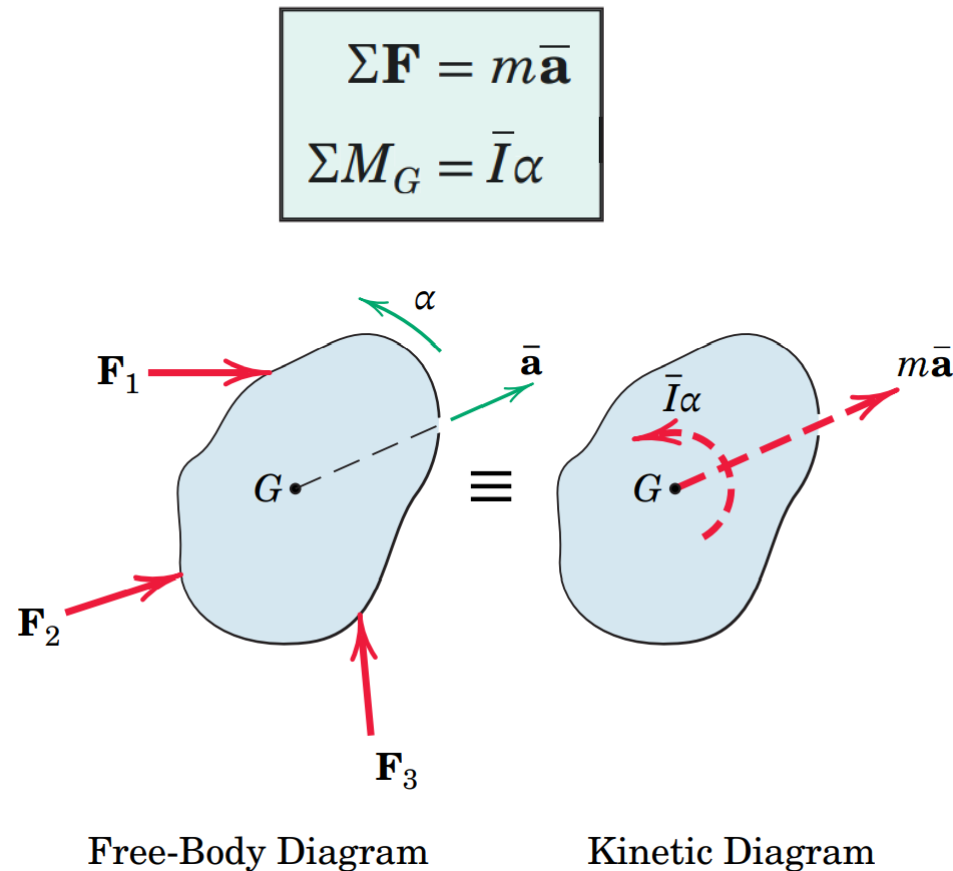
Free-Body Diagram

Kinetic Diagram

6/5

## General Plane Motion

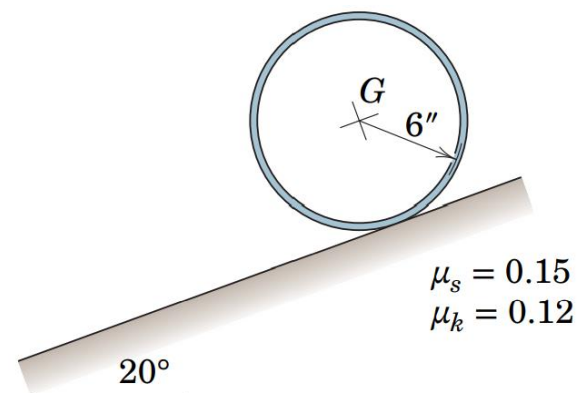
- ❖ The dynamics of a rigid body in general plane motion combines translation and rotation.





### SAMPLE PROBLEM 6/5

A metal hoop with a radius  $r = 6$  in. is released from rest on the  $20^\circ$  incline. If the coefficients of static and kinetic friction are  $\mu_s = 0.15$  and  $\mu_k = 0.12$ , determine the angular acceleration  $\alpha$  of the hoop and the time  $t$  for the hoop to move a distance of 10 ft down the incline.



Assume that the hoop rolls without slipping

$$\bar{a} = r\alpha$$

$$[\Sigma F_x = m\bar{a}_x]$$

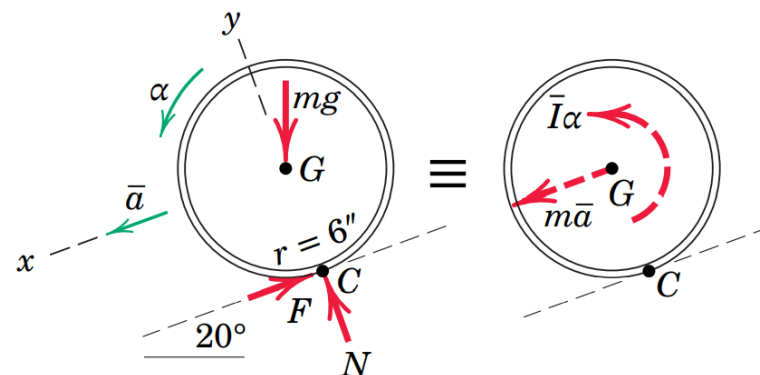
$$mg \sin 20^\circ - F = m\bar{a}$$

$$[\Sigma F_y = m\bar{a}_y = 0]$$

$$N - mg \cos 20^\circ = 0$$

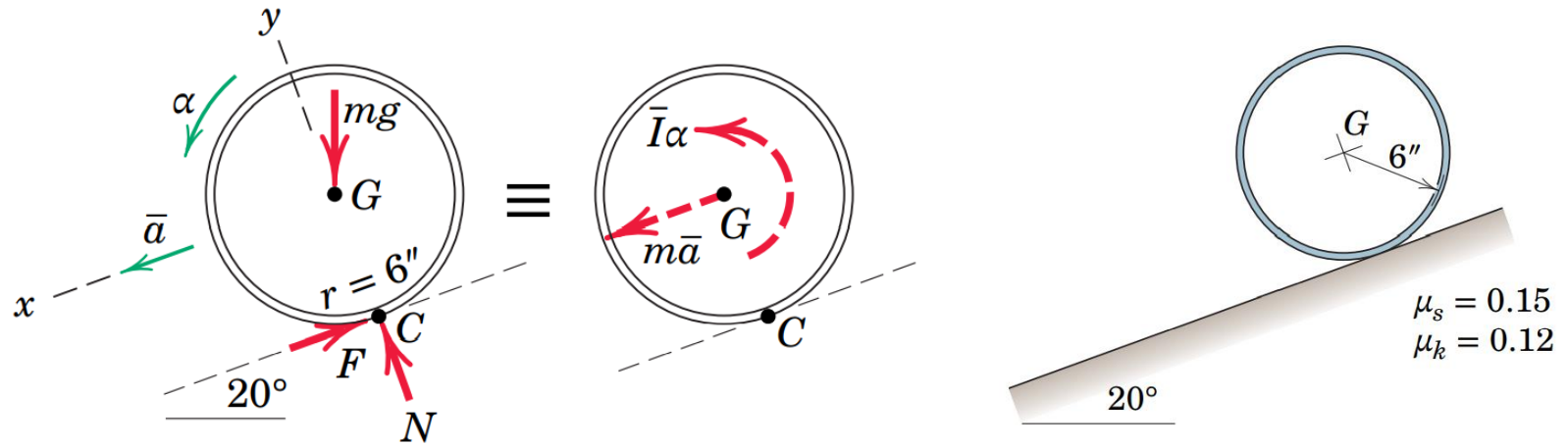
$$[\Sigma M_G = \bar{I}\alpha]$$

$$Fr = mr^2\alpha$$



$$\rightarrow \bar{a} = \frac{g}{2} \sin 20^\circ = \frac{32.2}{2} (0.342) = 5.51 \text{ ft/sec}^2$$





Alternatively:  $[\Sigma M_C = \bar{I}\alpha + m\bar{a}d] \quad mgr \sin 20^\circ = mr^2 \frac{\bar{a}}{r} + m\bar{a}r \quad \bar{a} = \frac{g}{2} \sin 20^\circ$

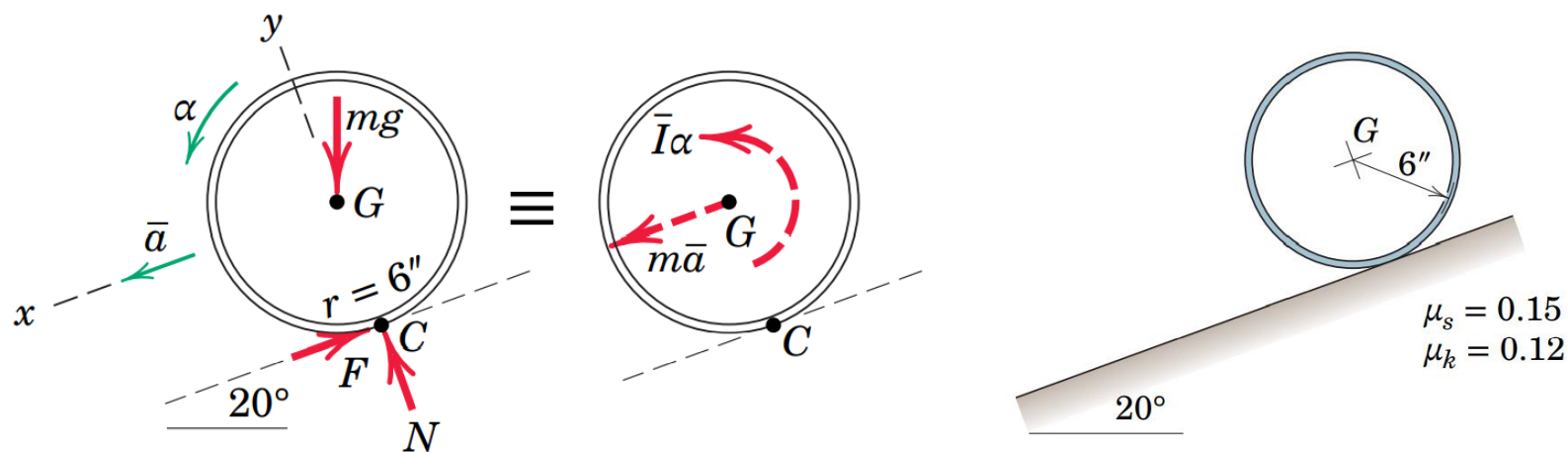
$$F = mg \sin 20^\circ - m \frac{g}{2} \sin 20^\circ = 0.1710mg$$

$$N = mg \cos 20^\circ = 0.940mg$$

→ hoop slips as it rolls

$$[F_{\max} = \mu_s N]$$

$$F_{\max} = 0.15(0.940mg) = 0.1410mg$$



$$[F = \mu_k N]$$

$$F = 0.12(0.940mg) = 0.1128mg$$

$$[\Sigma F_x = m\bar{a}_x]$$

$$mg \sin 20^\circ - 0.1128mg = m\bar{a}$$

$$\bar{a} = 0.229(32.2) = 7.38 \text{ ft/sec}^2$$

$$[\Sigma M_G = \bar{I}\alpha]$$

$$0.1128mg(r) = mr^2\alpha$$

$$\alpha = \frac{0.1128(32.2)}{6/12} = 7.26 \text{ rad/sec}^2$$

$$[x = \frac{1}{2}at^2] \quad \rightarrow \quad t = \sqrt{\frac{2x}{a}} = \sqrt{\frac{2(10)}{7.38}} = 1.646 \text{ sec}$$



## SECTION B Work and Energy

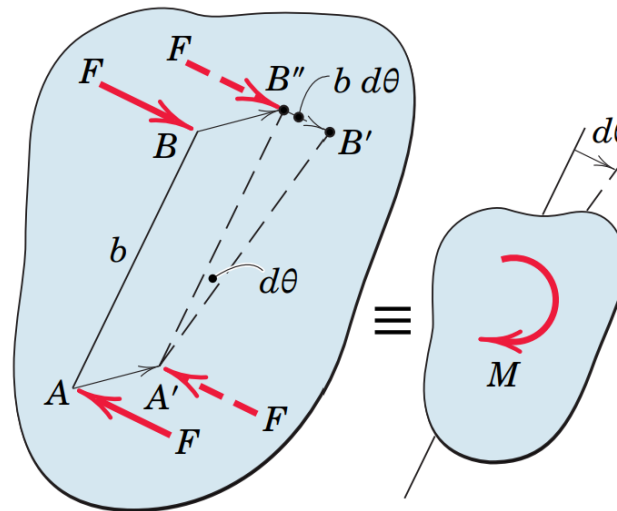
6/6

### Work-Energy Relations

#### Work of Forces and Couples

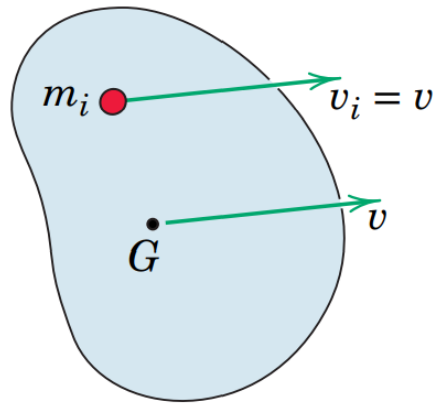
$$U = \int \mathbf{F} \cdot d\mathbf{r} \quad \text{or} \quad U = \int (F \cos \alpha) ds$$

$$U = \int M d\theta$$



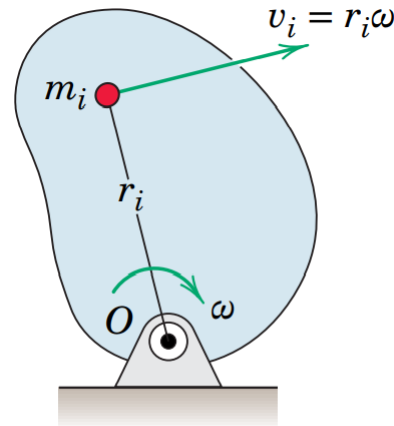
## 6/6 Work-Energy Relations

## Kinetic Energy



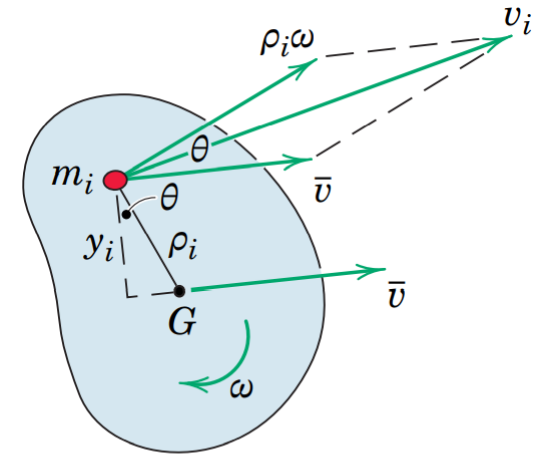
(a) Translation

$$T = \frac{1}{2}mv^2$$



(b) Fixed-Axis Rotation

$$T = \frac{1}{2}I_O\omega^2$$



(c) General Plane Motion

$$T = \frac{1}{2}m\bar{v}^2 + \frac{1}{2}\bar{I}\omega^2$$

$$T = \frac{1}{2}I_C\omega^2$$

C: instantaneous center  
of zero velocity

## 6/6 Work-Energy Relations

### Potential Energy and the Work-Energy Equation

- ❖ Gravitational potential energy  $V_g$  and elastic potential energy  $V_e$

$$T_1 + U_{1-2} = T_2$$

$$T_1 + V_1 + U'_{1-2} = T_2 + V_2$$

### Power

$$P = \frac{dU}{dt} = \frac{\mathbf{F} \cdot d\mathbf{r}}{dt} = \mathbf{F} \cdot \mathbf{v}$$

$$P = \frac{dU}{dt} = \frac{M d\theta}{dt} = M\omega$$

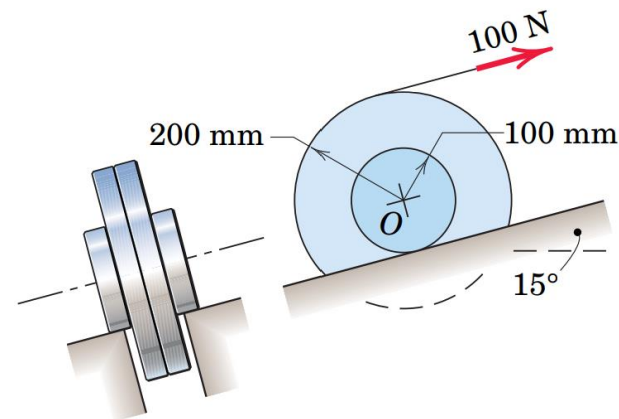


$$P = \mathbf{F} \cdot \mathbf{v} + M\omega$$



### SAMPLE PROBLEM 6/9

The wheel rolls up the incline on its hubs without slipping and is pulled by the 100-N force applied to the cord wrapped around its outer rim. If the wheel starts from rest, compute its angular velocity  $\omega$  after its center has moved a distance of 3 m up the incline. The wheel has a mass of 40 kg with center of mass at  $O$  and has a centroidal radius of gyration of 150 mm. Determine the power input from the 100-N force at the end of the 3-m motion interval.



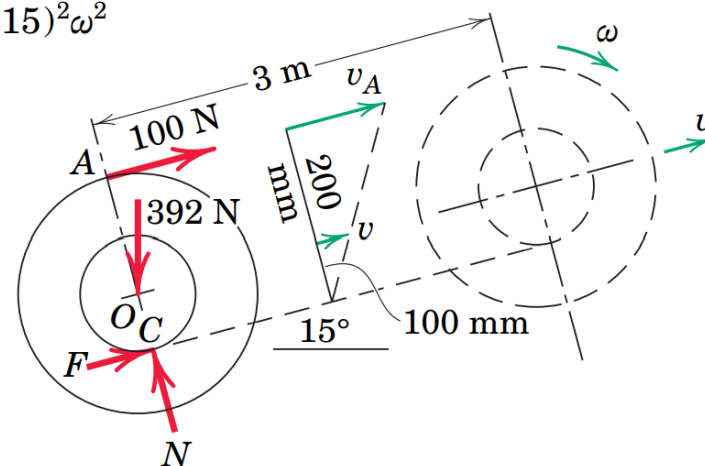
$$U_{1-2} = 100 \frac{200 + 100}{100} (3) - (392 \sin 15^\circ)(3) = 595 \text{ J}$$

$$[T = \frac{1}{2}m\bar{v}^2 + \frac{1}{2}\bar{I}\omega^2] \quad T_1 = 0 \quad T_2 = \frac{1}{2}40(0.10\omega)^2 + \frac{1}{2}40(0.15)^2\omega^2 = 0.650\omega^2$$

$$[T = \frac{1}{2}I_C\omega^2] \quad T = \frac{1}{2}40[(0.15)^2 + (0.10)^2]\omega^2 = 0.650\omega^2$$

$$[T_1 + U_{1-2} = T_2] \quad 0 + 595 = 0.650\omega^2 \quad \omega = 30.3 \text{ rad/s}$$

$$\rightarrow [P = \mathbf{F} \cdot \mathbf{v}] \quad P_{100} = 100(0.3)(30.3) = 908 \text{ W}$$



## SECTION c Impulse and Momentum

6/8

### Impulse-Momentum Equations

#### Linear Momentum

$$\mathbf{G} = m\bar{\mathbf{v}}$$

$$\Sigma \mathbf{F} = \dot{\mathbf{G}}$$

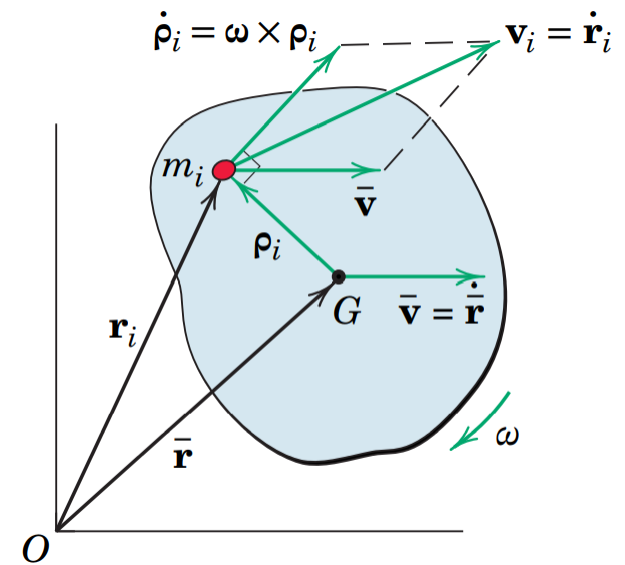
$$\mathbf{G}_1 + \int_{t_1}^{t_2} \Sigma \mathbf{F} dt = \mathbf{G}_2$$

$$\Sigma F_x = \dot{G}_x$$

$$\Sigma F_y = \dot{G}_y$$

$$(G_x)_1 + \int_{t_1}^{t_2} \Sigma F_x dt = (G_x)_2$$

$$(G_y)_1 + \int_{t_1}^{t_2} \Sigma F_y dt = (G_y)_2$$





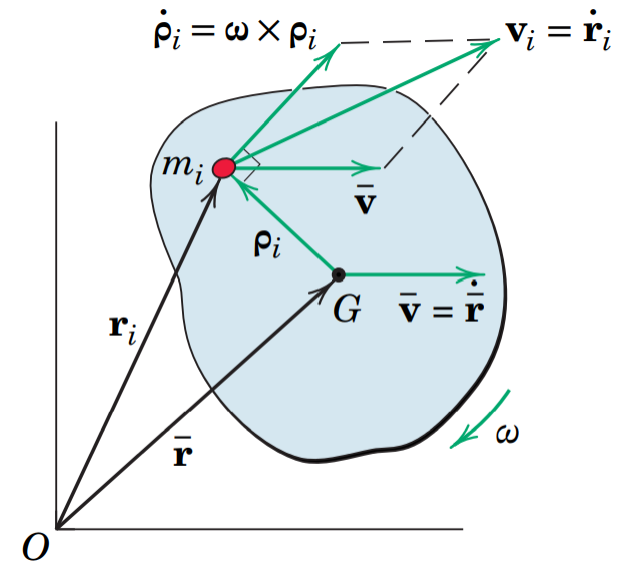
## 6/8 Impulse-Momentum Equations

### Angular Momentum

$$H_G = \bar{I}\omega$$

$$\Sigma M_G = \dot{H}_G$$

$$(H_G)_1 + \int_{t_1}^{t_2} \Sigma M_G dt = (H_G)_2$$



6/8

## Impulse-Momentum Equations

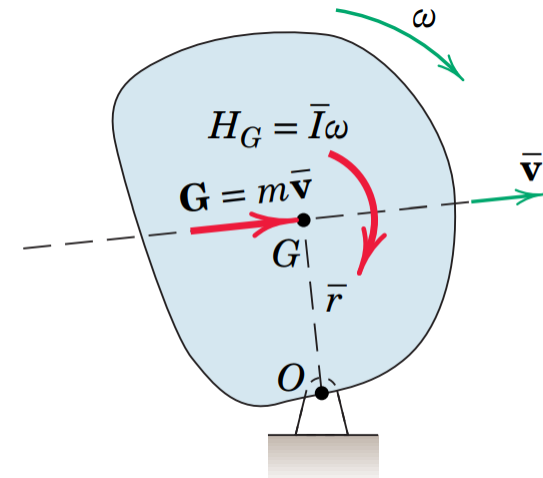
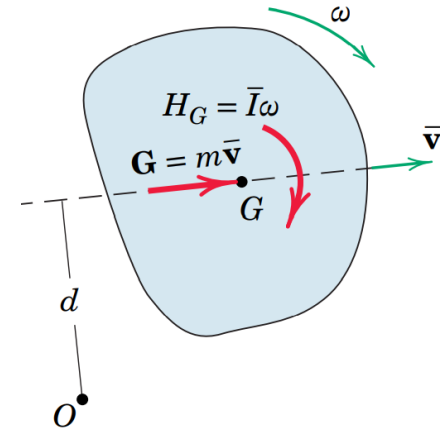
## Angular Momentum

$$H_O = \bar{I}\omega + m\bar{v}d$$

$$H_O = I_O\omega$$

$$\Sigma M_O = \dot{H}_O$$

$$(H_O)_1 + \int_{t_1}^{t_2} \Sigma M_O dt = (H_O)_2$$



6/8

## Impulse-Momentum Equations

### Conservation of Momentum

$$\mathbf{G}_1 = \mathbf{G}_2$$

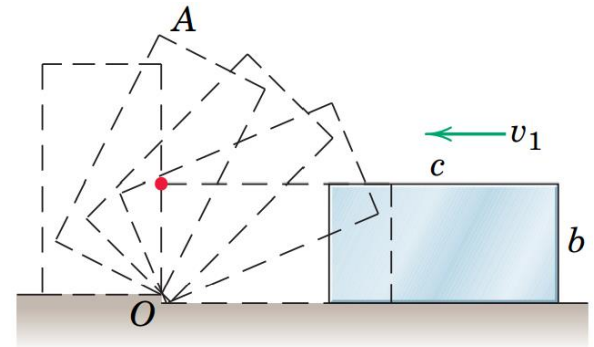
$$(\mathbf{H}_O)_1 = (\mathbf{H}_O)_2$$

$$(\mathbf{H}_G)_1 = (\mathbf{H}_G)_2$$



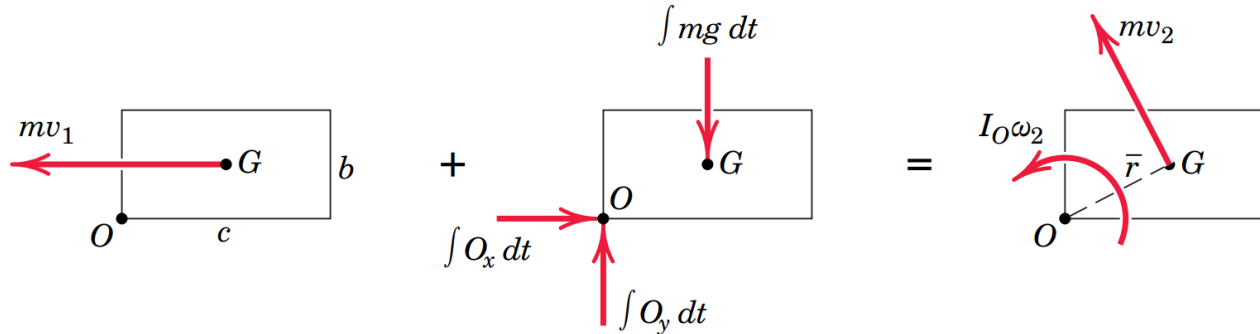
**SAMPLE PROBLEM 6/16**

The uniform rectangular block of dimensions shown is sliding to the left on the horizontal surface with a velocity  $v_1$  when it strikes the small step at  $O$ . Assume negligible rebound at the step and compute the minimum value of  $v_1$  which will permit the block to pivot freely about  $O$  and just reach the standing position  $A$  with no velocity. Compute the percentage energy loss  $n$  for  $b = c$ .



$$\begin{array}{c}
 \begin{array}{c}
 \text{Initial state:} \\
 \text{Block of width } c \text{ and height } b. \\
 \text{Center of mass } G \text{ is at } (c/2, b/2). \\
 \text{Point } O \text{ is at } (0, 0). \\
 \text{Velocity } mv_1 \text{ is directed to the left.}
 \end{array}
 +
 \begin{array}{c}
 \text{Impulse forces at } O: \\
 \int O_x dt \text{ (horizontal, right)} \\
 \int O_y dt \text{ (vertical, up)} \\
 \int mg dt \text{ (gravity, down at } G)
 \end{array}
 =
 \begin{array}{c}
 \text{Final state:} \\
 \text{Block pivoted about } O. \\
 \text{Center of mass } G \text{ is at } (c/2, b/2). \\
 \text{Velocity } mv_2 \text{ is directed up and to the right.} \\
 \text{Angular velocity } I_O \omega_2 \text{ is counter-clockwise.} \\
 \text{Position vector } \bar{r} \text{ is from } O \text{ to } G.
 \end{array}
 \end{array}$$





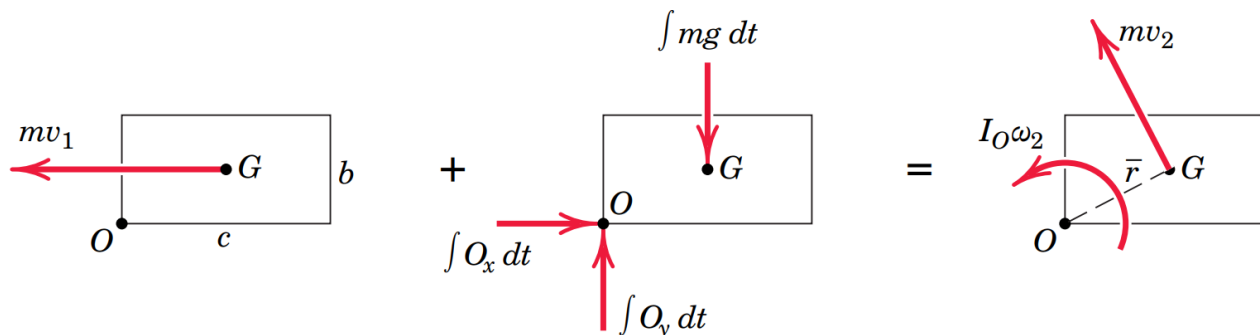
## I. Collision

$$[H_O = I_O\omega] \quad (H_O)_2 = \left\{ \frac{1}{12}m(b^2 + c^2) + m \left[ \left(\frac{c}{2}\right)^2 + \left(\frac{b}{2}\right)^2 \right] \right\} \omega_2$$

$$= \frac{m}{3}(b^2 + c^2)\omega_2$$

$$[(H_O)_1 = (H_O)_2] \quad mv_1 \frac{b}{2} = \frac{m}{3}(b^2 + c^2)\omega_2 \quad \omega_2 = \frac{3v_1 b}{2(b^2 + c^2)}$$





## II. Rotation about O

$$[T_2 + V_2 = T_3 + V_3] \quad \frac{1}{2}I_O\omega_2^2 + 0 = 0 + mg \left[ \sqrt{\left(\frac{b}{2}\right)^2 + \left(\frac{c}{2}\right)^2} - \frac{b}{2} \right]$$

$$\frac{1}{2} \frac{m}{3} (b^2 + c^2) \left[ \frac{3v_1 b}{2(b^2 + c^2)} \right]^2 = \frac{mg}{2} (\sqrt{b^2 + c^2} - b)$$

$$v_1 = 2 \sqrt{\frac{g}{3} \left( 1 + \frac{c^2}{b^2} \right) (\sqrt{b^2 + c^2} - b)}$$

$$n = \frac{|\Delta E|}{E} = \frac{\frac{1}{2}mv_1^2 - \frac{1}{2}I_O\omega_2^2}{\frac{1}{2}mv_1^2} = 1 - \frac{k_O^2\omega_2^2}{v_1^2} = 1 - \left( \frac{b^2 + c^2}{3} \right) \left[ \frac{3b}{2(b^2 + c^2)} \right]^2$$

$$= 1 - \frac{3}{4 \left( 1 + \frac{c^2}{b^2} \right)} \quad n = 62.5\% \text{ for } b = c$$